

Axion-like dark matter and Bose stars



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Light Axion-Like Particles (ALP)

$$a(x)$$

- (pseudo)scalars
 - low mass m_a
- } Pseudo-Nambu-Goldstone bosons

Low energies: $\Phi = f_a \cdot e^{i\theta(x)}$

$$\mathcal{L}_{\text{eff}} \sim \underbrace{f_a^2 (\partial_\mu \theta)^2}_{\text{Goldstone theorem}} - \underbrace{\Lambda^4 \mathcal{V}(\theta)}_{\text{explicit breaking}}$$

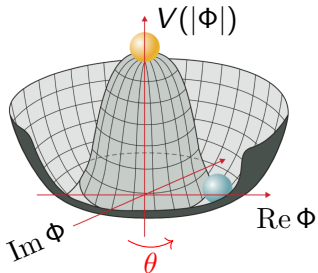
theory scale breaking scale
 \downarrow \downarrow

• $\theta = a(x)/f_a$ — canonical field

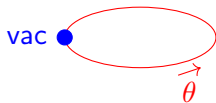
$$\mathcal{V} = \underbrace{(\Lambda^2/f_a)^2}_{m_a^2 - \text{mass}} a^2 - \underbrace{(m_a/f_a)^2}_{\lambda - \text{coupling}} a^4 + \dots$$

• Interactions are suppressed! ($\theta = a/f_a$):

$$\mathcal{L}_{\text{eff}} \supseteq \underbrace{g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}}_{f_a^{-1} \times \text{photons}} + \underbrace{g_{a\psi} \partial_\mu a \bar{\psi} i \gamma^\mu \gamma^5 \psi}_{f_a^{-1} \times \text{fermions}}$$



$$\text{Low E: } \Phi = f_a \cdot e^{i\theta(x)}$$



• Breaking: 1 vacuum

Popular models

QCD axions

Peccei, Quinn '77

- Strong CP–problem:

$$\Delta\mathcal{L}_{\text{QCD}} \sim \theta G_{\mu\nu} \tilde{G}^{\mu\nu}, \underbrace{\theta < 10^{-10}}_{\text{experiment}}$$

- Solution — field $\Phi(x)$!

low E : $\Phi = f_a \cdot e^{i\theta(x)}$

explicit breaking \equiv anomaly

$$\mathcal{V} \sim \theta(x) G_{\mu\nu} \tilde{G}^{\mu\nu}$$

$$\langle \mathcal{V} \rangle_{\text{QCD}} \sim f_\pi^2 m_\pi^2 [1 - \cos \theta]$$

- Axion field: $\theta = a(x)/f_a$

mass: $m_a \sim f_\pi m_\pi / f_a$

interactions: *e.g. di Cortona et al '16*

$$g_{a\gamma\gamma} = c f_a^{-1}, \quad 10^{-4} \lesssim c \lesssim 1$$

$$g_{ap} \sim g_{an} \sim f_a^{-1}$$

String axions

e.g. Arvanitaki et al '10

- Strings live in

$$10d = 4d \times$$



6d Calabi–Yau

- Axions = CY «symmetries» $\theta_i(x)$

typically ~ 30 axions

scale: $f_a \equiv a/\theta \sim 10^{-2} M_{\text{pl}}$

explicit breaking \equiv instantons

- Interactions are tiny!

mass: $m_a \propto \underbrace{e^{-S_{\text{inst}}/2}}_{\text{arbitrary!}}$

$$g_{a\gamma\gamma} \sim g_{ap} \sim g_{an} \propto f_a^{-1}$$

$$\lambda \sim m_a^2 / f_a^2 \sim 10^{-100}$$

Light axions form good dark matter!

