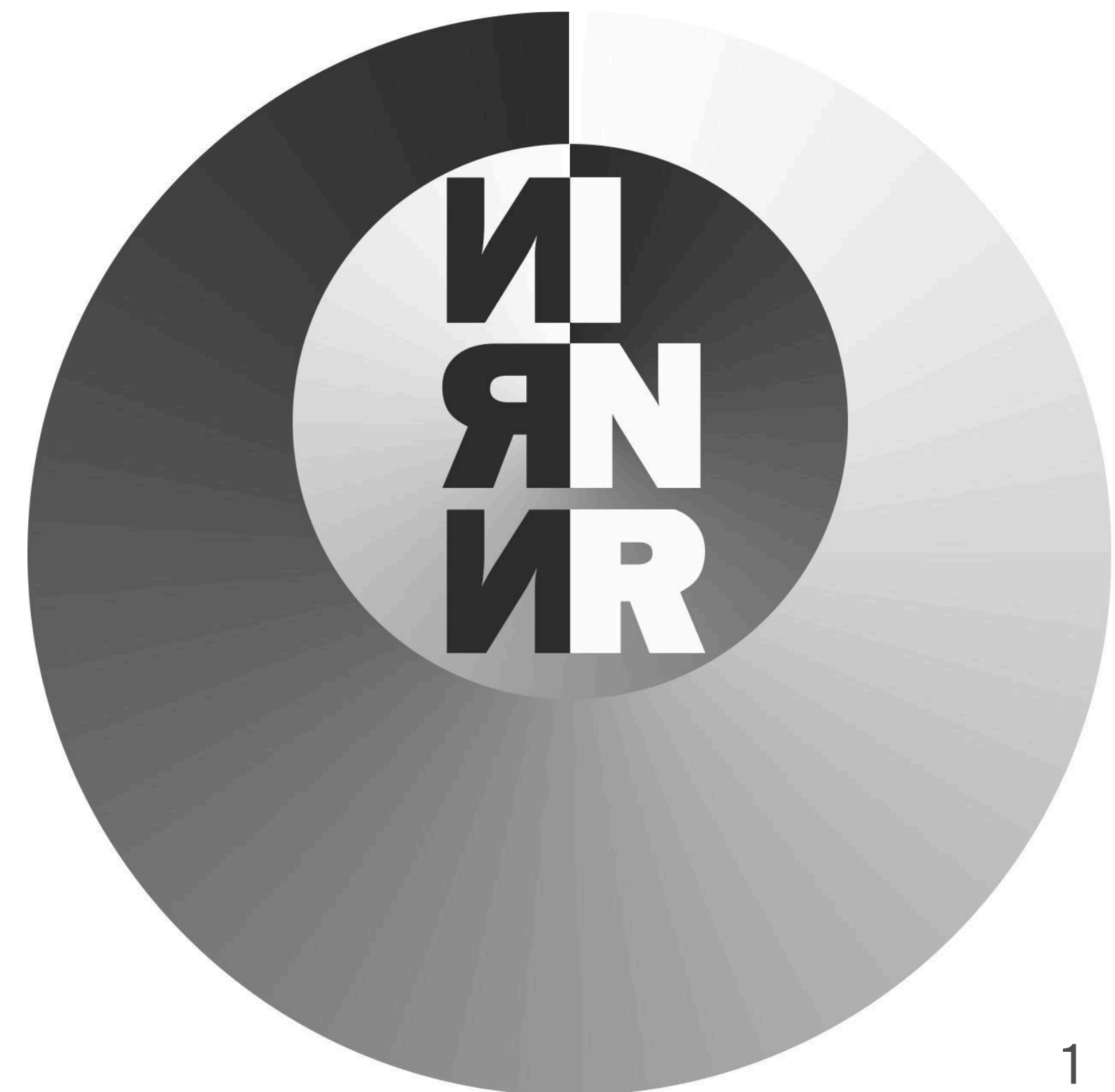


Axion–neutrino interactions in seesaw models and astrophysical probes

Polina Kivokurtseva, INR RAS



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Axion and ALP



Hypothetical pseudo-Nambu–Goldstone bosons that arise in a variety of extensions of the Standard Model, most famously in the Peccei–Quinn solution to the strong-CP problem.

