

# Extending Anisotropic Einstein-Dilaton-Maxwell Gravity for Holographic QCD

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with I.Ya. Aref'eva, V. Zlobin  
[arXiv:2604.08431](https://arxiv.org/abs/2604.08431) [hep-th]

A. Golubtsova, V. Nerovnova  
[arXiv:2604.08441](https://arxiv.org/abs/2604.08441) [hep-th]

# Black hole thermodynamics

0-th law          thermal equilibrium,  $T$           mechanical equilibrium,  $\kappa$

1-st law           $dE = TdS - pdv + \mu dN$            $dM = \frac{\kappa}{8\pi} dA + \Omega dJ + \Phi dQ$

2-nd law           $\delta S \geq 0$            $\delta A \geq 0$

3-rd law  
(in Planck form)           $S \rightarrow 0$  as  $T \rightarrow 0$            $S = \frac{A}{4} \rightarrow 0$  as  $T \rightarrow 0$  ???

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Schwarzschild BH — no 3-rd law  
SchwAdS (global coordinates) — no 3-rd law  
SchwPAdS (Poincare patch) — the 3-rd law



This is good for AdS/CFT applications



What for other models, that are interesting for holography?

# Application to holography

$$T = \left. \frac{|g'|}{4\pi} \right|_{z=z_h} \quad \longrightarrow \quad \text{QGP temperature (Maldacena)}$$

$$s = \left. \frac{1}{4} \sqrt{\prod_{i=1}^3 g_{ii}} \right|_{z=z_h} \quad \longrightarrow \quad \text{multiplicity of process (Landau)}$$

$$F = \int_{z_h}^{\infty} s dT \quad \longrightarrow \quad \text{background BH-BH phase transition}$$

$$A_t = \mu - \rho z^2 + \dots \quad \longrightarrow \quad \text{chemical potential, density}$$

# Overview

- 1 Maximal models
- 2 Particular models
  - Model I: 5D, Gauss anisotropy + Lifshitz anisotropy
  - Model II: 5D, Lifshitz anisotropy + Lifshitz anisotropy
  - Model III: 6D, Lifshitz anisotropy + Lifshitz anisotropy
- 3 D-dimensional models with electric and magnetic ansatz
- 4 Conclusions

# Holographic model of anisotropic plasma in magnetic field at nonzero chemical potential

I.Aref'eva, KR'18; IA, KR, P.Slepov'21

$$S = \int d^5x \sqrt{-g} \left[ R - \frac{f_0(\phi)}{4} F_{(0)}^2 - \frac{f_1(\phi)}{4} F_{(1)}^2 - \frac{f_B(\phi)}{4} F_{(3)}^2 - \frac{1}{2} \partial_M \phi \partial^M \phi - V(\phi) \right]$$

$$ds^2 = \frac{L^2}{z^2} \mathfrak{b}(z) \left[ -g(z) dt^2 + dx^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dy_1^2 + e^{c_B z^2} \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dy_2^2 + \frac{dz^2}{g(z)} \right]$$

$$A_{(1)\mu} = A_t(z) \delta_\mu^0 \quad A_t(0) = \mu \quad F_{(1)} = q_1 dy^1 \wedge dy^2 \quad F_{(B)} = q_3 dx \wedge dy^1$$

$$\mathfrak{b}(z) = e^{2\mathcal{A}(z)} \Leftrightarrow \text{quarks mass}$$

“Bottom-up approach”

**Heavy quarks (c, b)**

$$\mathcal{A}(z) = -cz^2/4$$

$$\mathcal{A}(z) = -cz^2/4 + p(c_B)z^4$$

*Andreev, Zakharov'06*

*IA, Hajilou, Rannu, Slepov' 23*

**Light quarks (u, d, s)**

$$\mathcal{A}(z) = -a \ln(bz^2 + 1)$$

$$\mathcal{A}(z) = -a \ln((bz^2 + 1)(dz^4 + 1))$$

*Li, Yang, Yuan'17*

*Zhu, Chen, Zhou, Zhang, Huang'25*

## Model extension

$$S = \int d^D x \sqrt{-g_D} \left[ R - \frac{f_{(0)}(\phi)}{2 \cdot 2!} F_{(0)}^2 - \sum_{\mathcal{M}} \frac{f_{(\mathcal{M})}(\phi)}{2 \cdot 2!} F_{(\mathcal{M})}^2 - \sum_{\mathcal{M}} \frac{f_{(\mathcal{M})}(\phi)}{2 \cdot 2!} \mathfrak{F}_{(\mathcal{M})}^2 - \sum_{\mathcal{N}} \frac{h_{(\mathcal{N})}(\phi)}{2 \cdot 3!} H_{(\mathcal{N})}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$ds^2 = \frac{L^2 \mathbf{b}(z)}{z^2} \left( -g(z) dt^2 + \sum_{i=1}^{D-2} \mathbf{g}_i(z) dx_i^2 + \frac{dz^2}{g(z)} \right)$$