- They are long-living:

$$\Gamma \left[a \rightarrow \gamma \gamma \right] \sim \underbrace{g_{a\gamma\gamma}^2 m_a^3}_{\propto m_a^5} \ll (10^{10} \text{ yr})^{-1}$$

$\Rightarrow m_a \ll 100 \text{ eV}$ (QCD axions)

- They should fit into galaxies:

$$(m_a v)^{-1} \lesssim \text{kpc} \Rightarrow m_a \gtrsim 10^{-22} \text{ eV}$$

- Quantum («fuzzy») dark matter:

$$(m_a v)^{-1} \sim \text{kpc} \Rightarrow m_a \sim 10^{-22} \text{ eV}$$

- Large occupation numbers:

$$m_a \ll 100 \text{ eV} \Rightarrow f \sim N_a / (R m_a v)^3 \gg 1$$

thermalization \Rightarrow Bose-Einstein condensation

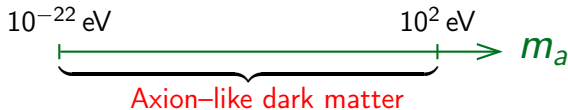
Fornax galaxy



$$\rho_{dm} \sim 0.1 M_{\odot} / \text{pc}^3$$

$$R \sim \text{kpc}$$

$$v \sim 10 \text{ km/s}$$



Generation: vacuum realignment at $f_a > H_{\text{infl}}$

- Works for **string axions** & **low-mass QCD axions**
- $\Phi = f_a \cdot e^{ia(t, \mathbf{x})/f_a}$ **prior to inflation**
- After inflation: $a = a(t)$

$$\underbrace{\ddot{a} + m_a^2 a}_{\text{pendulum}} + \underbrace{3H(t) \dot{a}}_{\substack{\text{friction} \\ H \propto t^{-1}}} = 0$$

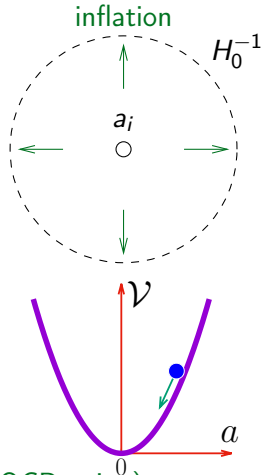
- Solution: $a \sim \begin{cases} \text{const}, & m_a < 3H \\ a_i t^{-3/4} \cos(m_a t), & m_a > 3H \end{cases}$

waves (particles) with $\mathbf{p} = 0$

- Present-day density: $\rho_a \propto \underbrace{a_i^2 m_a^2}_{\text{fixes } a_i} \sim 25\% \rho_c$

Antropic window: $\underbrace{a_i}_{\text{fine-tuning}} \ll 2\pi f_a$ if $m_a \ll 2 \cdot 10^{-5} \text{ eV}$ (QCD axion)

- Constraints from inflation: $H_{\text{infl}} \lesssim 10^{14} \text{ GeV} \left(\frac{10^{-11} \text{ eV}}{m_a} \right)^{1/4}$

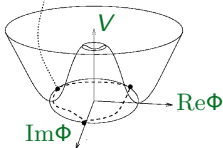
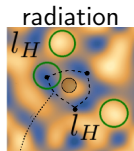


Generation: vacuum realignment at $f_a < H_{\text{infl}}$

- **Spontaneous breaking:** $\Phi = f_a \cdot e^{ia(x)/f_a}$ at $T \lesssim f_a$
 - volumes $l_H = H^{-1}$ are **causally disconnected**
 - ⇒ **different $a(x)$** inside each l_H
 - ⇒ strings of **thickness f_a^{-1}**
- $m_a < 3H$ — **massless $a(x)$**
 - ⇒ mix inside l_H
- $m_a \gtrsim 3H$ — **massive $a(x)$** & a single vacuum $a = 0$
 - **dust:** $a \propto \cos(m_a t)$
 - ⇒ **inhomogeneities of size l_H :** $m_a \sim 3H$!
 - **but also:** domain walls form $0 \rightarrow a \rightarrow 2\pi f_a$
strings & walls **annihilate!**

