

Astrophysical and cosmological Appearances of Intermediate- Mass Black Holes



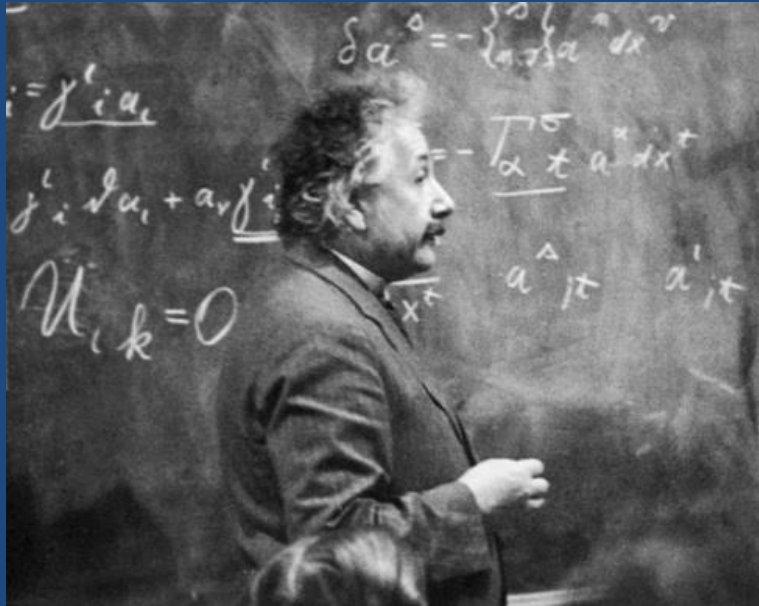
Konstantin Postnov
Sternberg Astronomical Institute, Moscow
University

Supported by RSF 23-42-00055

Plan

- Introduction
- BH Zoo
- Astrophysical IMBHs
- Cosmological IMBHs
 - Binary IMBHs from EPTA
 - Primordial binary IMBHs
- Prospects for binary IMBH detection

Black holes



A. Einstein, 1912-1915

Gravitation = curvature of spacetime

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}, \quad ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$

$$R_{\alpha\beta} \equiv R^\mu{}_{\alpha\mu\beta}$$

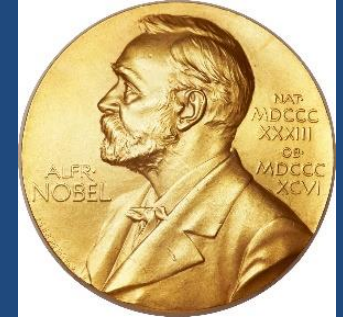
$$R \equiv R^\alpha{}_\alpha$$

From Schwarzschild to Penrose

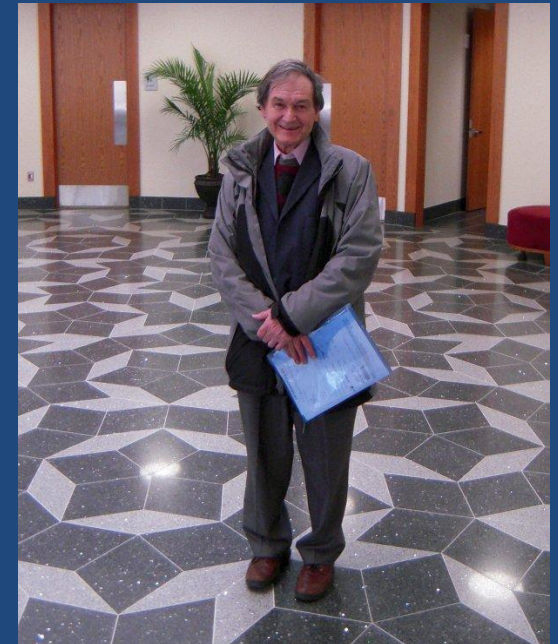


1916, K. Schwarzschild, metrics of spherically symmetric spacetime outside a massive body

$$ds^2 = \left(1 - \frac{r_s}{r}\right) c^2 dt^2 - \frac{dr^2}{\left(1 - \frac{r_s}{r}\right)} - r^2 (\sin^2 \theta d\varphi^2 + d\theta^2)$$



2020

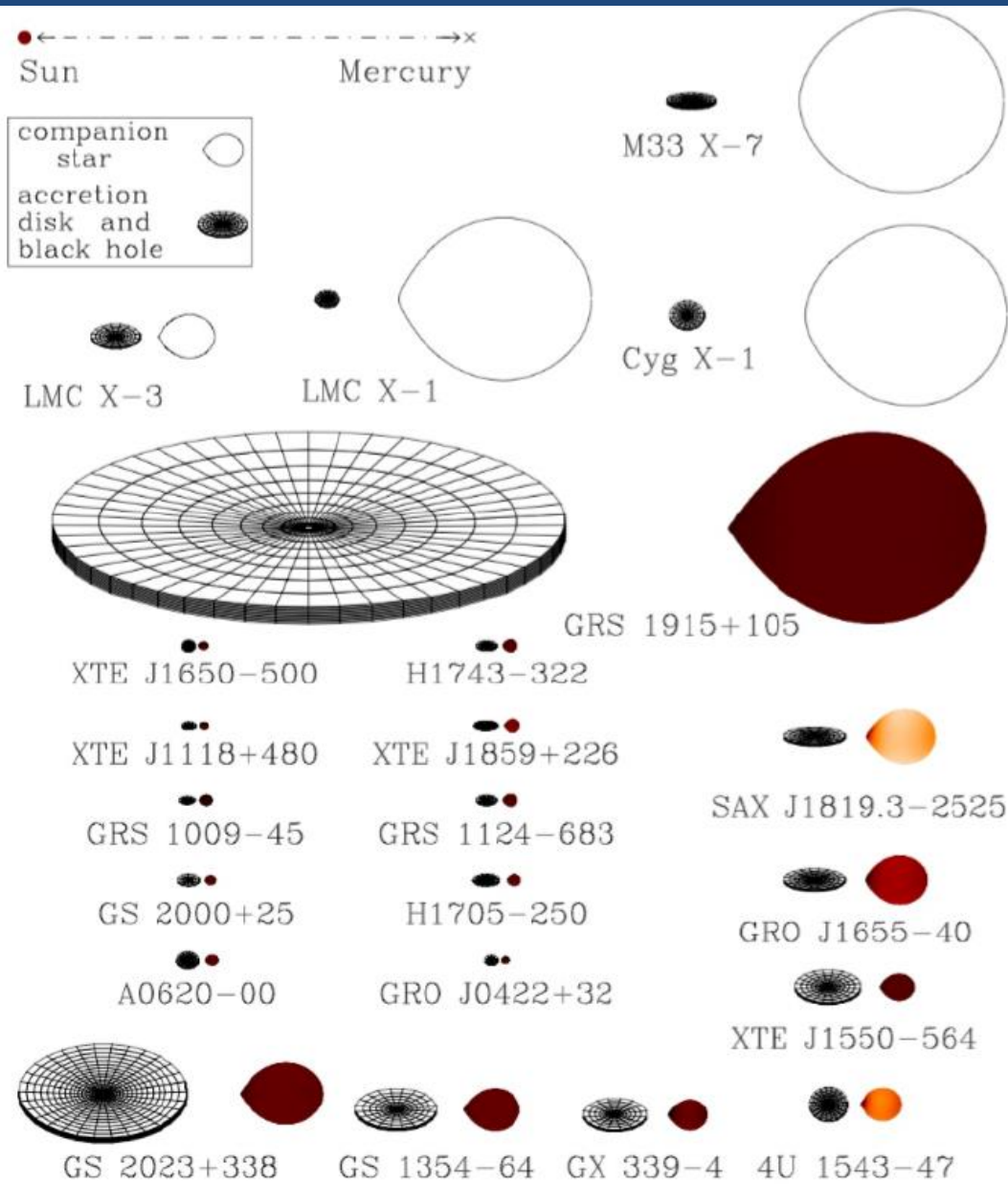


Evidence for stellar-mass BH

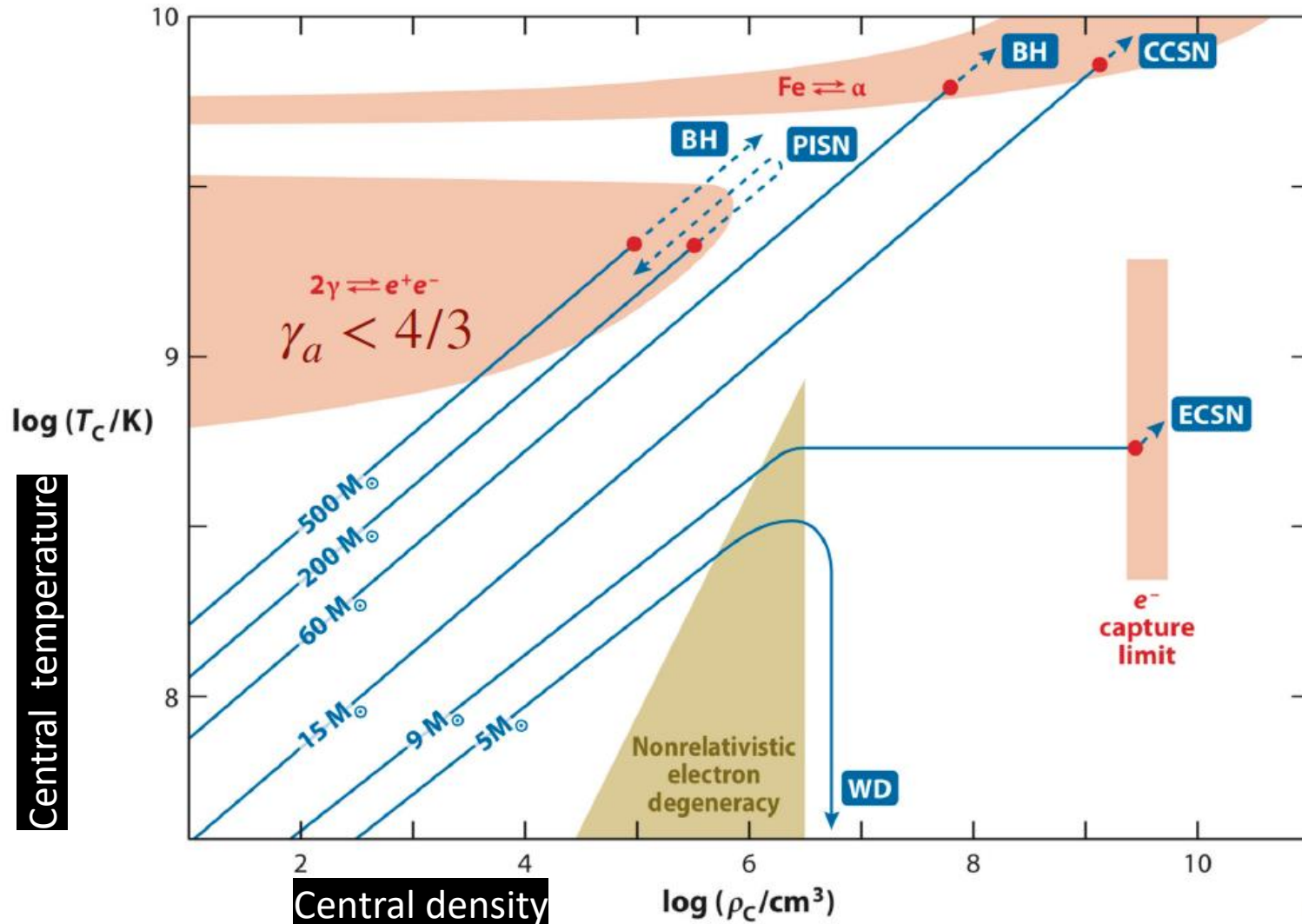
- 1939, Oppenheimer & Volkoff, maximum mass of neutron star (GR: $<3 M_{\odot}$)
- 1939, Oppenheimer & Snyder, collapse of a pressureless star
- 1970s, Search for BH in X-ray binaries
- Presently, ~ 100 BH binaries are known in the Galaxy and in nearby galaxies