Scalar field  $\phi = \phi(z)$ ,  $V(\phi) = V(\phi(z))$

Electric ansatz  $F_{\mu\nu}^{(0)} = \partial_\mu A_\nu^{(0)} - \partial_\nu A_\mu^{(0)}$ ,  $A_\mu^{(0)} = A_t(z) \delta_\mu^0$ ,  $\mu = \overline{0, D-1}$

Magnetic ansatz  $F_{(\mathcal{M})} = q_{(\mathcal{M})} \epsilon_{\mathcal{M}pk} dx^p \wedge dx^k$ ,  $\mathcal{M}, p, k = \overline{1, 3}$

Magnetic ansatz  $\mathfrak{F}_{(\mathcal{M})} = \mathbf{q}_{(\mathcal{M})} dx^{\mathcal{M}} \wedge dx^j$ ,  $\mathcal{M} = \overline{1, 3}$ ,  $j = \overline{4, D-2}$

Magnetic ansatz  $H_{(\mathcal{N})} = Q_{(\mathcal{N})} \epsilon_{\mathcal{N}pkl} dx^p \wedge dx^k \wedge dx^l$ ,  $\mathcal{N}, p, k, l = \overline{1, D-2}$

# Model extension

$$S = \int d^D x \sqrt{-g_D} \left[ R - \frac{f_{(0)}(\phi)}{2 \cdot 2!} F_{(0)}^2 - \sum_{\mathcal{M}} \frac{f_{(\mathcal{M})}(\phi)}{2 \cdot 2!} F_{(\mathcal{M})}^2 - \sum_{\mathcal{M}} \frac{\mathfrak{f}_{(\mathcal{M})}(\phi)}{2 \cdot 2!} \mathfrak{F}_{(\mathcal{M})}^2 - \right. \\ \left. D = 6 \quad - \sum_{\mathcal{N}} \frac{h_{(\mathcal{N})}(\phi)}{2 \cdot 3!} H_{(\mathcal{N})}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$ds^2 = \frac{L^2 \mathfrak{b}(z)}{z^2} \left( -g(z) dt^2 + \sum_{i=1}^{D-2} \mathfrak{g}_i(z) dx_i^2 + \frac{dz^2}{g(z)} \right)$$

Scalar field  $\phi = \phi(z)$ ,  $V(\phi) = V(\phi(z))$

Electric ansatz  $F_{\mu\nu}^{(0)} = \partial_\mu A_\nu^{(0)} - \partial_\nu A_\mu^{(0)}$ ,  $A_\mu^{(0)} = A_t(z) \delta_\mu^0$ ,  $\mu = \overline{0, 5}$

Magnetic ansatz  $F_{(\mathcal{M})} = q_{(\mathcal{M})} \epsilon_{\mathcal{M}pk} dx^p \wedge dx^k$ ,  $\mathcal{M}, p, k = \overline{1, 3}$

Magnetic ansatz  $\mathfrak{F}_{(\mathcal{M})} = \mathfrak{q}_{(\mathcal{M})} dx^{\mathcal{M}} \wedge dx^4$ ,  $\mathcal{M} = \overline{1, 3}$

Magnetic ansatz  $H_{(\mathcal{N})} = Q_{(\mathcal{N})} \epsilon_{\mathcal{N}pkl} dx^p \wedge dx^k \wedge dx^l$ ,  $\mathcal{N}, p, k, l = \overline{1, 4}$

# 6D with 10 magnetic ansatz: metric and dilaton equation

$$ds^2 = -g_{00}(z)dt^2 + g_{11}(z)dx_1^2 + g_{22}(z)dx_2^2 + g_{33}(z)dx_3^2 + g_{44}(z)dx_4^2 + g_{55}(z)dx_5^2$$

$$\phi'' + \frac{\phi'}{2} \sum_{\mu} (-1)^{\delta_{\mu 5}} \frac{g'_{\mu\mu}}{g_{\mu\mu}} -$$

$$- g_{55} \left( \sum_{\mathcal{K}} \frac{q_{\mathcal{K}}^2}{2} \frac{\partial f_{\mathcal{K}}}{\partial \phi} \prod_{i \neq \mathcal{K}} \frac{1}{g_{ii}} + \sum_{\mathcal{M}} \frac{q_{\mathcal{M}}^2}{2g_{\mathcal{M}\mathcal{M}g_{44}}} \frac{\partial f_{\mathcal{M}}}{\partial \phi} + \sum_{\mathcal{N}} \frac{Q_{\mathcal{N}}^2}{2} \frac{\partial H_{\mathcal{N}}}{\partial \phi} \prod_{j \neq \mathcal{N}} \frac{1}{g_{jj}} + \frac{\partial V(\phi)}{\partial \phi} \right) = 0$$

$$\mathcal{K}, \mathcal{M} = \overline{1,3}, \quad \mathcal{N} = \overline{1,4}, \quad \mu, \nu = \overline{0,5}, \quad i = \overline{1,3}, \quad j = \overline{1,4}, \quad k, l = \overline{1,5}$$