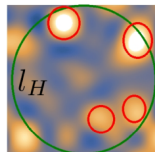
Kibble '76

$m_a < 3H(t)$



$m_a > 3H(t)$

dust = «sand»



inhomogeneities

Inhomogeneities ⇒ axion miniclusters *Kolb, Tkachev '94*



grav/bound

Now (QCD axions): 70% in miniclusters,

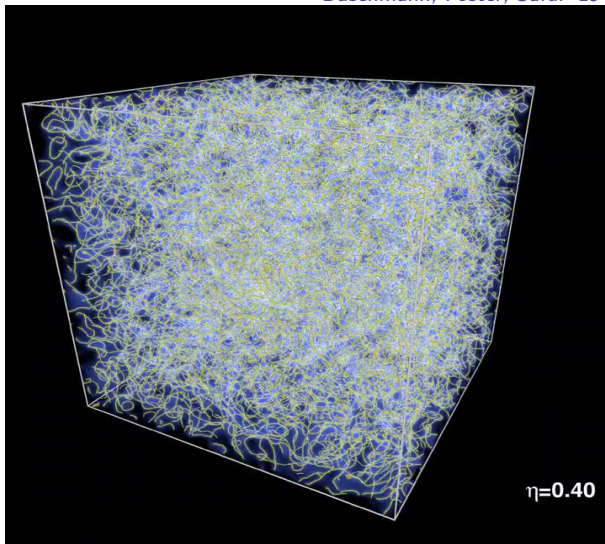
$R_{\text{mc}} \sim 10 \div 500 \text{ AU}$, $M_{\text{mc}} \sim 10^{-(10 \div 18)} M_{\odot}$

Pierobon et al '23

- **QCD axions:** $T \sim \text{GeV}$ (QCD phase transition)

Simulation of QCD phase transition

Buschmann, Foster, Safdi '19



Simulations:

Kolb, Tkachev '94

Klaer, Moore '17

Buschmann et al '19 '22

Ghorghetto et al '20

$$\rho_a = 25\% \rho_c$$



$$m_a = \underbrace{10^{-5} \dots 10^{-3}} \text{ eV}$$

Systematics: $\frac{f_a}{H} \sim 10^{30}$



minicluster uncertainties

Bose–Einstein condensation via gravitational scattering

DL, Panin, Tkachev '18

- Gravitational relaxation:

Rutherford cross section

$$\sigma_{gr} = \left| \begin{array}{c} \phi \text{---} \phi \\ \updownarrow g \\ \phi \text{---} \phi \end{array} \right|^2 \propto \frac{(m_a G)^2}{v^4}$$

$$t_{gr} \sim \frac{1}{\sigma_{gr} v n f}$$

Bose amplification

(phase density)

$$f \propto \frac{n}{m_a^3 v^3}$$

large factors at $v \ll 1, f \gg 1$

- «Fuzzy» DM in dwarf galaxies:

$$m_a \sim 10^{-22} \text{ eV} \Rightarrow t_{gr} \gtrsim 10^6 \text{ yr}$$

$$\text{Ly-}\alpha \text{ bound: } m_a \gtrsim 2 \cdot 10^{-20} \text{ eV} \Rightarrow t_{gr} \gtrsim 10^{14} \text{ yr}$$

Rogers, Peiris '20

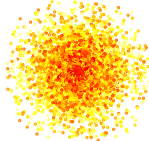
dwarf galaxy



- QCD axions in miniclusters:

$$m_a \sim 10^{-4} \text{ eV}, \quad t_{gr} \gtrsim \text{hr}$$

miniclusters



- $f \gg 1$ $\left\{ \begin{array}{l} \text{relaxation} \Rightarrow \text{Bose condensation} \\ \text{classical field } a(x) \end{array} \right.$