BH candidates in binary systems

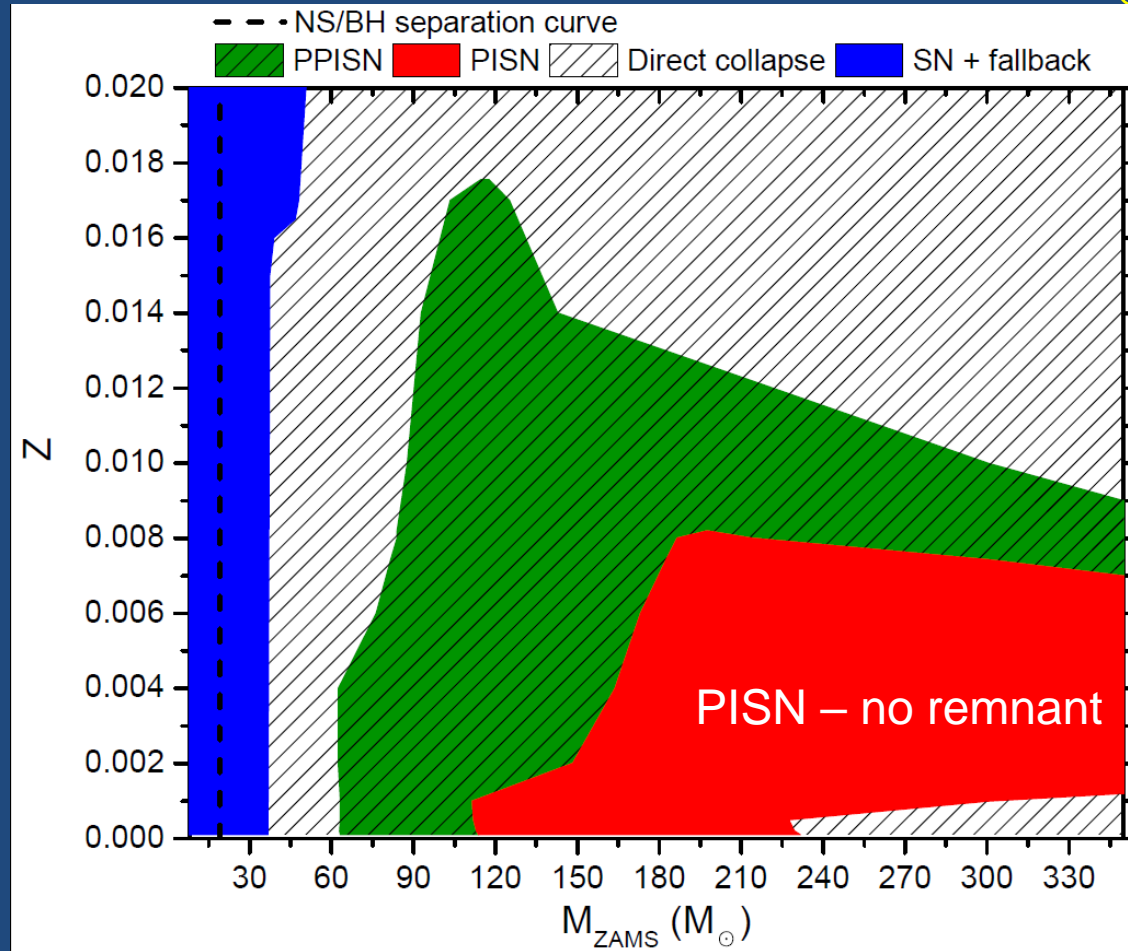
- $M_x > 3M_\odot$
- No signatures of surface
 - *No periodic pulsations (X-ray pulsar)*
 - *No explosions on the surface (X-ray type I bursts)*
 - *No magnetic fields (radiopulsar)*



BH formation from stars (solar metallicity)



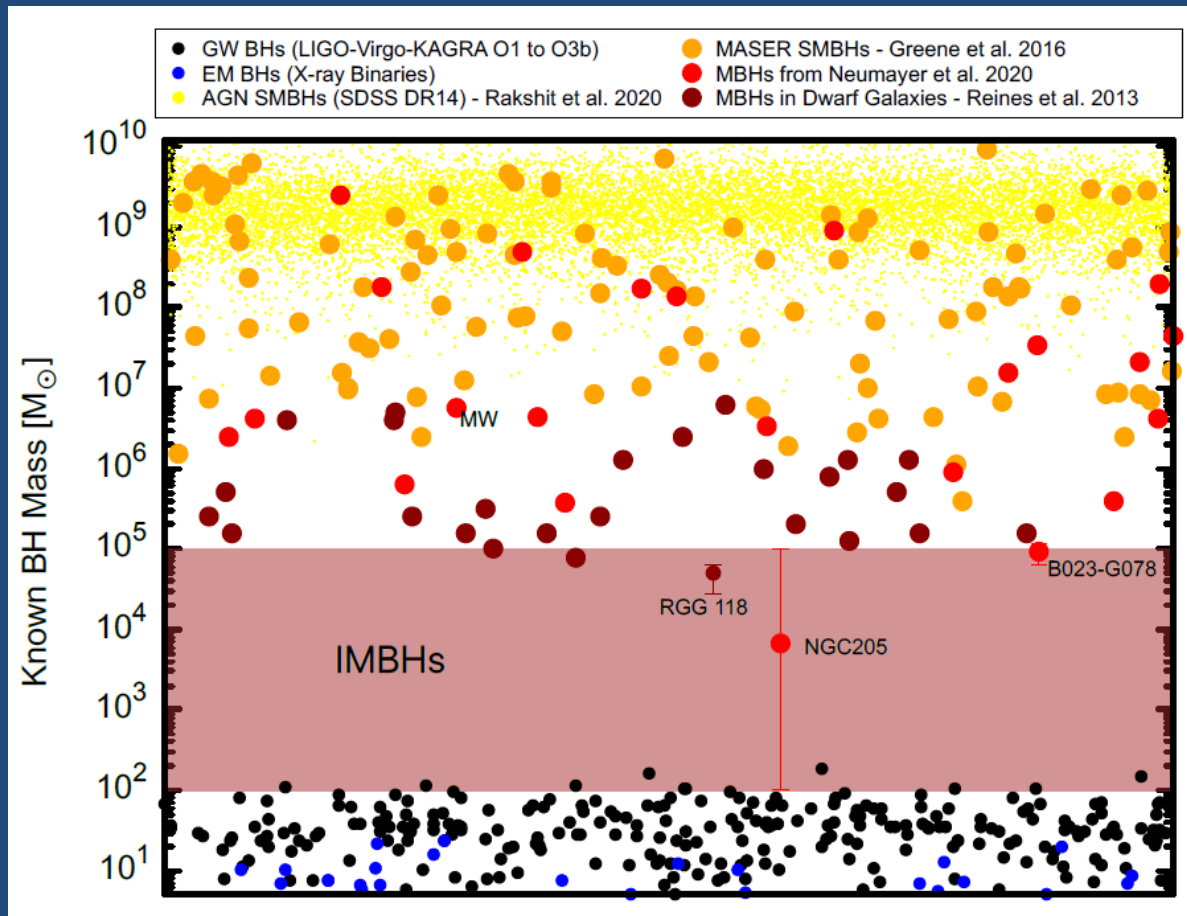
e⁺/e⁻ pair creation in massive stars → stellar BH 'upper mass gap' (60-130 M_⊙)



Spera, Mapelli
2017

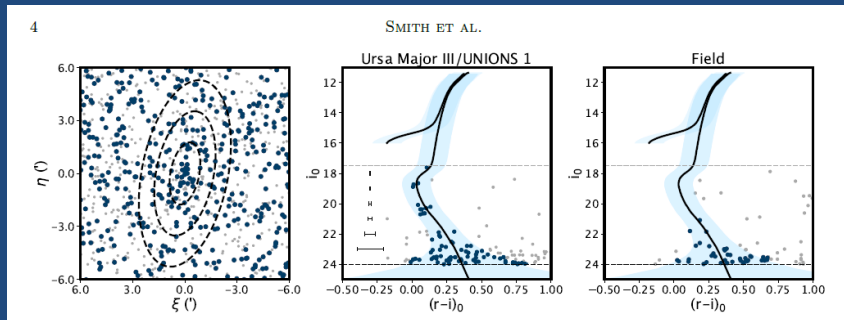
Pair-instability SNe (PISNe, Fowler & Hoyle 1964),
pulsational PISNe (PPISNe) (Woosley 2017)

BH: Astronomical Observations

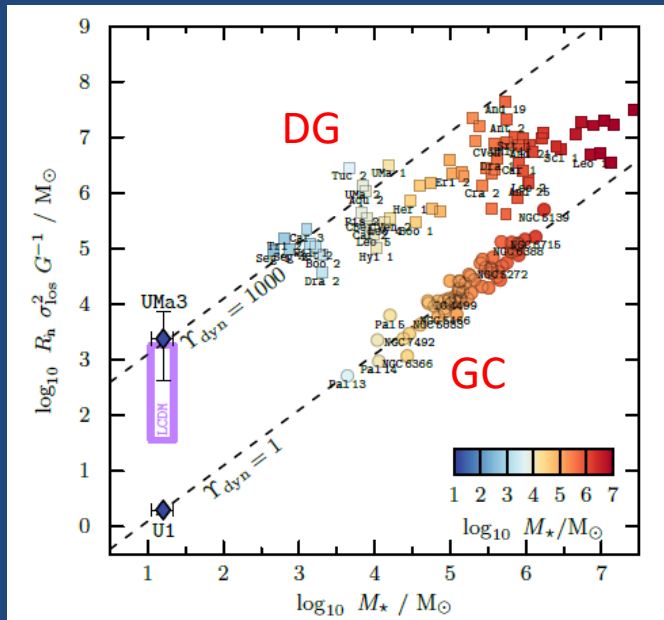


IMBH in dwarf galaxies?

Smith et al 2311.10147



Errani et al 2311.10134

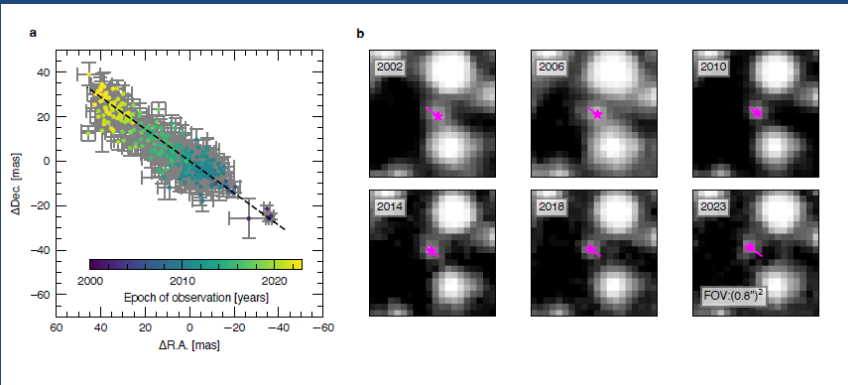
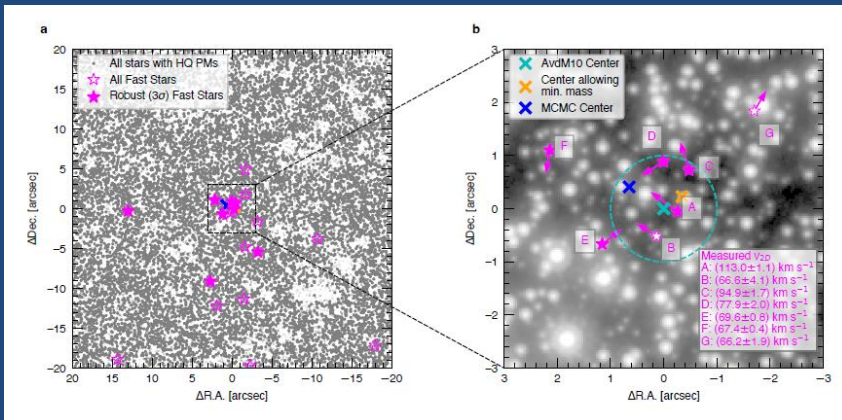