# 6D with 10 magnetic ansatz: Einstein equations

$$\begin{aligned}
 & \sum_j \left( \frac{2g''_{jj}}{g_{jj}} - \frac{g'^2_{jj}}{g^2_{jj}} \right) + \sum_{k<l} (-1)^{\delta_{5k}} (-1)^{\delta_{5l}} \frac{g'_{kk}}{g_{kk}} \frac{g'_{ll}}{g_{ll}} + \\
 & + g_{55} \left( \sum_{\mathcal{K}} f_{\mathcal{K}} q_{\mathcal{K}}^2 \prod_{i \neq \mathcal{K}} \frac{1}{g_{ii}} + \sum_{\mathcal{M}} \frac{f_{\mathcal{M}} q_{\mathcal{M}}^2}{g_{\mathcal{M}\mathcal{M}} g_{44}} + \sum_{\mathcal{N}} H_{\mathcal{N}} Q_{\mathcal{N}}^2 \prod_{j \neq \mathcal{N}} \frac{1}{g_{jj}} + \frac{\phi'^2}{g_{55}} + 2V(\phi) \right) = 0 \\
 & \sum_{\mu \neq j} (1 - \delta_{\mu 5}) \left( \frac{2g''_{\mu\mu}}{g_{\mu\mu}} - \frac{g'^2_{\mu\mu}}{g^2_{\mu\mu}} \right) + \sum_{\mu < \nu \neq j} (-1)^{\delta_{5\mu}} (-1)^{\delta_{5\nu}} \frac{g'_{\mu\mu}}{g_{\mu\mu}} \frac{g'_{\nu\nu}}{g_{\nu\nu}} + \\
 & + g_{55} \left( \sum_{\mathcal{K}} (-1)^{\delta_{\mathcal{K}j}} (-1)^{1+\delta_{4j}} f_{\mathcal{K}} q_{\mathcal{K}}^2 \prod_{i \neq \mathcal{K}} \frac{1}{g_{ii}} + \sum_{\mathcal{M}} (-1)^{\delta_{\mathcal{M}j}} (-1)^{\delta_{4j}} \frac{f_{\mathcal{M}} q_{\mathcal{M}}^2}{g_{\mathcal{M}\mathcal{M}} g_{44}} + \right. \\
 & \left. + \sum_{\mathcal{N}} (-1)^{1+\delta_{\mathcal{N}j}} H_{\mathcal{N}} Q_{\mathcal{N}}^2 \prod_{i \neq \mathcal{N}} \frac{1}{g_{ii}} + \frac{\phi'^2}{g_{55}} + 2V(\phi) \right) = 0, \quad j = \overline{1,4} \\
 & \frac{g'_{00}}{g_{00}} \sum_j \frac{g'_{jj}}{g_{jj}} + \prod_{k<l} (1 - \delta_{l5}) \frac{g'_{kk}}{g_{kk}} \frac{g'_{ll}}{g_{ll}} + \\
 & + g_{55} \left( \sum_{\mathcal{K}} f_{\mathcal{K}} q_{\mathcal{K}}^2 \prod_{i \neq \mathcal{K}} \frac{1}{g_{ii}} + \sum_{\mathcal{M}} \frac{f_{\mathcal{M}} q_{\mathcal{M}}^2}{g_{\mathcal{M}\mathcal{M}} g_{44}} + \sum_{\mathcal{N}} H_{\mathcal{N}} Q_{\mathcal{N}}^2 \prod_{j \neq \mathcal{N}} \frac{1}{g_{jj}} - \frac{\phi'^2}{g_{55}} + 2V(\phi) \right) = 0
 \end{aligned}$$

## 6D solution

$$ds^2 = \mathfrak{B}^2(z) \left( -g(z) dt^2 + \mathfrak{g}_1(z) dx_1^2 + \mathfrak{g}_2(z) dx_2^2 + \mathfrak{g}_3(z) dx_3^2 + \mathfrak{g}_4(z) dx_4^2 + \frac{dz^2}{g(z)} \right)$$

$$\phi(z) = \pm \int \sqrt{2g_{zz} \left( \frac{G_{zz}}{g_{zz}} + \frac{G_{tt}}{g_{tt}} \right)} dz$$

$$A_t(z) = \mu \left( 1 - \frac{I(z)}{I(z_h)} \right), \quad \mathfrak{s}(z) = \prod_{\gamma \neq z, t} \sqrt{g_{\gamma\gamma}(z)}, \quad I(z) = \int_0^z \frac{\mathfrak{B}^2(\xi)}{\mathfrak{s}(\xi) f_0(\phi(\xi))} d\xi$$