Extensions of Standard Model

Massive neutrinos

Seesaw I

Inverse seesaw

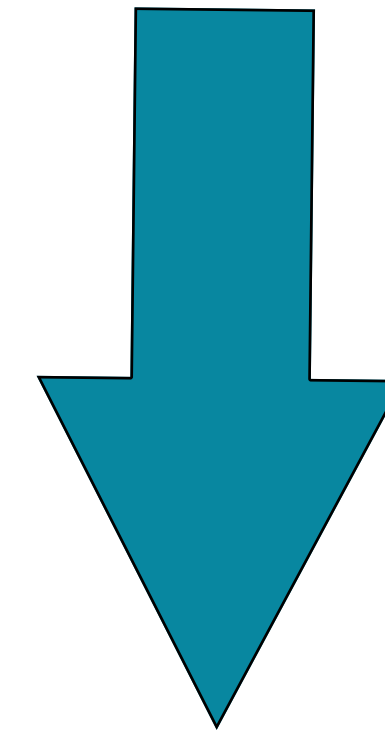
m_ν

Axion and ALP

$$\mathcal{L} \supset i g_{af} a \bar{f} \gamma_5 f,$$

$$g_{af} \propto \frac{m_f}{f_a}$$

Neutrino oscillations demonstrate that neutrinos are massive, implying that an axion–neutrino coupling is generically allowed in mass-proportional frameworks.



Can we detect the result of such interaction in different astrophysical scenarios ?

Seesaw I

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + i\bar{N}_R \partial N_R - (Y_\nu)_\alpha \bar{L}_\alpha \tilde{H} N_R - \frac{1}{2} (\mathcal{M}_R) \overline{N_R^c} N_R + \mathcal{L}_{aNN} + \mathcal{L}_{ae} + \text{h.c.}$$



$$\mathcal{L}_{aNN} = \frac{C_\nu}{f_a} (\partial_\mu a) \bar{N}_R \gamma^\mu \gamma_5 N_R = -\frac{2iC_\nu}{f_a} m_{N_R} a \bar{N}_R \gamma_5 N_R$$

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & \mathcal{M}_D \\ \mathcal{M}_D^\top & \mathcal{M}_R \end{pmatrix}$$

$$\Gamma(a \rightarrow \nu\nu) = \frac{m_N^2 m_a |U_{\nu N}|^4}{2\pi f_a^2} \sqrt{1 - \frac{4m_\nu^2}{m_a^2}}$$

$$U_{N\nu} \simeq \mathcal{M}_D \mathcal{M}_R^{-1}$$

Inverse Seesaw

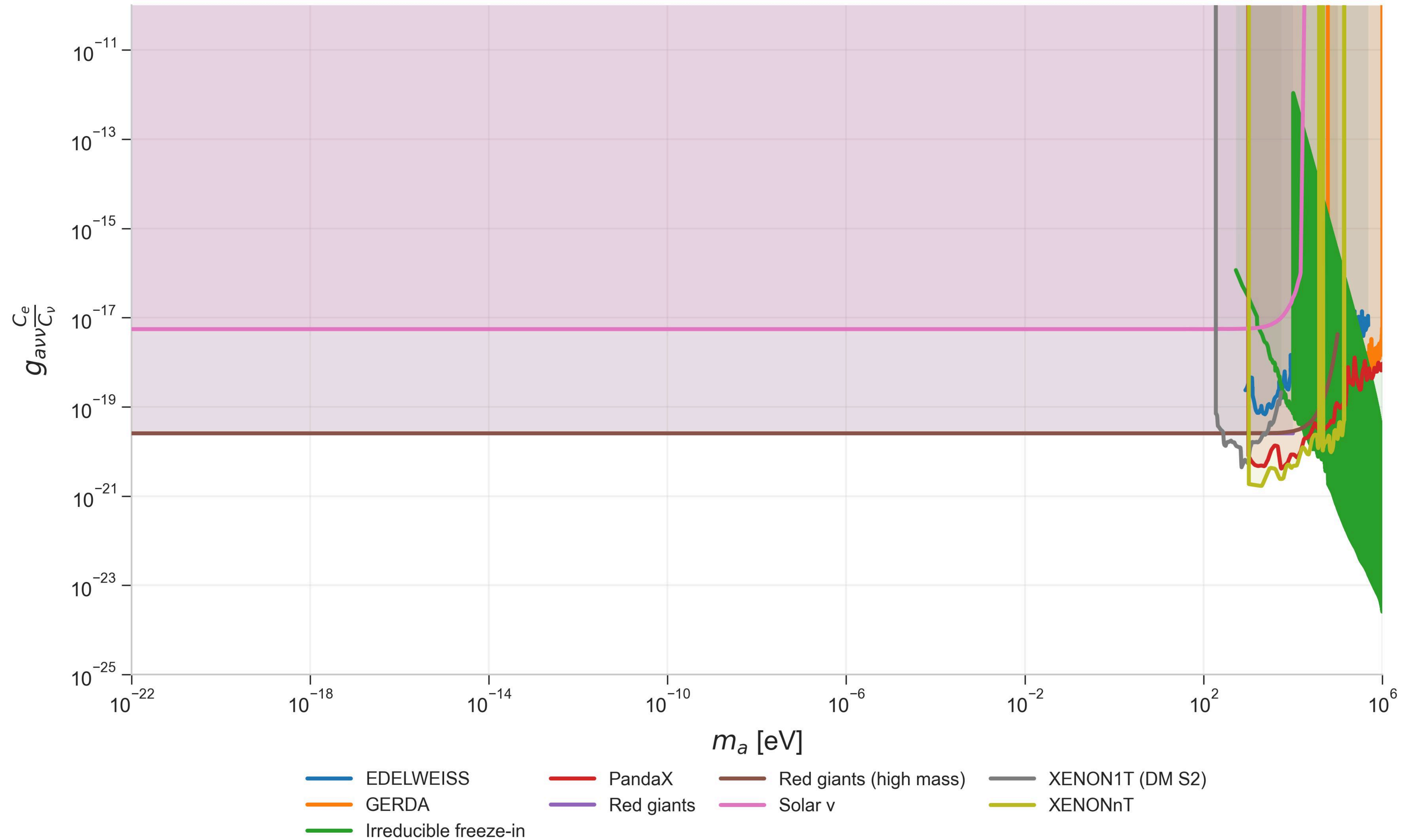
$$\mathcal{L}_{\text{mass}} = -\overline{N_R} m_D \nu_L - \overline{S_R} M N_R^c - \frac{1}{2} \overline{S_R} \mu S_R^c + \text{h.c.}$$