- Uma II/UNIONS 1 – the lightest MW satellite galaxy
- $M^* = 16 \pm 5 M_{\odot}$
- $r_h = 3 \pm 1$ pc
- Dynamical evidence – DM-dominated or an IMBH?

$$\Upsilon_{\text{dyn}} \equiv \frac{M_{\text{dyn}}(< r_h)}{M_{\star}(< r_h)} \approx 8 R_h \langle \sigma_{\text{los}}^2 \rangle G^{-1} M_{\star}^{-1}$$

IMBH in GC ω Cen (stripped nucleus of accreted dwarf galaxy)

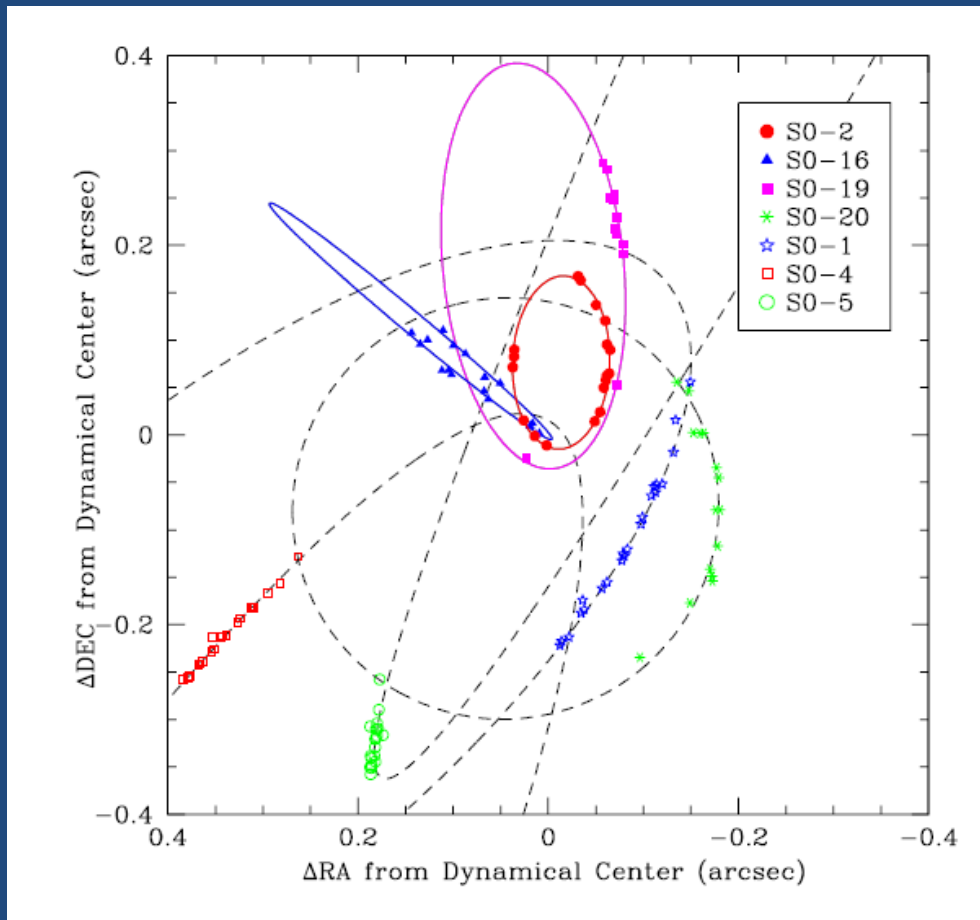
Haeberle+ 2024, 2405.06015



- Fast-moving stars with $v > 3\sigma$ above $v_{esc} = 62$ km/s in the center of ω Cen suggest IMBH with $M > 8200 M_{\odot}$
- Similar IMBH in stripped nuclei of satellites around M31 (G1: $\sim 20000 M_{\odot}$, B023-G078 $\sim 10^5 M_{\odot}$)

Supermassive BH (SMBH)

Orbits of stars around Sgr A*



Chez et al. 2005

$$\theta \sim \frac{r_g}{d} \sim 0.05 \text{ milliarcseconds}$$

$$M_{\text{SgrA}^*} = 4.10 \pm 0.03 \times 10^6 M_{\odot}$$

(GRAVITY collab. 2018)

Deviation of S0-2/S2 orbit
($P=11.5$ yr) from Keplerian motion
(Do et al 2019)

Einstein Horizon Telescope

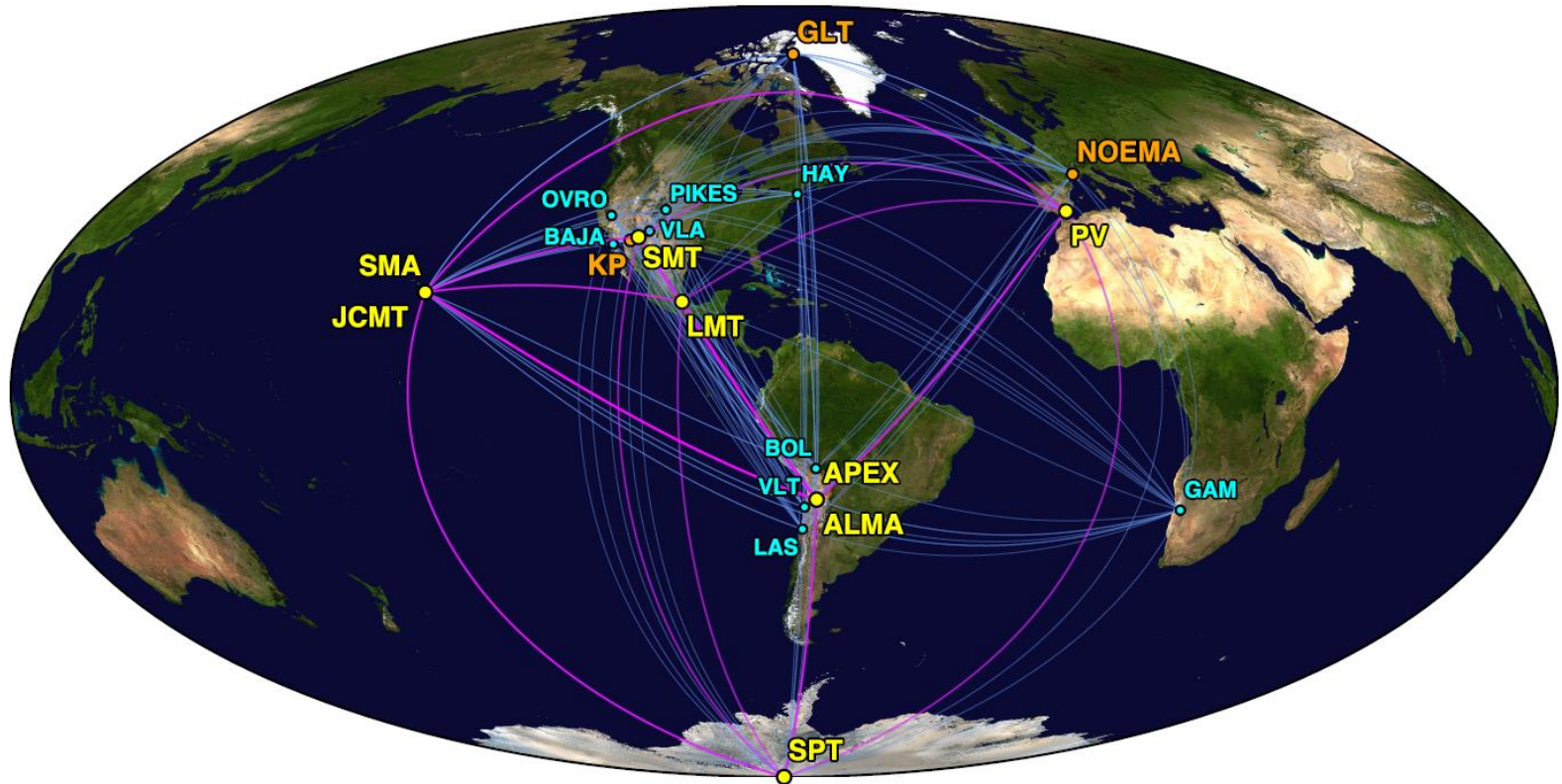
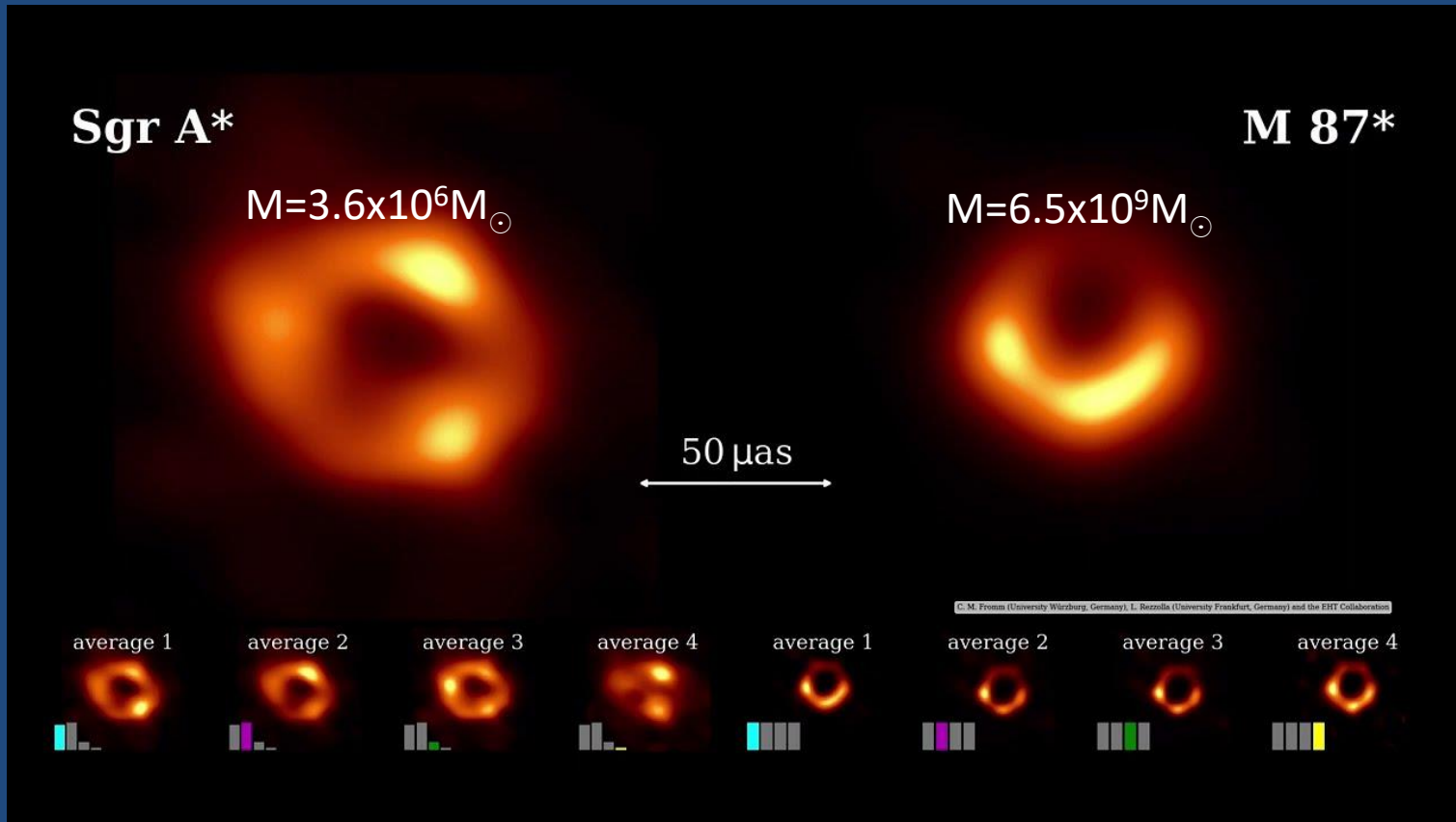