$$g(z) = \mathfrak{G}_i(z) \left[ C \int_{z_h}^z \frac{d\xi}{\mathfrak{s}(\xi) \mathfrak{G}_i(\xi)} + \right. \\ \left. + 2q_i^2 \int_{z_h}^z \frac{d\xi}{\mathfrak{s}(\xi) \mathfrak{G}_i(\xi)} \int_{z_h}^{\xi} \frac{\mathfrak{s}(\eta)}{\mathfrak{B}^2(\eta)} f_i(\phi(\eta)) \prod_{j=1, j \neq i}^3 \frac{1}{\mathfrak{g}_j(\eta)} d\eta \right. \\ \left. + \left( \frac{\mu}{I(z_h)} \right)^2 \int_{z_h}^z \frac{d\xi}{\mathfrak{s}(\xi) \mathfrak{G}_i(\xi)} \int_{z_h}^{\xi} \frac{\mathfrak{B}^2(\eta)}{\mathfrak{s}(\eta) f_0(\phi(\eta))} d\eta \right], \quad i = 1, 2, 3$$

Talk by V. Zlobin  $\mathfrak{G}_i = \frac{1}{\mathfrak{g}_i} \prod_{j=1, j \neq i}^3 \mathfrak{g}_j \quad \lim_{z \rightarrow 0^+} g(z) = 1$

# Particular models

$$\text{Model I} \quad D = 5 \quad F_3 = q_3 dx^1 \wedge dx^2, \quad F_1 = q_1 dx^2 \wedge dx^3$$

$$ds^2 = \frac{L^2}{z^2} \left( -g(z)dt^2 + dx_1^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dx_2^2 + e^{c_B z^2} \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dx_3^2 + \frac{dz^2}{g(z)} \right)$$



# Particular models

**Model I**  $D = 5$   $F_3 = q_3 dx^1 \wedge dx^2$ ,  $F_1 = q_1 dx^2 \wedge dx^3$

$$ds^2 = \frac{L^2}{z^2} \left( -g(z)dt^2 + dx_1^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dx_2^2 + e^{c_B z^2} \left(\frac{z}{L}\right)^{2-\frac{2}{\nu}} dx_3^2 + \frac{dz^2}{g(z)} \right)$$

**Model II**  $D = 5$  **i**  $F_2 = q_2 dx^1 \wedge dx^3$ ,  $F_3 = q_3 dx^1 \wedge dx^2$   
**ii**  $F_1 = q_1 dx^2 \wedge dx^3$ ,  $F_3 = q_3 dx^1 \wedge dx^2$   
 $\mathfrak{b}(z) = e^{cz^n}$  **iii**  $F_1 = q_1 dx^2 \wedge dx^3$ ,  $F_2 = q_2 dx^1 \wedge dx^3$

$$ds^2 = \frac{L^2}{z^2} \mathfrak{b}(z) \left( -g(z)dt^2 + dx_1^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu_1}} dx_2^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu_2}} dx_3^2 + \frac{dz^2}{g(z)} \right)$$

**Model III**  $D = 6$   $\mathfrak{b}(z) = 1$ ,  $\mathfrak{F}_3 = q_3 dx^3 \wedge dx^4$ ,  $H_1 = Q_1 dx^2 \wedge dx^3 \wedge dx^4$

$$ds^2 = \frac{L^2}{z^2} \mathfrak{b}(z) \left( -g(z)dt^2 + dx_1^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu_1}} dx_2^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu_2}} dx_3^2 + \left(\frac{z}{L}\right)^{2-\frac{2}{\nu_2}} dx_4^2 + \frac{dz^2}{g(z)} \right)$$

# Model I solution

$$g(z) = e^{c_B z^2} \frac{\int_z^{z_h} e^{-\frac{3c_B}{2}\xi^2} \xi^{1+\frac{2}{\nu}} d\xi}{\int_0^{z_h} e^{-\frac{3c_B}{2}\xi^2} \xi^{1+\frac{2}{\nu}} d\xi}, \quad 1 + \frac{1}{\nu} > 0$$

$$s(z_h) = \frac{1}{4} \left( \frac{L}{z_h} \right)^{1+\frac{2}{\nu}} e^{\frac{c_B z_h^2}{2}}, \quad T(z_h) = \begin{cases} \frac{1}{4\pi} \frac{e^{-\frac{c_B}{2}z_h^2} z_h^{1+\frac{2}{\nu}}}{\int_0^{z_h} e^{-\frac{3c_B}{2}\xi^2} \xi^{1+\frac{2}{\nu}} d\xi} \\ \frac{1}{2\pi z_h} \left( 1 + \frac{1}{\nu} \right), \quad c_B = 0 \end{cases}$$

$$f_1(\phi(z)) = -2 \left( \frac{L}{z} \right)^{4/\nu} \frac{e^{c_B z^2}}{L^2 q_1^2} \left( 1 - \frac{1}{\nu} \right) \left[ g'(z) z - \left( 2 + \frac{2}{\nu} - c_B z^2 \right) g(z) \right]$$

$$f_3(\phi(z)) = 2 \left( \frac{L}{z} \right)^{2/\nu} \frac{c_B}{q_3^2} \left[ g'(z) z - \left( \frac{2}{\nu} - c_B z^2 \right) g(z) \right]$$

$$\phi'(z)^2 = -\frac{2c_B^2}{z^2} \left( z^4 + \frac{3\nu - 2}{c_B \nu} z^2 + \frac{2 - 2\nu}{c_B^2 \nu^2} \right)$$