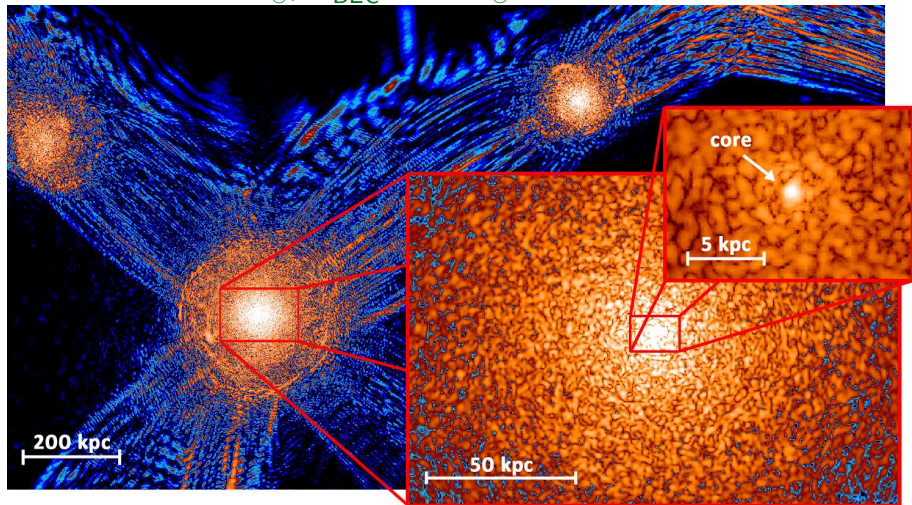
Bose–Einstein Condensate in miniclusters!

Condensate in «fuzzy» dark matter, $m_a \sim 10^{-22}$ eV

Dwarf galaxy \sim Fornax:

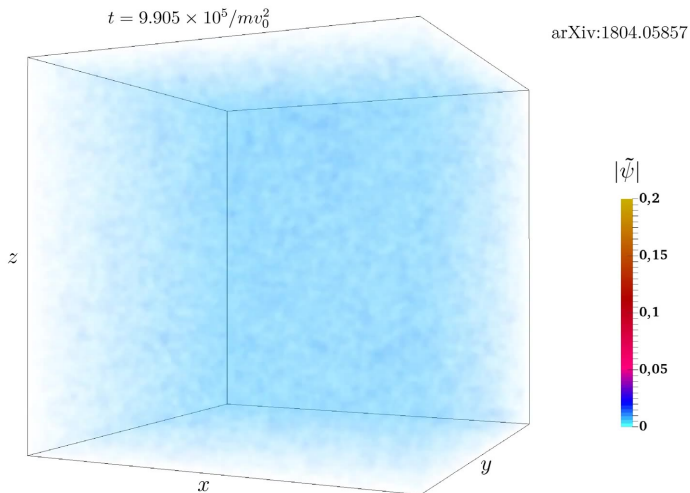
$$M \sim 4 \times 10^9 M_\odot, M_{\text{BEC}} \sim 10^8 M_\odot$$

Schive, Chiueh, Broadhurst '14



Simulation of Bose condensation

Start: virialized (random) gas in the box

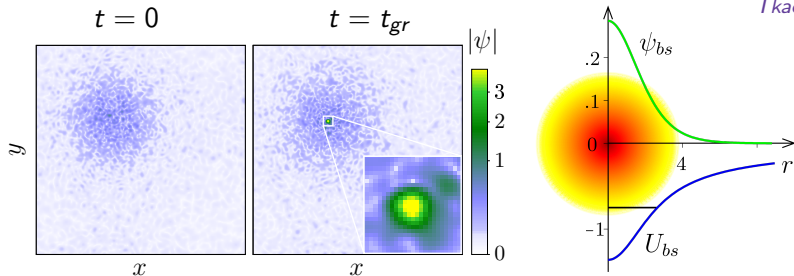


DL, Panin, Tkachev '18

Condensate + gravity = Bose (axion) star

Ruffini, Bonazzola '69

Tkachev '86



- All axions \in ground level of $U_{bs}(x)$

Growth of Bose stars

see A. Dmitriev's talk

→ **Simulations:** stop growing at $M_{bs} \propto \underbrace{\sqrt{M_{mc}/R_{mc}}}_{\text{"core-halo"}} - \text{small!}$

Schive et al '14

→ **Self-similar kinetics:** $M_{bs} \propto t^{1/3}$ — slow growth continues!