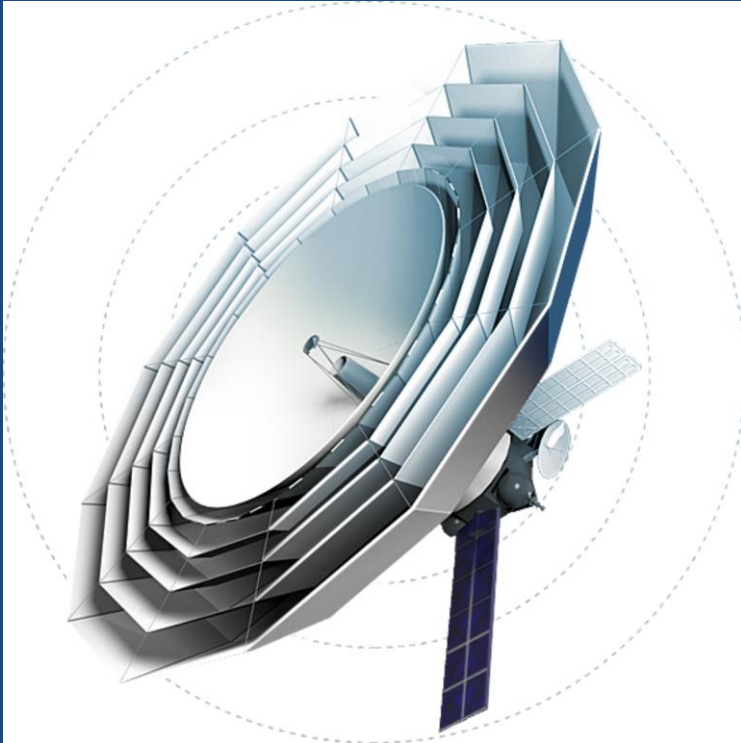
Figure 5: Distribution of stations around the globe. Stations that participated in the EHT2017 observing campaign are labeled in yellow, while the additional stations that will be present in the EHT2020 array are labeled in orange. Several possible new site locations for the ngEHT are labeled in cyan. Current EHT2017 baselines are shown in magenta.

EHT image of local SMBH

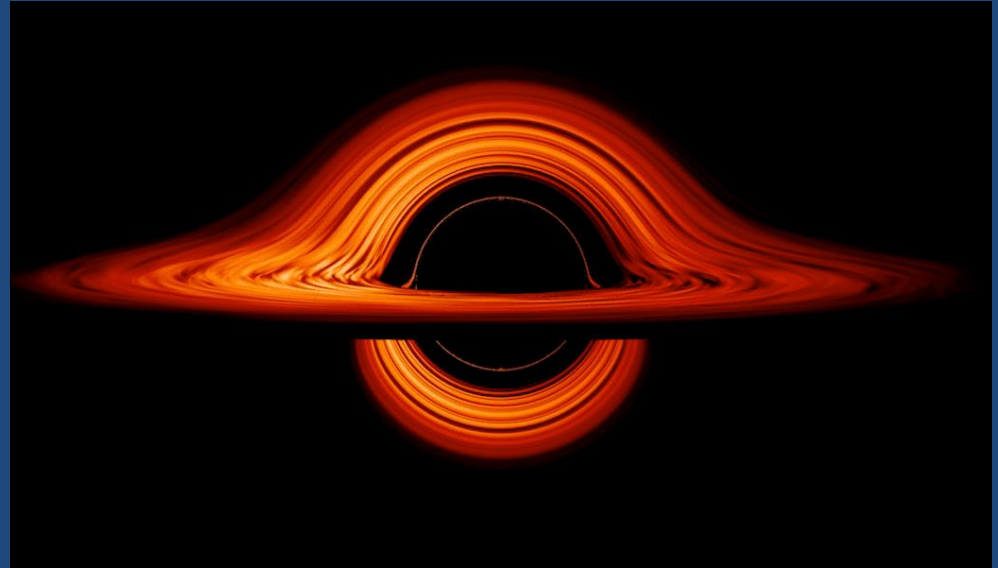


$\lambda=1.3 \text{ mm}, \theta \sim 20 \text{ mas}$

Future prospects: Millimetron space mission

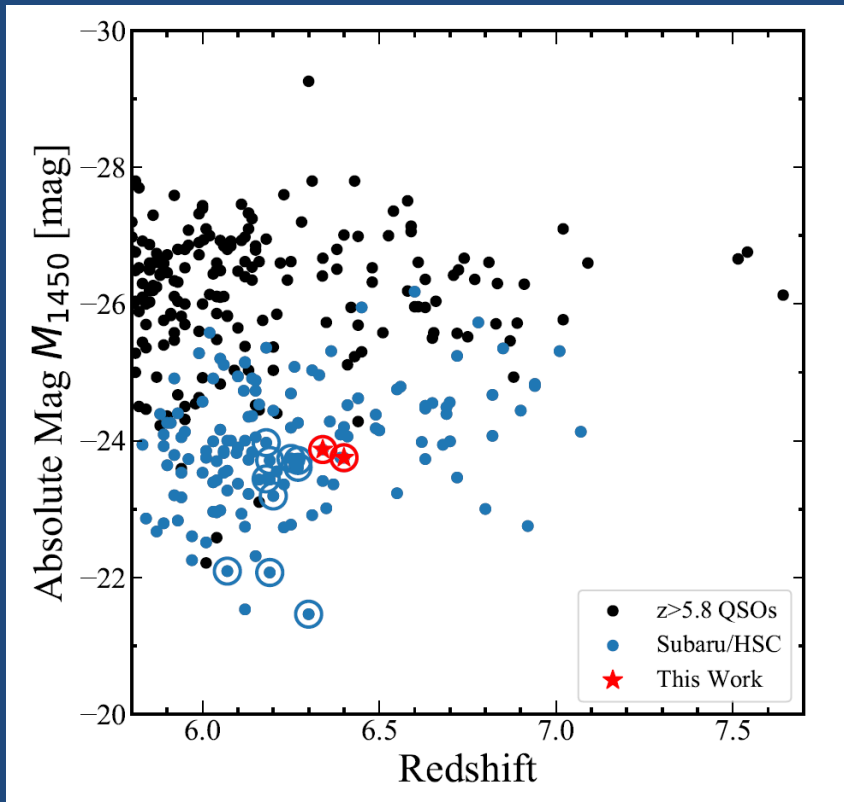


$\lambda=0.07-10$ mm
VLBI L2-Earth
 $\theta\sim 0.4-0.1$ mas
~2030



Credits: NASA's GSFC/Jeremy Schnittman

High-redshift QSOs and SMBH growth

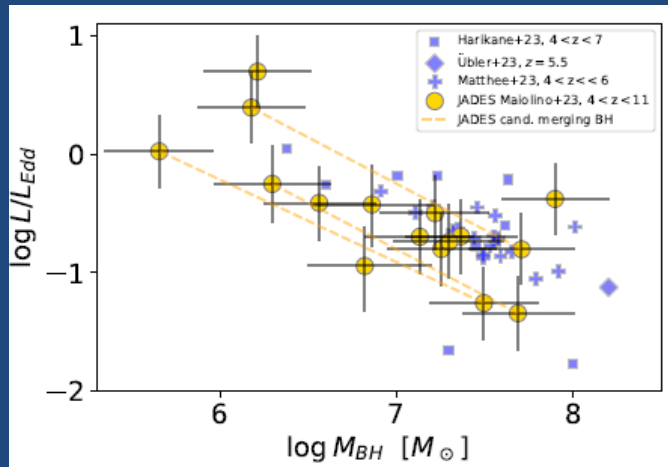


- >200 QSO $z > 6$
- Central SMBH $M \sim 10^9 M_{\odot}$
- Inefficient merging
- Eddington-limited accretion
- **Need for 'massive 'seeds'**
 $M \sim 10^3 - 10^4 M_{\odot}$
- Population III star?
Primordial massive BHs?

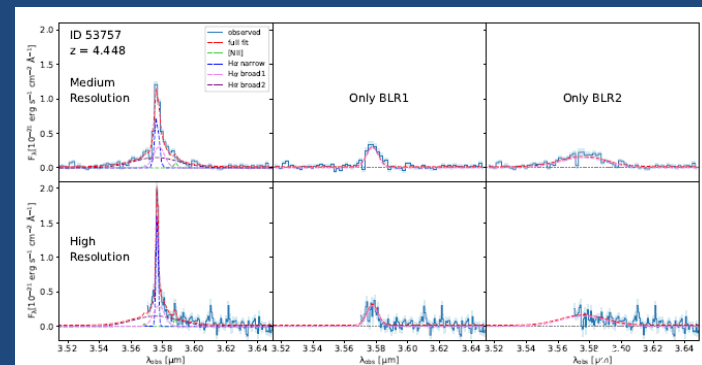
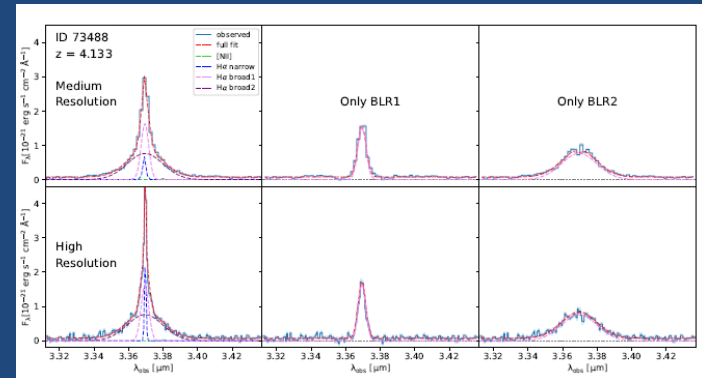
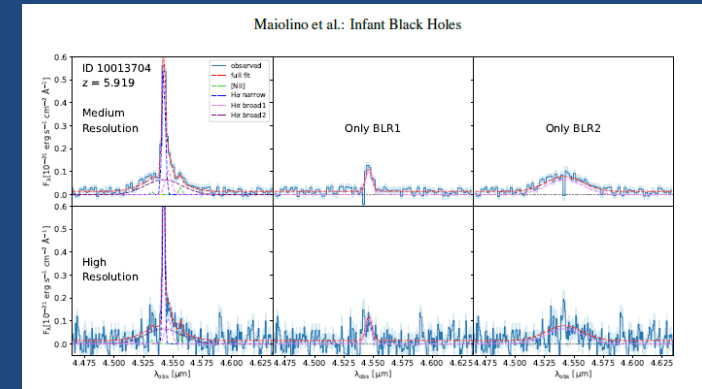
2211.14329

Double SMBH in JADES JWST

- Dual BLR H_α lines in three out of 11 AGNs $4 < z < 6 \rightarrow$ merging accreting double BH