# Model I results

## 3-rd law

$$c_B = 0 : s(T) = \frac{1}{4} \left( \frac{2\pi L T}{1 + \frac{1}{\nu}} \right)^{1 + \frac{2}{\nu}} \rightarrow 0 \text{ as } T \rightarrow 0, 1 + \frac{2}{\nu} > 0, 1 + \frac{1}{\nu} > 0$$

$$c_B < 0 : s(T) \rightarrow 0 \text{ as } T \rightarrow 0, 1 + \frac{1}{\nu} > 0$$

$$c_B > 0 : s(z_h) \rightarrow +\infty \text{ as } T(z_h) \rightarrow 0, z_h \rightarrow +\infty$$

# Model I results

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## 3-rd law

$$c_B = 0 : s(T) = \frac{1}{4} \left( \frac{2\pi LT}{1 + \frac{1}{\nu}} \right)^{1 + \frac{2}{\nu}} \rightarrow 0 \text{ as } T \rightarrow 0, \quad 1 + \frac{2}{\nu} > 0, \quad 1 + \frac{1}{\nu} > 0$$

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## NEC

$$f_1(\phi) \geq 0, \quad f_3(\phi) \geq 0, \quad \phi'^2 \geq 0 \implies \begin{cases} c_B = 0, \nu \geq 1 \\ c_B < 0, \nu \geq 1, z_h \leq z_h^{max} \end{cases}$$

$$z_h^{max} = \sqrt{\frac{2 - 3\nu - \sqrt{9\nu^2 - 4\nu - 4}}{2c_B\nu}}$$

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$$z_h^{max} = \sqrt{\frac{2 - 3\nu - \sqrt{9\nu^2 - 4\nu - 4}}{2c_B\nu}}$$

$\Downarrow$

$$c_B < 0 : z_h \leq z_h^{max} \implies T \geq T_{min} = T(z_h^{max})$$

# Model I results

## 3-rd law

$$c_B = 0: \quad s(T) = \frac{1}{4} \left( \frac{2\pi L T}{1 + \frac{1}{\nu}} \right)^{1 + \frac{2}{\nu}} \rightarrow 0 \text{ as } T \rightarrow 0, \quad 1 + \frac{2}{\nu} > 0, \quad 1 + \frac{1}{\nu} > 0$$

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$$z_h^{max} = \sqrt{\frac{2 - 3\nu - \sqrt{9\nu^2 - 4\nu - 4}}{2c_B\nu}}$$

$\Downarrow$

$$c_B < 0: \quad z_h \leq z_h^{max} \implies T \geq T_{min} = T(z_h^{max})$$

$$c_B = 0: \quad \nu \geq 1$$

## Model II solution

$$b(z) = e^{cz^n}, \quad g_a(z) = C_a z^{\alpha_a} \int_z^{z_h} b(\xi)^{-\frac{3}{2}} \xi^{1+\frac{1}{\nu_1}+\frac{1}{\nu_2}-\alpha_a} d\xi, \quad a = \text{i, ii, iii,}$$

$$\alpha_{\text{i}} = 4 - \frac{2}{\nu_1} - \frac{2}{\nu_2}, \quad \alpha_{\text{ii}} = \frac{2}{\nu_1} - \frac{2}{\nu_2}, \quad \alpha_{\text{iii}} = \frac{2}{\nu_2} - \frac{2}{\nu_1}$$

$$s_{div, \text{i}}(T) = \frac{1}{2\pi L} \cdot \frac{1}{T}, \quad s_{div, \text{ii}}(T) = \frac{1+1/\nu_1}{4\pi L} \cdot \frac{1}{T}$$

$$c = 0, \quad s_{conv, \text{i}}(T) = \frac{1}{4} (\pi L T)^3, \quad s_{conv, \text{ii}}(T) = \frac{1}{4} \left( \frac{2\pi L T}{1 + \frac{1}{\nu_1}} \right)^{1 + \frac{2}{\nu_1}}$$

$$c < 0, \quad 0 < n < 1, \quad X = \frac{3n|c|}{8\pi T}$$

$$s_{conv, \text{i}}(T) = \frac{L^3 e^{\frac{4-n}{n-1}}}{4} X^{-\frac{3}{1-n}} \exp\left(-\frac{3|c|}{2} X^{n/(1-n)}\right) \left[1 + O\left(T^{n/(1-n)}\right)\right]$$