⇒ $O(1)$ of miniclusters is eaten by Bose stars

Dmitriev et al '23

Bose stars form 10 ÷ 50% of dark matter

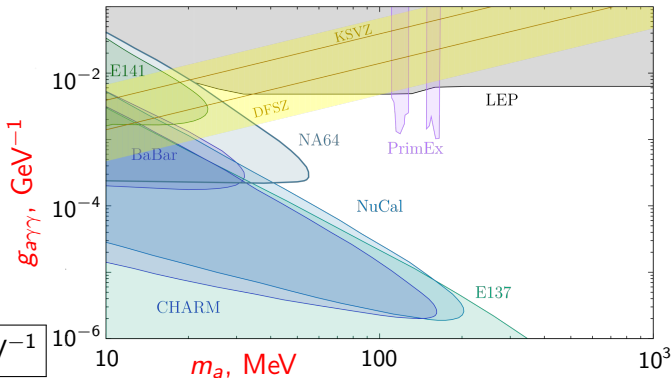
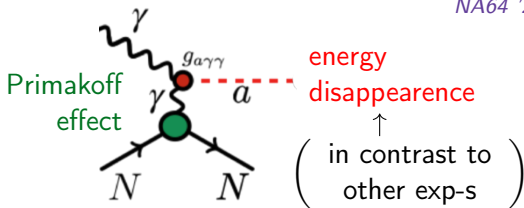
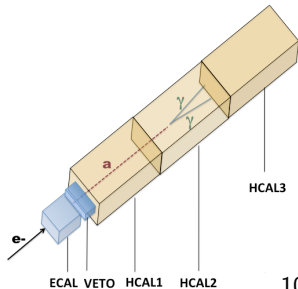
We need precise & realistic simulations!

Direct constraints: NA64 experiment

see I. Tliso, D. Kirpichnikov's talks

→ 16 institutes + CERN SPS

NA64 '20



at $m_a \ll 10$ MeV

$$g_{a\gamma\gamma} \lesssim 2 \cdot 10^{-4} \text{ GeV}^{-1}$$

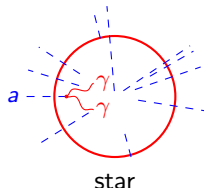
Some experiments: axions

cm. Irastorza, Redondo '18

• Star cooling

Axion emission from $\text{Vol}_{\text{star}} \Rightarrow$ faster evolution
 \Rightarrow too many red giants

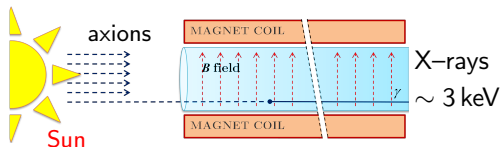
SN1987a: shorter ν signal



$$g_{a\gamma\gamma} \lesssim 10^{-10} \text{ GeV}^{-1}$$

• Helioscopes [CAST, IAXO] (planned)

Axions from the Sun

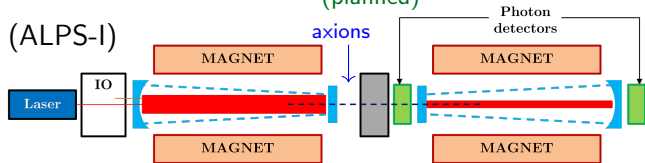


$$g_{a\gamma\gamma} \lesssim 10^{-10} \text{ GeV}^{-1} \text{ (CAST)}$$

• Light through the wall [ALPS, OSQAR, ALPS-II, III] (planned)

$$g_{a\gamma\gamma} \lesssim 10^{-7} \text{ GeV}^{-1} \text{ (ALPS-I)}$$

visible light



Some experiments: dark matter axions

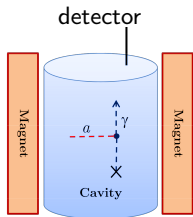
cm. Irastorza, Redondo '18

- $f_a \sim 10^{27} (m_a/\mu\text{eV})^{-4} \gg 1$ — coherent amplification!