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m_D^\top & 0 \\ m_D & 0 & M^\top \\ 0 & M & \mu \end{pmatrix} \quad \begin{aligned} m_\nu &\simeq \mu \frac{m_D^2}{M^2}, \\ m_\pm &\simeq \pm \sqrt{M^2 + m_D^2} + \mu \frac{M^2}{2(M^2 + m_D^2)}. \end{aligned}$$

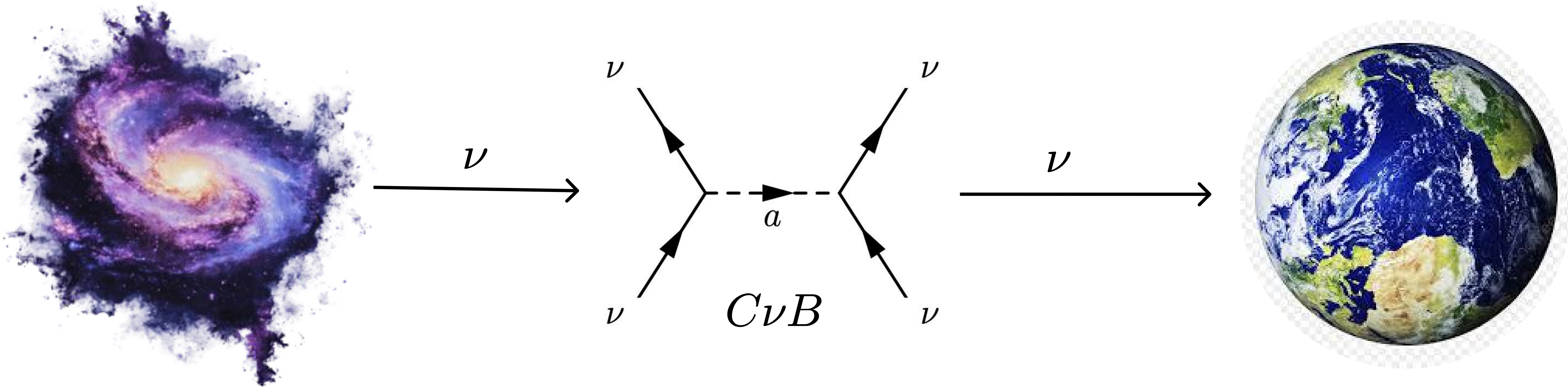
$$\mathcal{L}_{a\psi\psi} = \frac{1}{f_a} (\partial_\mu a) \bar{\psi} \gamma^\mu \gamma_5 Q \psi = \frac{1}{f_a} (\partial_\mu a) \bar{\nu} \gamma^\mu \gamma_5 G \nu$$