$$s_{conv, \text{ii}}(T) = \frac{L^{1+\frac{2}{\nu_1}} e^{\frac{2+2/\nu_1-n}{n-1}}}{4} X^{-\frac{1+2/\nu_1}{1-n}} \exp\left(-\frac{3|c|}{2} X^{n/(1-n)}\right) \left[1 + O\left(T^{n/(1-n)}\right)\right]$$

$$c > 0, \quad s_{conv, \text{i}}(T) = \frac{L^3}{16\pi} \frac{n(3c/2)^{\frac{4}{n}}}{\Gamma\left(\frac{4}{n}\right)} \cdot \frac{1}{T} + O\left(\left(\ln \frac{1}{T}\right)^{(1-n)/n}\right)$$

$$s_{conv, \text{ii}}(T) = \frac{L^{1+\frac{2}{\nu_1}}}{16\pi} \frac{n(3c/2)^{\frac{2+2/\nu_1}{n}}}{\Gamma\left(\frac{2+2/\nu_1}{n}\right)} \cdot \frac{1}{T} + O\left(\left(\ln \frac{1}{T}\right)^{(1-n)/n}\right)$$

# Model II results

## 3-rd law

$$s_{div, \mathbf{i}}(T), s_{div, \mathbf{ii}}(T) \sim \frac{1}{T} \rightarrow \infty \quad \text{as } T \rightarrow 0$$

$$c = 0: \quad s_{conv, \mathbf{i}}(T) \rightarrow 0, \quad s_{conv, \mathbf{ii}}(T) \rightarrow 0, \quad 1 + \frac{2}{\nu_1} > 0$$

$$c < 0: \quad s_{conv}(T) \rightarrow 0, \quad 0 < n < 1, \quad \mathfrak{b}(z) = e^{cz^n}$$

$$c > 0: \quad s_{conv, \mathbf{i}/\mathbf{ii}}(z_h) \rightarrow +\infty, \quad T_{conv, \mathbf{i}/\mathbf{ii}}(z_h) \rightarrow 0 \quad \text{as } z_h \rightarrow +\infty$$

# Model II results

## 3-rd law

$$s_{div, \mathbf{i}}(T), s_{div, \mathbf{ii}}(T) \sim \frac{1}{T} \rightarrow \infty \quad \text{as } T \rightarrow 0$$

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## NEC

$$c = 0: \quad \begin{cases} \mathbf{i} & \nu_1 = \nu_2 = 1 \\ \mathbf{ii}, \mathbf{iii} & \nu_1 = \nu_2 \geq 1 \end{cases} \quad c < 0: \quad \begin{cases} \mathbf{i} & \nu_1 = \nu_2 = 1 \\ \mathbf{ii}, \mathbf{iii} & \nu_1 = \nu_2 \geq 1 \end{cases}$$

$$c > 0: \quad \nu_1 = \nu_2 \geq 1, \quad z_h \leq \mathbf{min}(z_1, z_2) = \left( \frac{3(1+n) \pm \sqrt{3} \sqrt{\frac{8-8\nu_1+3(1+n)^2\nu_1^2}{\nu_1^2}}}{3cn} \right)^{1/n} > 0$$

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↓

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# Model III solution

$$g(z) = C_1 \int_z^{z_h} \frac{\xi^{1+\frac{1}{\nu_1}+\frac{2}{\nu_2}}}{b^2(\xi)} d\xi$$

$$b(z) = 1: \quad g(z) = 1 - \left(\frac{z}{z_h}\right)^{2+\frac{1}{\nu_1}+\frac{2}{\nu_2}}, \quad 2 + \frac{1}{\nu_1} + \frac{2}{\nu_2} > 0$$

$$\phi(z) = \phi_0 \pm \log\left(\frac{z}{L}\right) \sqrt{\frac{2}{\nu_1^2}(\nu_1 - 1) + \frac{4}{\nu_2^2}(\nu_2 - 1)}$$

$$h_1(\phi) = \frac{2}{L^2 Q_1^2} \left(1 - \frac{1}{\nu_1}\right) \left(2 + \frac{1}{\nu_1} + \frac{2}{\nu_2}\right) \exp\left[\left(\frac{2}{\nu_1} + \frac{4}{\nu_2}\right) \frac{\mp(\phi - \phi_0)}{\sqrt{\frac{2}{\nu_1^2}(\nu_1 - 1) + \frac{4}{\nu_2^2}(\nu_2 - 1)}}\right]$$

$$f_3(\phi) = \frac{2}{L^2 q_3^2} \left(\frac{1}{\nu_1} - \frac{1}{\nu_2}\right) \left(2 + \frac{1}{\nu_1} + \frac{2}{\nu_2}\right) \exp\left[\frac{4}{\nu_2} \frac{\mp(\phi - \phi_0)}{\sqrt{\frac{2}{\nu_1^2}(\nu_1 - 1) + \frac{4}{\nu_2^2}(\nu_2 - 1)}}\right]$$

$$V(\phi) \equiv -\frac{1}{L^2} \left(4 + \frac{4}{\nu_1} + \frac{6}{\nu_2} + \frac{1}{\nu_1^2} + \frac{2}{\nu_2^2} + \frac{3}{\nu_1 \nu_2}\right)$$

$$s(T) = 4^{\frac{1}{\nu_1} + \frac{2}{\nu_2}} \left(\frac{L \pi T}{2 + \frac{1}{\nu_1} + \frac{2}{\nu_2}}\right)^{1 + \frac{1}{\nu_1} + \frac{2}{\nu_2}}$$