● Haloscopes (ADMX)

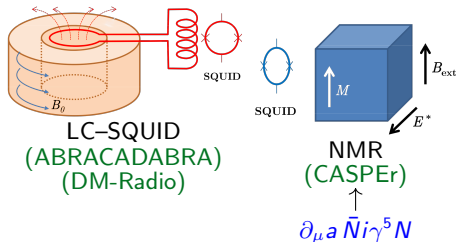
- + resonance: $m_a = \omega \sim L^{-1} \sim \mu\text{eV}$
 \Rightarrow high sensitivity!
- resonance: scan over $\omega = m_a$

$$g_{a\gamma\gamma} \lesssim 10^{-15} \text{ GeV}^{-1} \quad \text{at} \quad m_a \sim 2 - 3 \mu\text{eV} \quad (\text{ADMX})$$

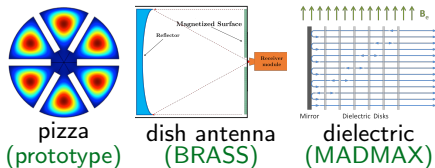


● Plethora of other ideas!

$m_a \ll \mu\text{eV}$ — large volume

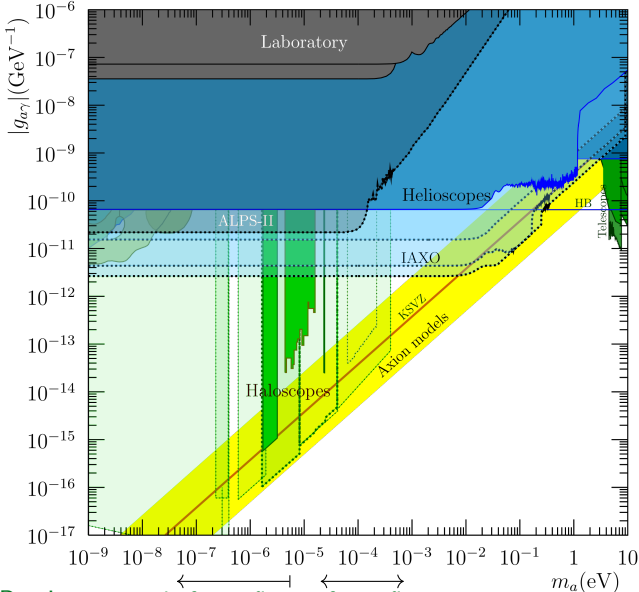


$m_a \gg \mu\text{eV}$ — small volume



Experimental constraints on axion-like dark matter

из Irastorza, Redondo '18



QCD axion:

before infl after infl

Manifestations of miniclusters

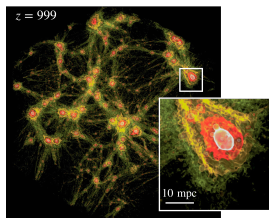
- **Suppression of the signal**

Eggemeier et al '22

miniclusters \Rightarrow mini-voids

\rightarrow mean signal $\times \underbrace{\sqrt{30\%}}_{0.5}$

\rightarrow in mini-voids: signal/3.5



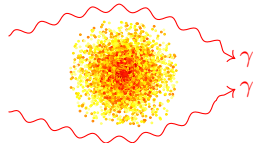
- **Microlensing?**

Kolb, Tkachev '94

\rightarrow only if $M_{\text{mc}} \gtrsim 10^{-11} M_{\odot}$, $m_a \gtrsim 10^{-3} \text{ eV}$

not for QCD axions? *Ellis et al '22*

\Rightarrow challenging for future expts.