Maiolino et al. 2308.01230



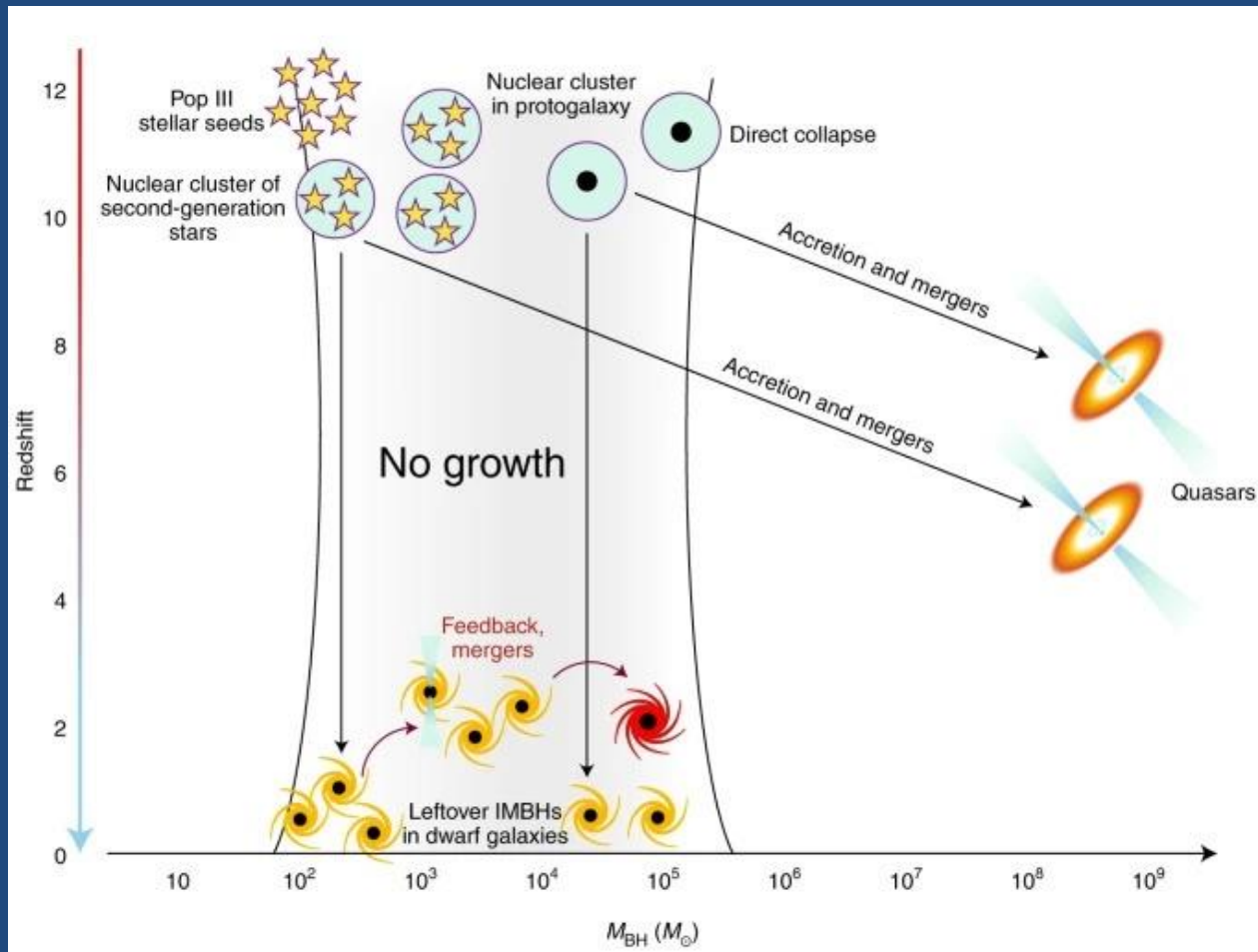
Merging binary SMBH candidates

Table 3: Parameters inferred for the broad line AGN presented in this JADES sub-sample, and including GN-z11 from [Maiolino et al. \(2023a\)](#).

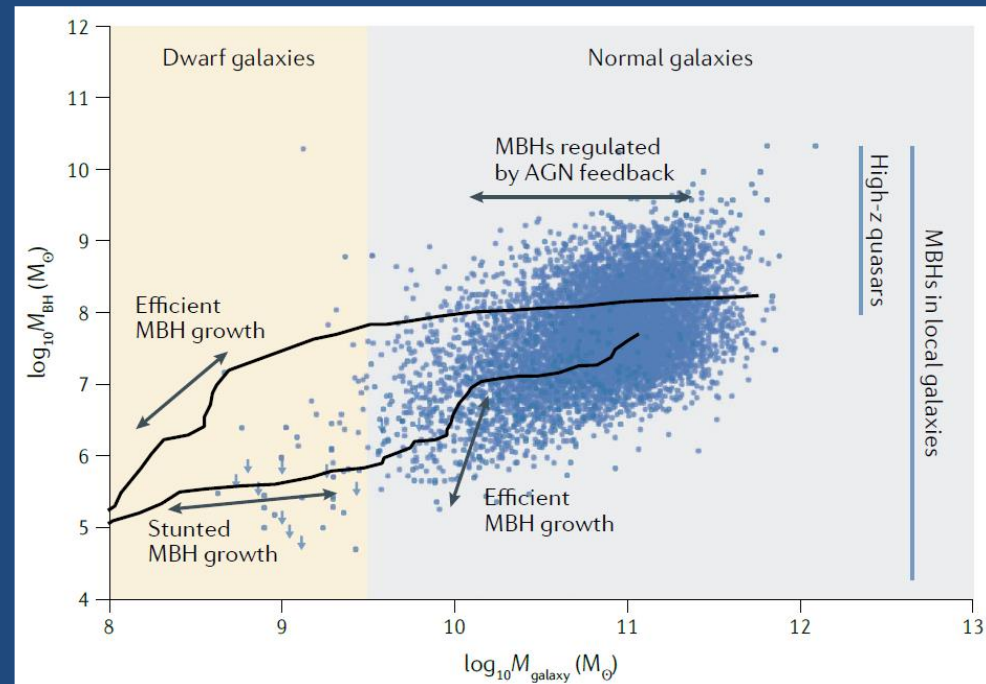
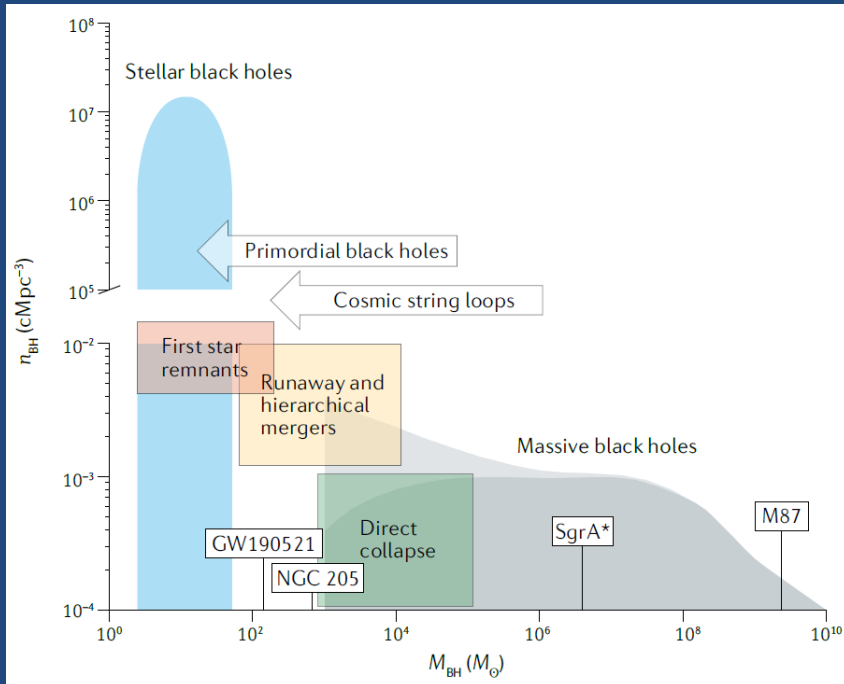
ID	Comp.	$\lg(M_{\text{BH}})$ [M_{\odot}]	$\lg(L_{\text{bol}})$ [erg s^{-1}]	L/L_{Edd}	$\lg(M_{\text{star}})$ [M_{\odot}]	$\lg(\sigma)^a$ [km s^{-1}]	R_* [kpc]	Sersic index	$\lg(M_{\text{dyn}})$ [M_{\odot}]	M_{UV} [mag]	A_V^b [mag]
10013704	BLR1	$5.65^{+0.31}_{-0.31}$	43.8	1.06	$8.88^{+0.66}_{-0.66}$	$1.93^{+0.05}_{-0.06}$	0.15	0.8	$9.23^{+0.1}_{-0.13}$	-18.89	0.27
	BLR2	$7.5^{+0.31}_{-0.31}$	44.3	0.06							
8083		$7.25^{+0.31}_{-0.31}$	44.6	0.16	$8.45^{+0.03}_{-0.03}$	$1.9^{+0.06}_{-0.07}$	0.11	5.7	$8.84^{+0.11}_{-0.15}$	-18.67	0.64
1093		$7.36^{+0.32}_{-0.31}$	44.8	0.2	$8.34^{+0.2}_{-0.2}$	$1.95^{+0.05}_{-0.06}$	-	-	-	-17.48	0.99
3608		$6.82^{+0.38}_{-0.33}$	44.0	0.11	$8.38^{+0.11}_{-0.15}$	$1.92^{+0.06}_{-0.07}$	-	-	-	-19.5	0.48
11836		$7.13^{+0.31}_{-0.31}$	44.5	0.2	$7.79^{+0.3}_{-0.3}$	$1.96^{+0.05}_{-0.06}$	0.48	0.8	$9.81^{+0.1}_{-0.14}$	-18.75	0.68
20621		$7.3^{+0.31}_{-0.31}$	44.7	0.18	$8.06^{+0.7}_{-0.7}$	$1.93^{+0.06}_{-0.07}$	-	-	-	-18.27	0.67
73488	BLR1	$6.18^{+0.3}_{-0.3}$	44.7	2.48	$9.78^{+0.2}_{-0.2}$	$1.64^{+0.11}_{-0.15}$	0.59	0.8	$9.28^{+0.21}_{-0.41}$	-18.73	0.45
	BLR2	$7.71^{+0.3}_{-0.3}$	45.0	0.16							
77652		$6.86^{+0.35}_{-0.34}$	44.5	0.38	$7.87^{+0.16}_{-0.28}$	$1.95^{+0.06}_{-0.07}$	-	-	-	-18.28	0.39
61888		$7.22^{+0.31}_{-0.31}$	44.8	0.32	$8.11^{+0.92}_{-0.92}$	$1.85^{+0.07}_{-0.09}$	0.09	0.9	$8.92^{+0.22}_{-0.46}$	-19.0	0.69
62309		$6.56^{+0.32}_{-0.31}$	44.2	0.39	$8.12^{+0.12}_{-0.13}$	$1.87^{+0.07}_{-0.08}$	0.21	1.3	$9.27^{+0.14}_{-0.2}$	-18.67	0.74
53757	BLR1	$6.29^{+0.33}_{-0.32}$	44.1	0.56	$10.18^{+0.13}_{-0.12}$	$1.77^{+0.09}_{-0.11}$	-	-	-	-18.9	0.36
	BLR2	$7.69^{+0.32}_{-0.31}$	44.4	0.05							
954 ^c		$7.9^{+0.3}_{-0.3}$	45.6	0.42	$10.66^{+0.09}_{-0.1}$	$1.91^{+0.06}_{-0.06}$	0.35	0.8	$9.56^{+0.13}_{-0.18}$	-19.78	0.64
GN-z11 ^d		$6.2^{+0.3}_{-0.3}$	45.0	5.5	$8.9^{+0.2}_{-0.3}$	-	0.20	0.9	-	-21.79	0.0

Maiolino et al. 2308.01230

Growth of SMBHs from astrophysical seeds



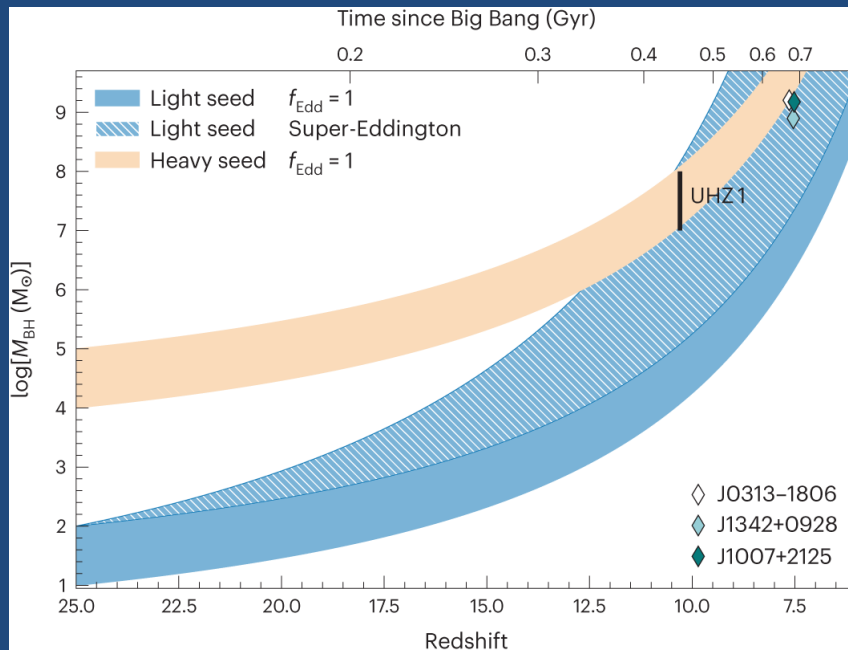
Massive BH formation



Whether and how the bridge between stellar and supermassive black holes was established when the first galaxies were forming is currently unknown (Volonteri et al. 2021).