ALP constraints

$$g_{a\nu} = g_{aee} \frac{C_\nu}{C_e} \frac{m_\nu}{m_e}$$



Time delay



$$\tau = n \sigma_{\nu\nu} D$$

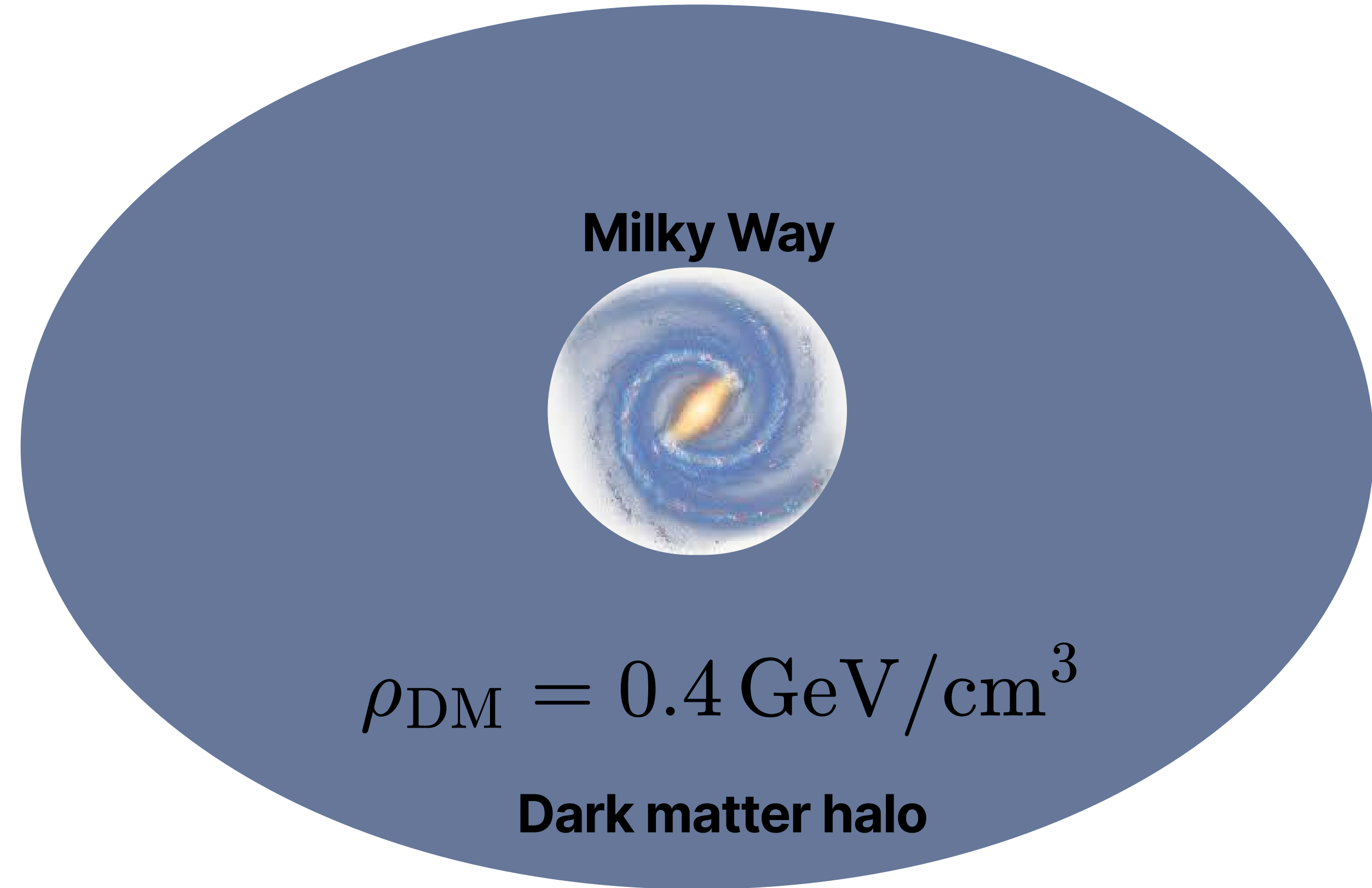
$$\tau \approx 10^{-53}$$

The probability for a neutrino to scatter on a $C\nu B$ neutrino along the line of sight is $1 - e^{-\tau}$

Axion flux



ν



Milky Way

$$\rho_{\text{DM}} = 0.4 \text{ GeV/cm}^3$$

Dark matter halo

$$\tau = n \sigma_{\nu\nu} D$$

$$\tau \approx 10^{-15}$$

Conclusions

- **Within minimal and phenomenologically viable Type-I and inverse-seesaw models, none of our benchmark scenarios can yield an observable effect**

- **Detectable signals would require either non-minimal model-building that decouples g_{ν} from charged-lepton constraints and astrophysical settings with drastically enhanced effective target densities**

Thank you!