# Model III results

## 3-rd law

$$s(T) \sim T^{1+\frac{1}{\nu_1}+\frac{2}{\nu_2}} \rightarrow 0, \quad 1 + \frac{1}{\nu_1} + \frac{2}{\nu_2} > 0, \quad 2 + \frac{1}{\nu_1} + \frac{2}{\nu_2} > 0$$

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## NEC

$$\begin{aligned} \phi'^2 \geq 0, \quad h_1(\phi) \geq 0, \quad \frac{h_1(\phi) Q_1^2}{g_{x_2 x_2}} + f_3(\phi) q_3^2 \geq 0 \\ \frac{1}{\nu_1} \left(1 - \frac{1}{\nu_1}\right) + \frac{2}{\nu_2} \left(1 - \frac{1}{\nu_2}\right) \geq 0, \\ \left(1 - \frac{1}{\nu_1}\right) \left(2 + \frac{1}{\nu_1} + \frac{2}{\nu_2}\right) \geq 0, \quad \left(1 - \frac{1}{\nu_2}\right) \left(2 + \frac{1}{\nu_1} + \frac{2}{\nu_2}\right) \geq 0 \end{aligned}$$

$$\nu_1, \nu_2 \geq 0 \Rightarrow \nu_1, \nu_2 \geq 1$$

$$\text{sgn}(h_1(\phi)) = \text{sgn}\left(1 - \frac{1}{\nu_1}\right) = 0, +1,$$

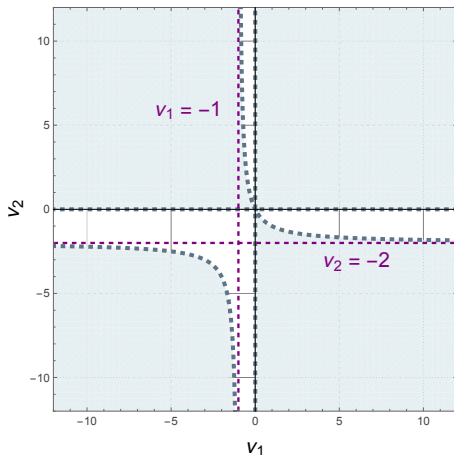
$$\text{sgn}(f_3(\phi)) = \text{sgn}\left(\frac{1}{\nu_1} - \frac{1}{\nu_2}\right) = -1, 0, +1$$

# Model III results

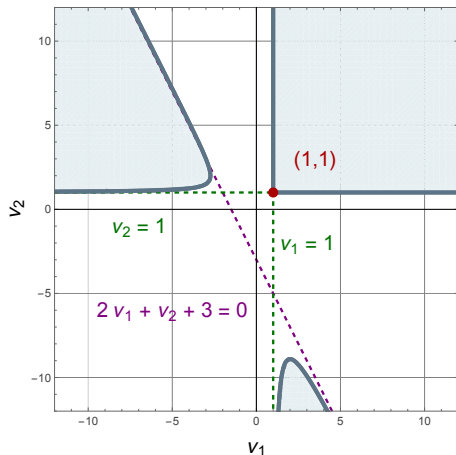
$$1 + \frac{1}{\nu_1} + \frac{2}{\nu_2} > 0$$

NEC  $\Rightarrow$  3-rd law

Third law in the  $(\nu_1, \nu_2)$  plane



NEC in the  $(\nu_1, \nu_2)$  plane



# Results

## Model I $D = 5$ Lifshitz + Gauss

NEC  $c_B \leq 0, \nu \geq 1$ :  $c_B < 0, z_h \leq z_h^{max}, T_{min} = T(z_h^{max}) > 0$

$c_B = 0$  horizon unbounded

3-rd law + NEC  $c_B = 0, \nu \geq 1$

## Model II $D = 5$ 2 Lifshitz

NEC **i**  $\nu_1 = \nu_2 = 1$ , **ii, iii**  $\nu_1 = \nu_2, z_h^{max}$ :  $\nu_1 = \nu_2 > 1, c > 0$

3-rd law + NEC  $\left\{ \begin{array}{ll} \mathbf{i} & \nu_1 = \nu_2 = 1, \quad c < 0, \quad 0 < n < 1 \\ \mathbf{ii, iii} & \nu_1 = \nu_2 \geq 1, \quad c < 0, \quad 0 < n < 1 \end{array} \right.$