SMBH growth from massive seeds

- A. Bogdan et al. 2023 NatAstronomy. UHZ1 $z=10.3$
- Growth from light seeds require super-Eddington accretion
- Growth from massive seeds is more natural



[2403.14745v1](#)

SUPERMASSIVE BLACK HOLE

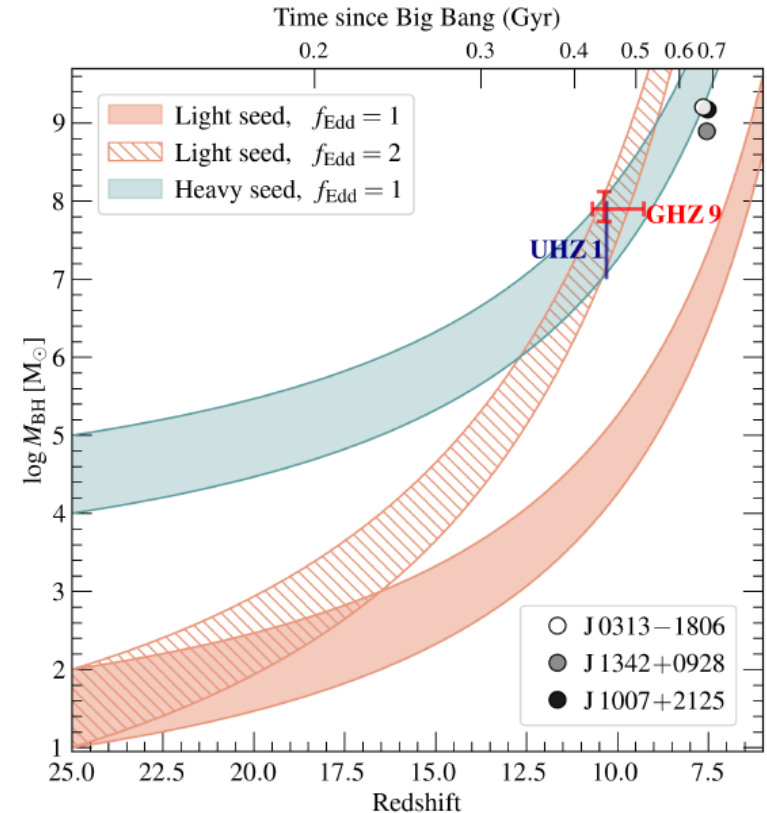
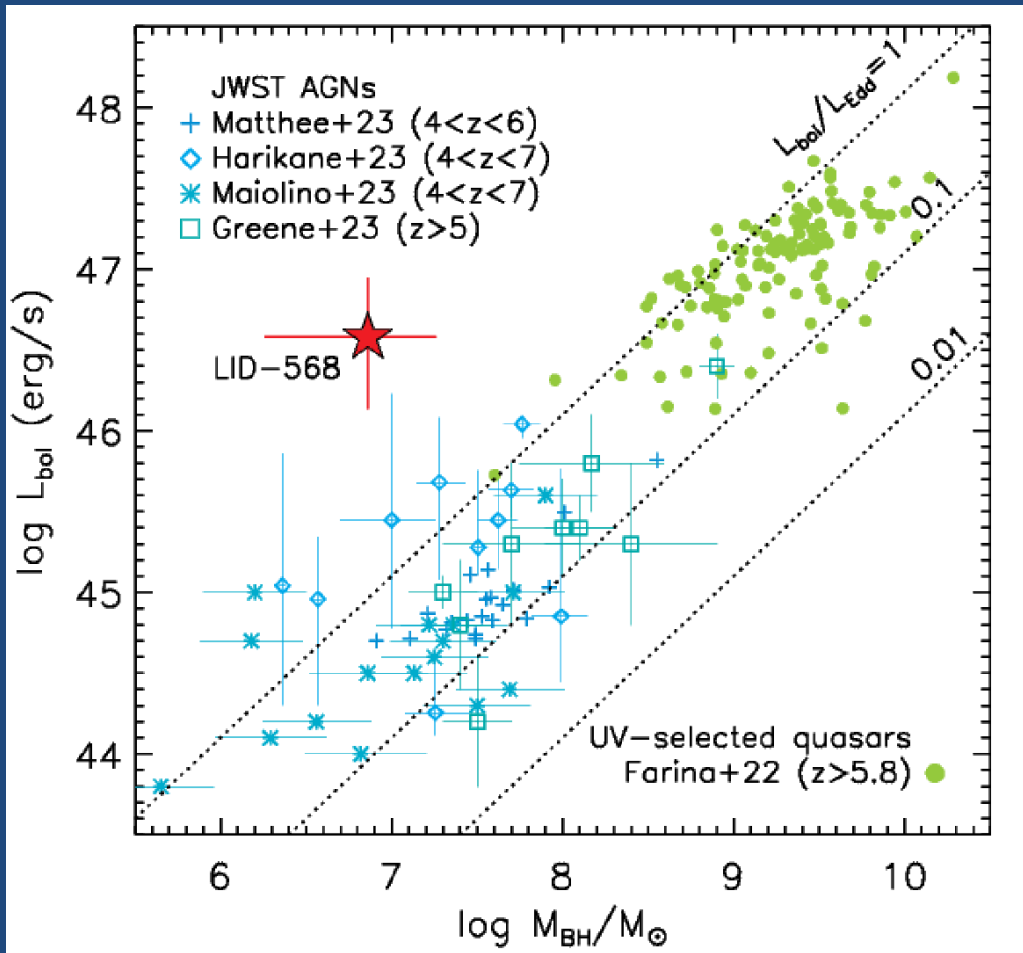


FIG. 3.— BH growth tracks for different initial seed masses and accretion rates assuming continuous accretion with a radiative efficiency of 10%. Light BH seeds only reach $10^4 - 10^5 M_{\odot}$ mass by $z \approx 10.4$ accreting at their Eddington limit. The estimated $\sim 8.0 \times 10^7 M_{\odot}$ mass for its redshift of $z \approx 10.37$, however, places the candidate BH of GHZ 9 as originating from a heavy seed, similar to UHZ 1 (Bogdán et al. 2024). While sustained accretion above the Eddington limit is unlikely (Willott et al. 2010; Smith et al. 2018), we also show the growth curves assuming $f_{\text{Edd}} = 2$ for light seeds with the hatched region, which would also be able to produce BH masses estimated for GHZ 9 and UHZ 1. We show three

Superluminal accretion onto SMBH?



$$L_{\text{Edd}} = 10^{38} [\text{erg s}^{-1}] \left(\frac{M}{M_{\odot}} \right)$$

Astrophysical BH formation: summary

- **Stellar-mass BHs**: from stellar evolution (core collapse of massive stars). $\sim 4\text{-}50 M_{\odot}$ (But: in dense stellar clusters, masses can be made higher due to previous coalescences)
- **Supermassive BHs** (in galactic centers, appear as AGNs and QSOs) $\sim 10^5 M_{\odot} - 10^{10} M_{\odot}$. Seeds are highly debated, growing in galaxy mergings

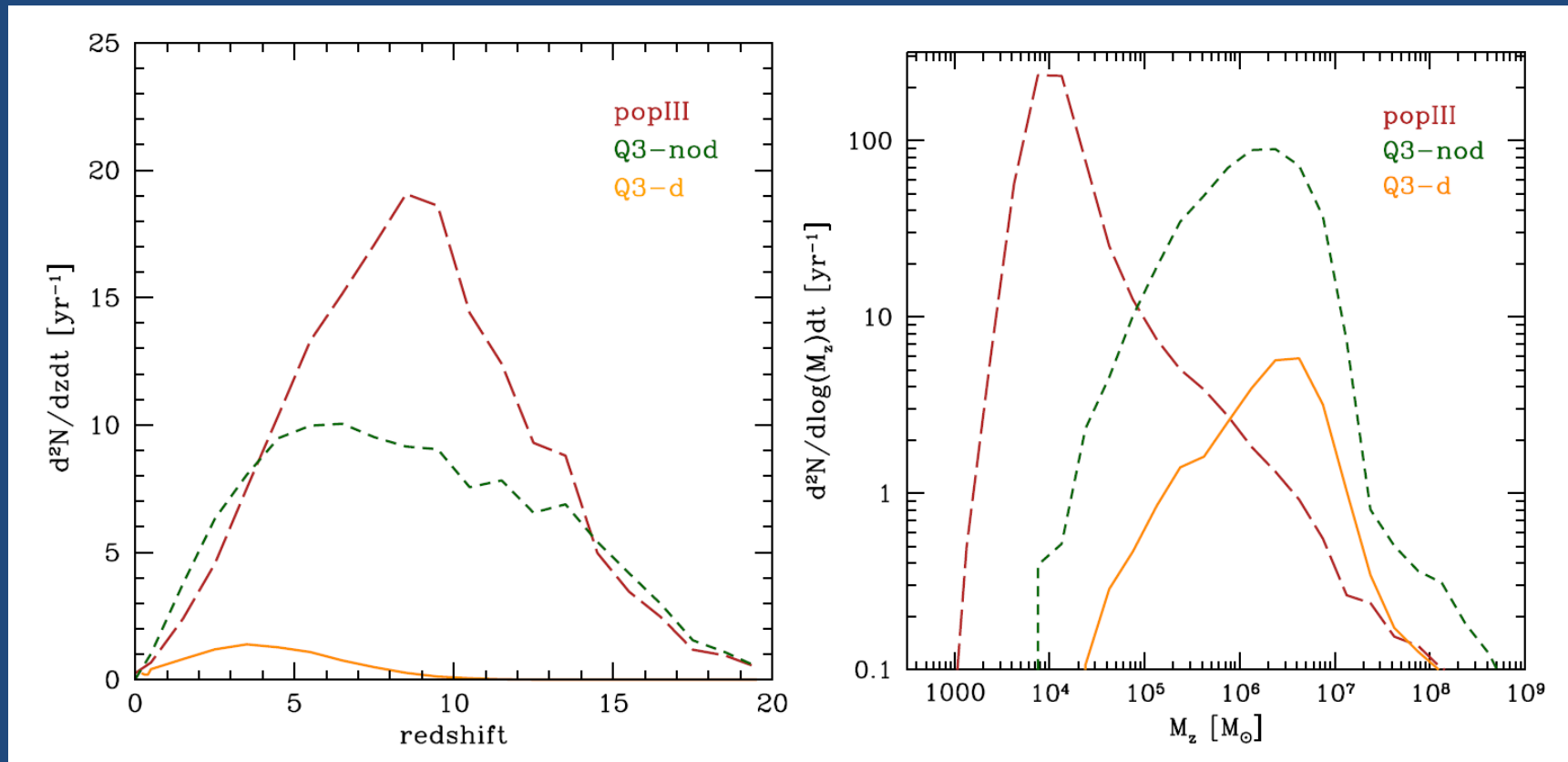
Why IMBH are interesting?