- **Tidal streams**

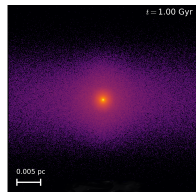
Tinyakov et al '15; Shen et al '22

\rightarrow meet the minicluster: once per 10^5 yr

\rightarrow But: minicluster + star = tidal stream

\rightarrow meet the stream: once per ~ 20 yr

\Rightarrow haloscopes may be lucky!



Manifestations of Bose stars (A. Panin's talk)

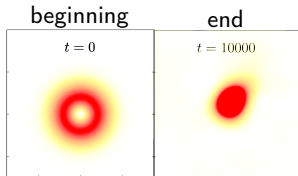
- **Rotating Bose stars decay**

shed off angular momentum

⇒ **bad black hole impostors**

Dmitriev et al '21

Sanchis-Gual '19



- **Heavy stars collapse**

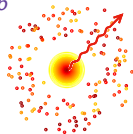
$-\lambda a^4$ attraction ⇒ **bosenova**

DL, Panin, Tkachev '16

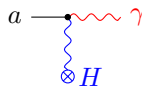
QCD axions: $M_{bs} \gtrsim 10^{-13} M_{\odot}$

+ magnetic field ⇒ **flashes!**

Arakawa '24



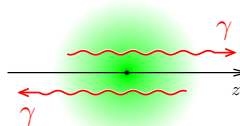
relativistic axions



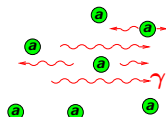
- **Axion laser = avalanche $a \rightarrow \gamma\gamma$**

→ $g_{a\gamma\gamma} \gtrsim 10^2 \cdot g_{a\gamma\gamma}^{\text{QCD}}$

DL, Panin, Tkachev '20



Tkachev '87



⇒ **Flashes in the sky!** — Fast Radio Bursts?

→ Can be triggered by a bosenova

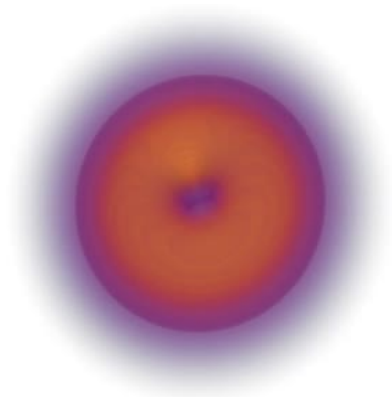
⇒ **Reionization of the Universe**

Escudero et al '23

⇒ **Contributes into diffuse radiobackground**

Eby, Takhistov '24

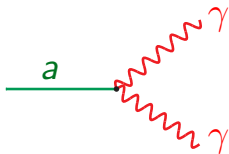
Instability of a rotating Bose star



arXiv: 2104.00962

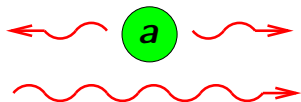
Interaction of QCD axions with photons

$$\mathcal{L}_{\text{int}} \propto g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$