## Model III $D = 6$ 2 Lifshitz

NEC  $\Rightarrow$  3-rd law:  $\mathfrak{b} = 1, 1 + \frac{1}{\nu_1} + \frac{2}{\nu_2} > 0$

$\mathfrak{b} \neq 1, ???$

# D-dimensional models: electric + magnetic ansatz

## Electric and 2-form magnetic ansatz

$$S = \int d^D x \sqrt{-g_D} \left( R - \frac{f_{(0)}(\phi)}{4} F_{(0)}^2 - \frac{f_{(1)}(\phi)}{4} F_{(1)}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right)$$

$$ds^2 = \frac{L^2 b(z)}{z^2} \left( -g(z) dt^2 + \sum_{i=1}^d dx_i^2 + z^{2-\frac{2}{\nu}} \sum_{j=1}^2 dy_j^2 + \frac{dz^2}{g(z)} \right)$$

## 2- and 3-forms magnetic ansatz

$$S = \int d^D x \sqrt{-g_D} \left( R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{4} f_{(2)}(\phi) F_{(2)}^2 - \frac{1}{4} f_{(H)}(\phi) H^2 - V(\phi) \right)$$

$$ds^2 = \frac{L^2 b(z)}{z^2} \left( -g(z) dt^2 + \sum_{i=1}^d dx_i^2 + \frac{e^{c_B z^2}}{(z/L)^{\frac{2}{\nu}-2}} \sum_{j=1}^3 dy_j^2 + \frac{dz^2}{g(z)} \right)$$

Scalar field  $\phi = \phi(z)$ ,  $V(\phi) = V(\phi(z))$

Electric ansatz  $F_{\mu\nu}^{(0)} = \partial_\mu A_\nu^{(0)} - \partial_\nu A_\mu^{(0)}$ ,  $A_\mu^{(0)} = A_t(z) \delta_\mu^0$ ,  $\mu = \overline{0, D-1}$

Magnetic ansatz  $F_{(1)} = q dy_1 \wedge dy_2$ ,  $F_{(2)} = q dy_2 \wedge dy_3$ ,  $H = q_H dy_1 \wedge dy_2 \wedge dy_3$

[arXiv:2604.08441 \[hep-th\]](https://arxiv.org/abs/2604.08441)

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Scalar field  $\phi = \phi(z)$ ,  $V(\phi) = V(\phi(z))$  Analytic solutions

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# 3-rd law conditions

## Electric and 2-form magnetic ansatz

$$\mathbf{b}(z) = 1$$

$$F_{(0)}^2 = 0: \quad s(T) = \frac{1}{4} \left( \frac{\frac{2}{\nu} + D - 3}{4\pi T} \right)^{4-D-\frac{2}{\nu}}, \quad D \geq 5, \nu \geq 1, 4 - D - \frac{2}{\nu} < 0$$

$$F_{(0)}^2 \neq 0: \quad s \rightarrow 0 \text{ as } T \rightarrow 0, \quad f_0 = e^{k(\phi-\phi_0)} = z^\kappa, \quad \kappa < -1$$

$$\mathbf{b}(z) = e^{-cz^2}, \quad D = 5: \quad c = 0$$

## 2- and 3-forms magnetic ansatz

$$\mathbf{b}(z) = e^{-cz^2}$$

$$s(z_h) = \frac{1}{4} \left( \frac{z_h}{L} \right)^{5-D-\frac{3}{\nu}} e^{\frac{z_h^2}{2}(c_B - c(D-2))}, \quad c_B = c(D-2), \nu \geq 1, D \geq 5$$

$$\mathbf{b}(z) = 1: \quad g(z) = 1 + C z^{1+\frac{3}{\nu}} - \left( \frac{z}{z_h} \right)^{2+\frac{4}{\nu}} \left( 1 + C z_h^{1+\frac{3}{\nu}} \right), \quad C = 0; \quad V = 0$$

arXiv:2604.08441 [hep-th]

# Conclusions

- Analytic solution can be obtained in quadratures
- In 5D with Gauss 3-rd law fulfills for  $\nu \geq 1$ ,  $c_B = 0$
- Magnetic anisotropy in Gauss form  $c_B > 0$  spoils the 3-rd law
- Anisotropy spoils NEC for 2 Lifshitz
- In 5D with Gauss  $c_B < 0$  and for 2 Lifshitz  $c < 0$ ,  $n \geq 1$  spoil the 3-rd law for NEC fulfillment
- 6D with 2 Lifshitz allows 3-rd law and NEC for  $\nu < 1$ ,  $\nu_1 \neq \nu_2$ , but for  $\mathfrak{b} = 1$

In 5D Gaussian warp-factor and anisotropy affect the 3-rd law

D > 5 models can solve these problems  
for **some** warp-factors and anisotropies

Thank you  
for your attention