- Can seed the formation of SMB in early galaxies
- Coalescing binary IMBH can be primary sources for future space GW observatories (eLISA, TianQin, Taiji etc.)

How much binary IMBH can we expect?

- Astrophysical formation channel: galaxy mergings at $z < 20$
- Cosmological formation channel:
 - Primordial binary IMBH
 - Binary IMBH from mergings of DM halos with central primordial BH

Astrophysics: model binary MBH populations



Klein+ 2016

Constraints from recent PTA results

- NANOGrav
- EPTA
- InPTA
- CPTA

Stochastic GW signal from collection of individual sources

$$\Omega(f) = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f)$$

Energy density per logarithmic frequency

PSD in timing residuals

$$S(f) = \frac{h_c^2(f)}{12\pi^2 f^3}$$

RMS timing residuals in frequency bin i/T

$$\text{RMS}_i = \left(\int_{\Delta f_i} S(f) df \right)^{1/2} \approx (S(f_i) \Delta f_i)^{1/2} = \left(\frac{S(f_i)}{T} \right)^{1/2}$$

$$h_c(f) = A \left(\frac{f}{f_0} \right)^\alpha$$



$$S(f) = \frac{h_c^2(f)}{12\pi^2 f^3} = \frac{A^2}{12\pi^2 f_0^{2\alpha}} f^{-\gamma}$$

$$\gamma = 3 - 2\alpha$$

SMBH mergings: stochastic GW signal

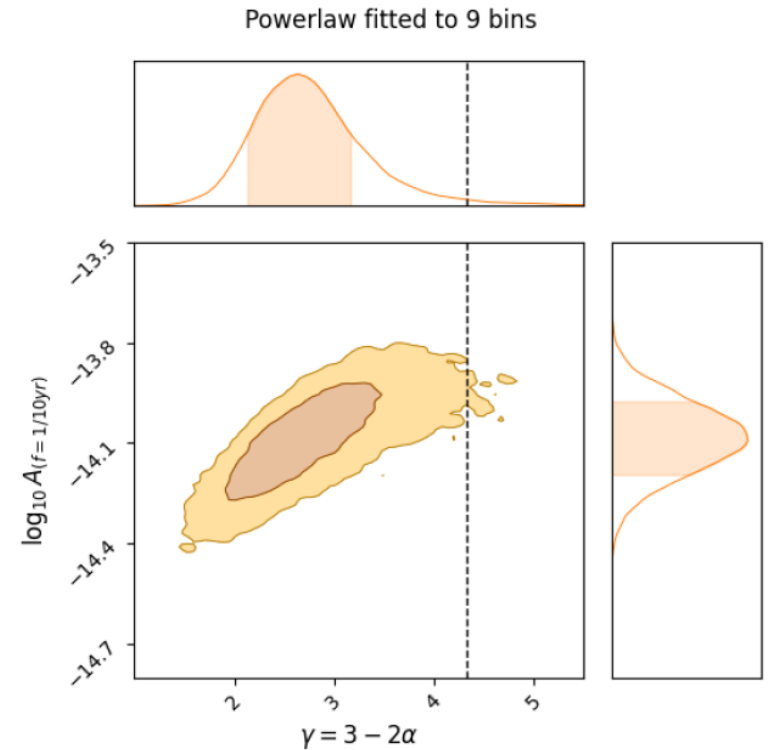
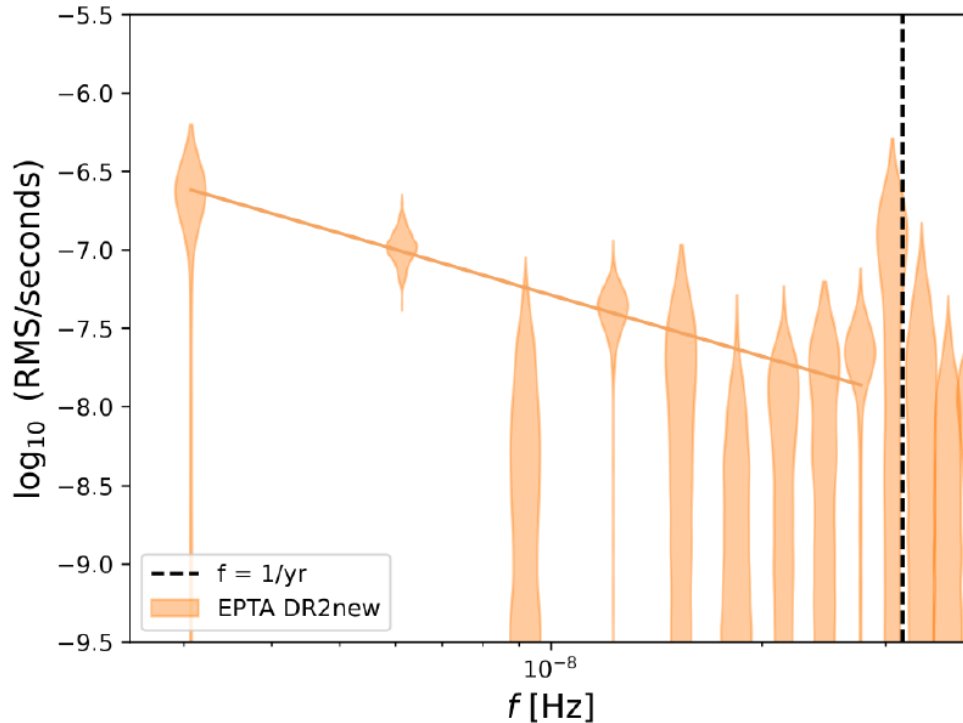
$$h_c^2(f) = \int_0^\infty dz \int_0^\infty dm_1 \int_0^1 dq \frac{d^5 N}{dz dm_1 dq dt_r} \frac{dt_r}{d \ln f_{K,r}} \times h^2(f_{K,r}) \sum_{n=1}^{\infty} \frac{g[n, e(f_{K,r})]}{(n/2)^2} \Big|_{f_{K,r}=f(1+z)/n}$$

Circular binaries:

$$h_c^2(f) = \int_0^\infty dz \int_0^\infty d\mathcal{M} \frac{d^3 N}{dz d\mathcal{M} d \ln f} h^2(f)$$

$$h_c^2(f) = \frac{4G^{5/3}}{3\pi^{1/3}c^2} f^{-4/3} \int d\mathcal{M} \int dz (1+z)^{-1/3} \mathcal{M}^{5/3} \frac{d^2 n}{dz d\mathcal{M}} \quad h_c \propto f^{-2/3}$$

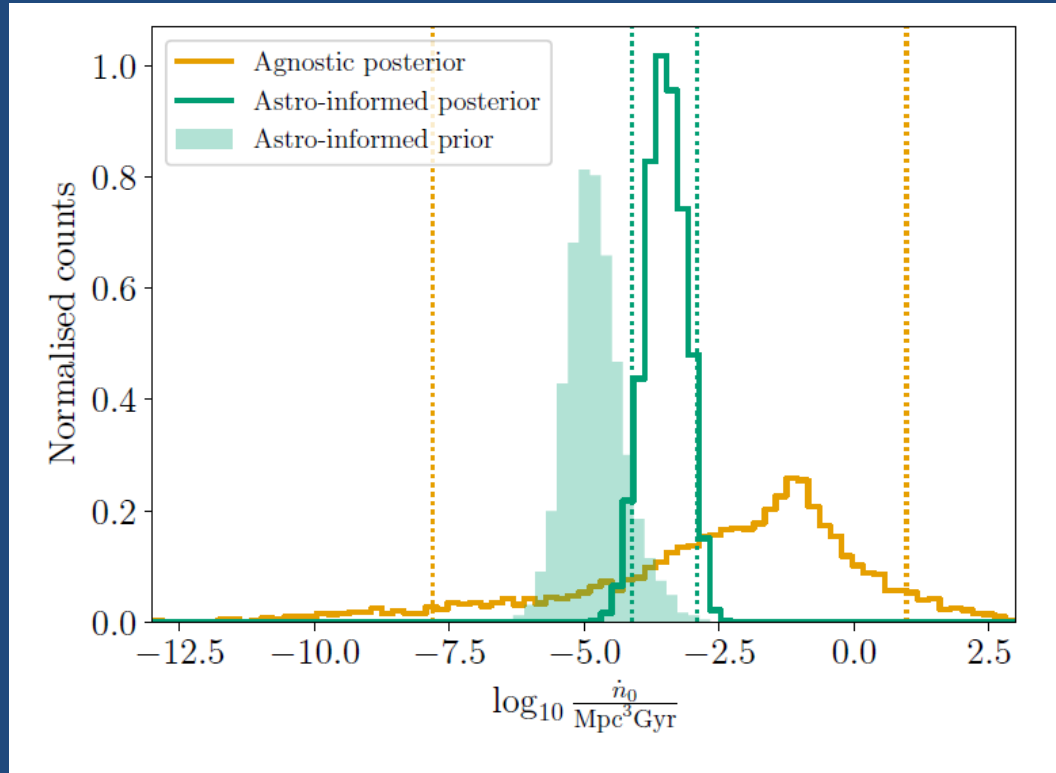
IPTA+InPTA results



$\gamma=13/3$ for binary mergings

Antoniadis+ 2306.16227 (AA 2024)

Implications from EPTA+InPTA



Antoniadis+ 2306.16227

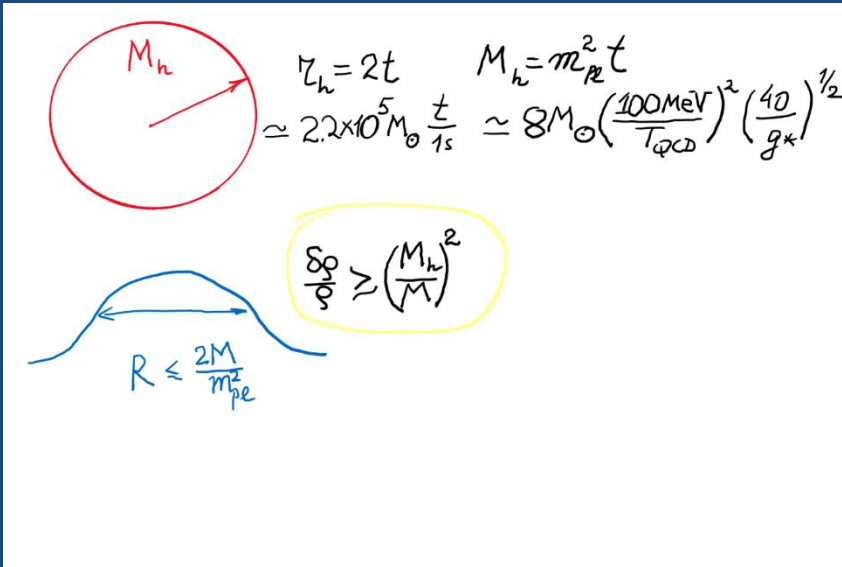
$$\frac{d^2n}{dzd \log_{10} \mathcal{M}} = \dot{n}_0 \left[\left(\frac{\mathcal{M}}{10^7 M_\odot} \right)^{-\alpha_{\mathcal{M}}} e^{-\mathcal{M}/M_\star} \right] \left[(1+z)^{\beta_z} e^{-z/z_0} \right] \frac{dt_R}{dz}$$

Cf: SMBH merger rate $10^{-2} - 10^{-3} \text{ yr}^{-1} \text{ Gpc}^{-3}$ up to $z=4$ (Wang, Zhang 2403.06416)

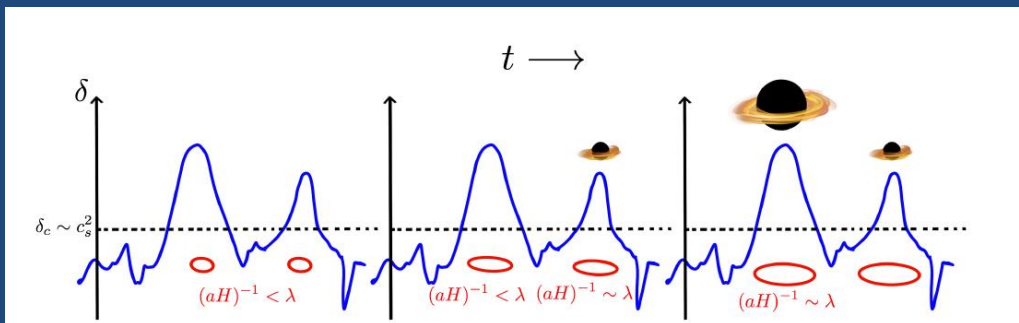
Primordial BH formation

Zeldovich, Novikov'67, Carr, Hawking'74....