spontaneous decay

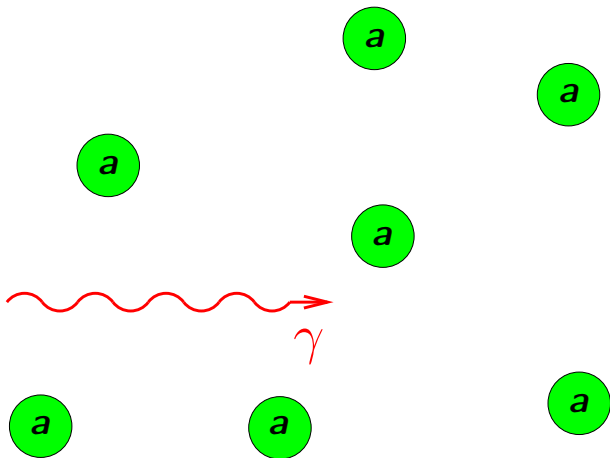
\mathcal{P}



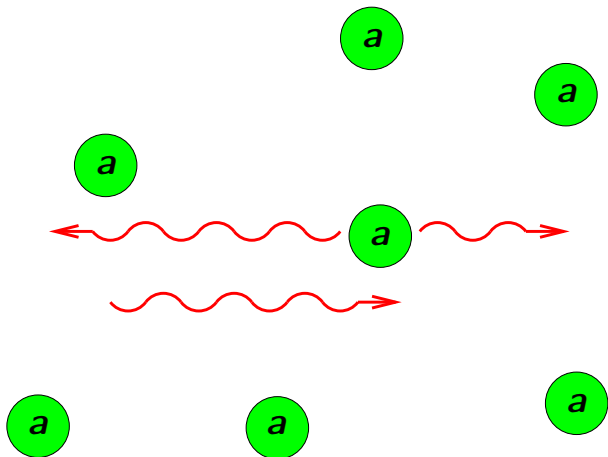
stimulated decay

$\mathcal{P} \times N_\gamma$

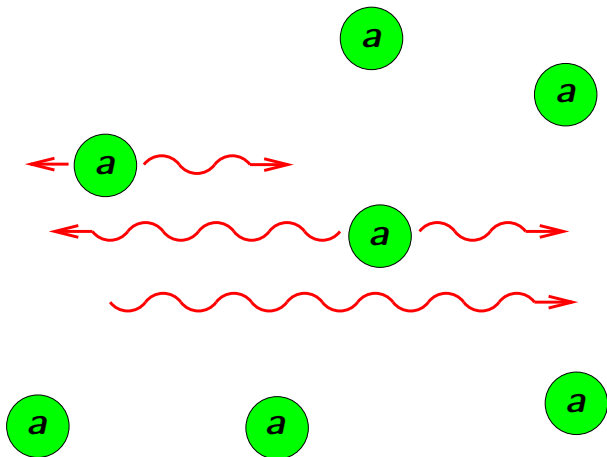
Axion laser: avalanche $a \rightarrow \gamma\gamma$



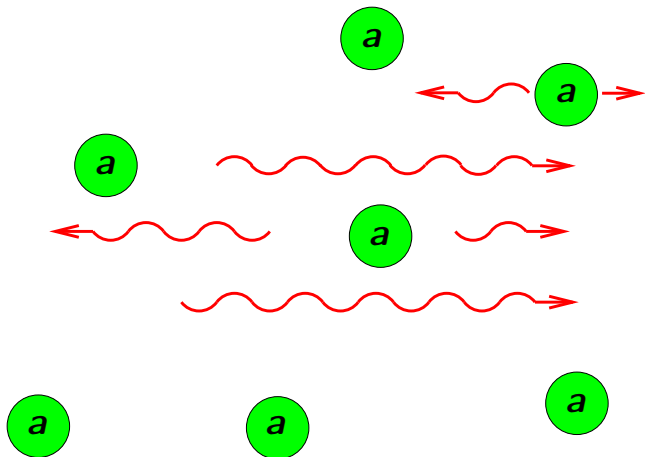
Axion laser: avalanche $a \rightarrow \gamma\gamma$



Axion laser: avalanche $a \rightarrow \gamma\gamma$



Axion laser: avalanche $a \rightarrow \gamma\gamma$



Condition: at least one decay per γ pass

Tkachev '87; Riotto, Tkachev '00; Hertzberg, Schiappacasse '18

Conclusions

Ultralight (axion-like) dark matter is:

- Motivated theoretically
- Leads to unusual cosmology with miniclusters and Bose stars
- Search for axions is a creative business!

Heavy dark matter
is classics



While the light one
is a heavy metal!



Thank you for attention!

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