- PBH mass = mass inside cosmological horizon



Dolgov+KP 2004.1699



Some PBH formation scenarios

- Primordial density fluctuations, $\delta > \delta_c \sim 0.45$ (model-dependent)
- Collapse of scale-invariant fluctuations \rightarrow power-law mass spectra $dn/dM \sim M^{-\alpha}$, $\alpha = 2(1+2w)/(1+w) = 2.5$ at RD ($w=1/3$)
- Inflationary models \rightarrow **log-normal mass spectrum**

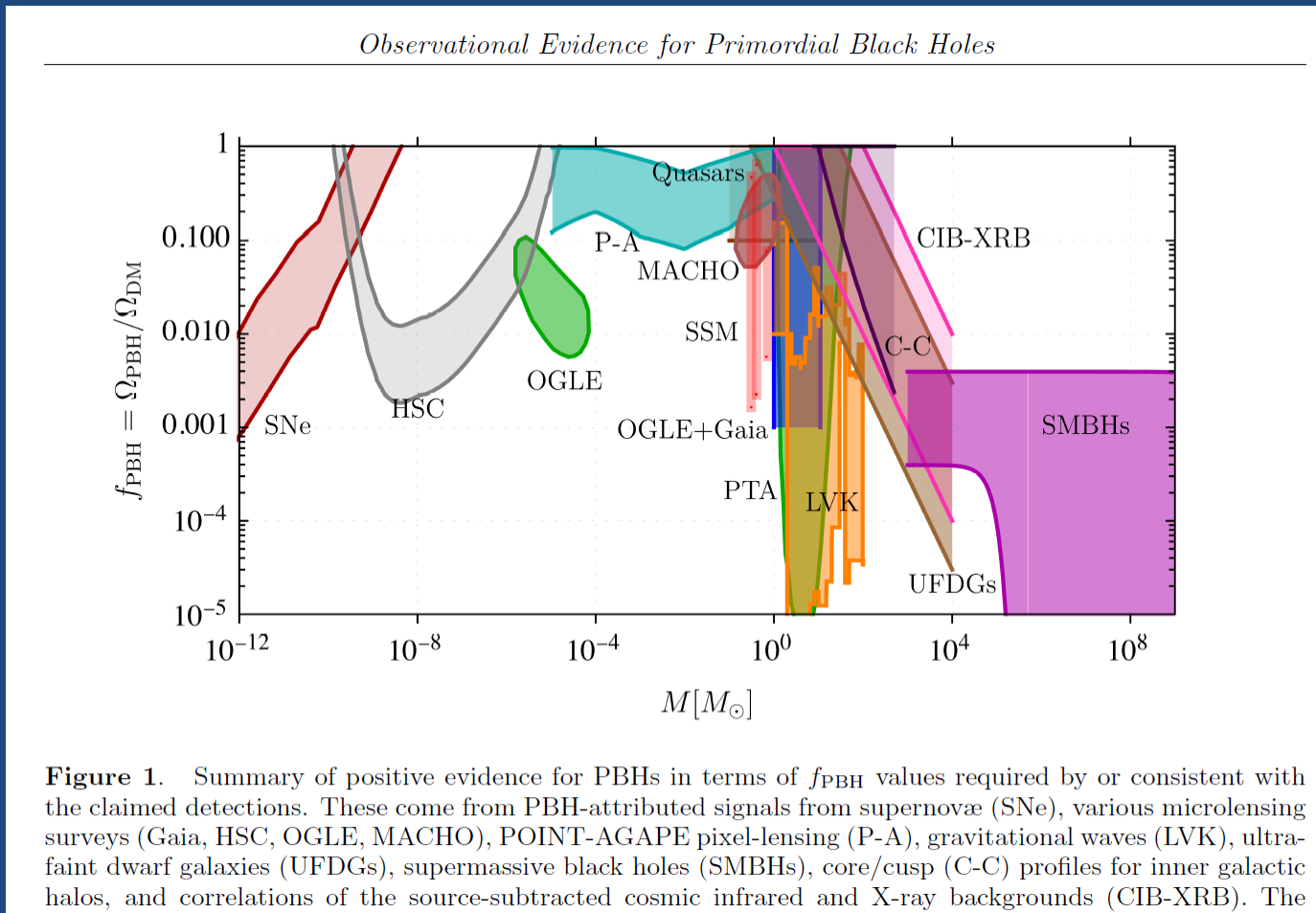
$$\frac{dn}{dM} = \mu^2 \exp[-\gamma \ln^2(M/M_{max})] \quad (\text{Dolgov, Silk 1993})$$

- QCD phase transition at $T=100-200$ MeV:

$$M \sim \left(\frac{m_{pl}}{m_p}\right)^3 m_p \sim M_{ch} \sim 1M_{\odot}$$

- Bubble collisions, cosmic strings ...

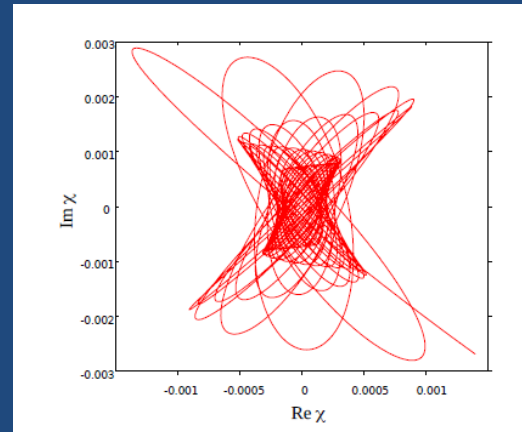
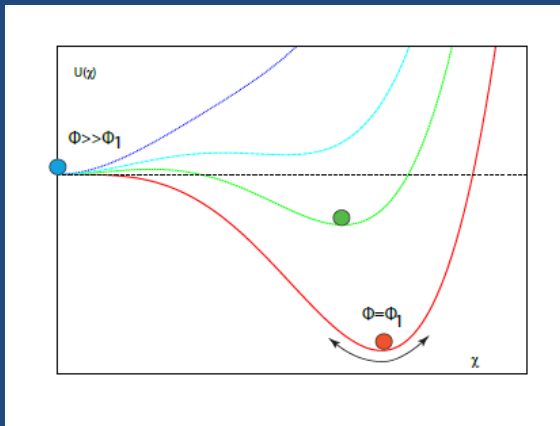
Positive evidence for PBHs?



Modified Affleck-Dine baryogenesis (Dolgov and Silk , 1993)

- Complex scalar baryon field χ + Inflaton field Φ

$$U(\chi, \Phi) = U_{\Phi}(\Phi) + U_{\chi}(\chi) + \lambda_1(\Phi - \Phi_1)^2|\chi|^2$$



- Along 'flat directions', field χ rotates \rightarrow regions with baryonic charge (HBB) by the end of inflation
- HBB turns into density fluctuations after QCD phase transition at $T \sim 100$ MeV with **log-normal mass spectrum**

$$\frac{dn}{dM} = \mu^2 \exp[-\gamma \ln^2(M/M_{max})]$$

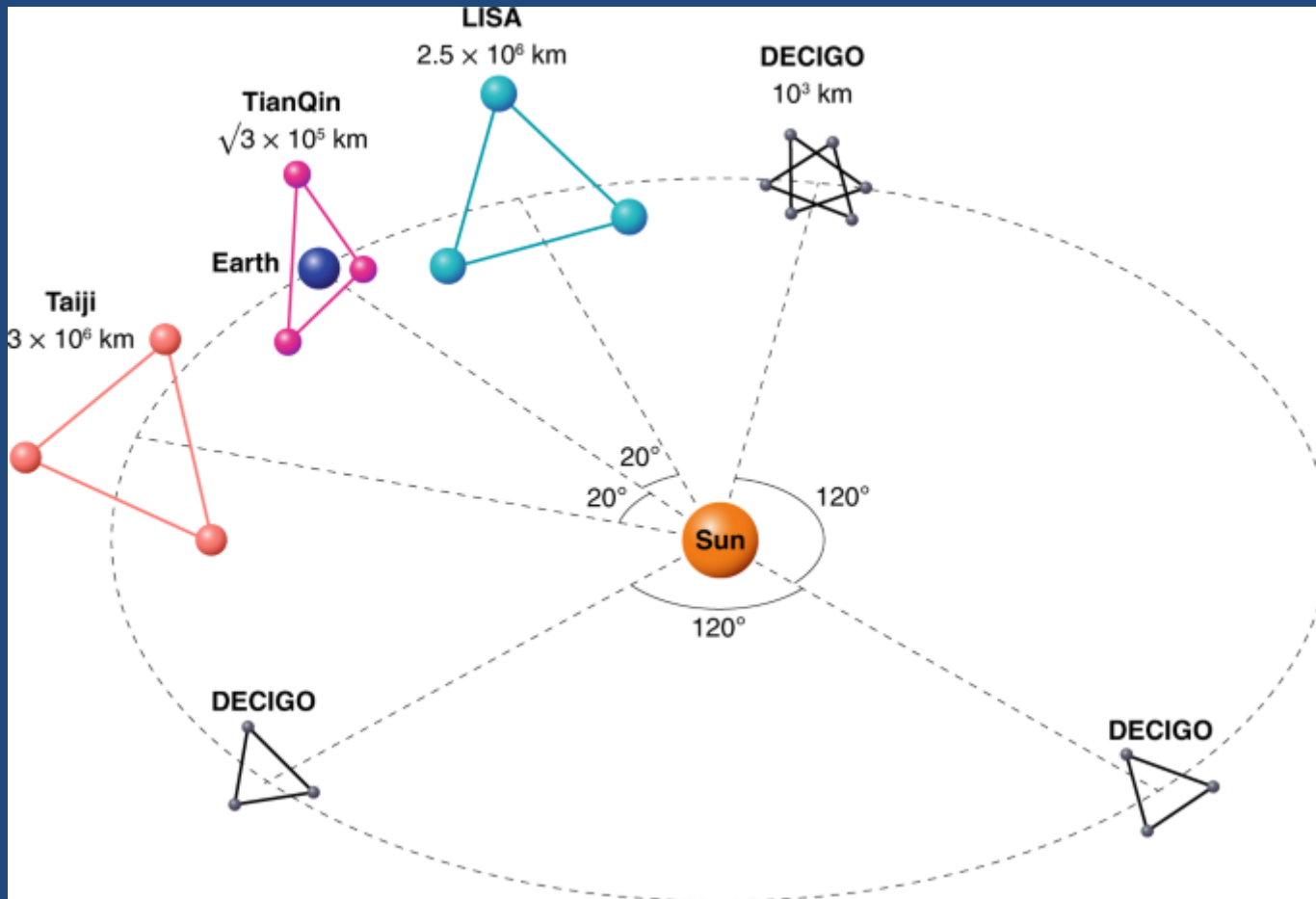
Primordial IMBH

- Maximum BH mass depends on the duration of inflation after ‘opening’ of flat directions of the χ potential

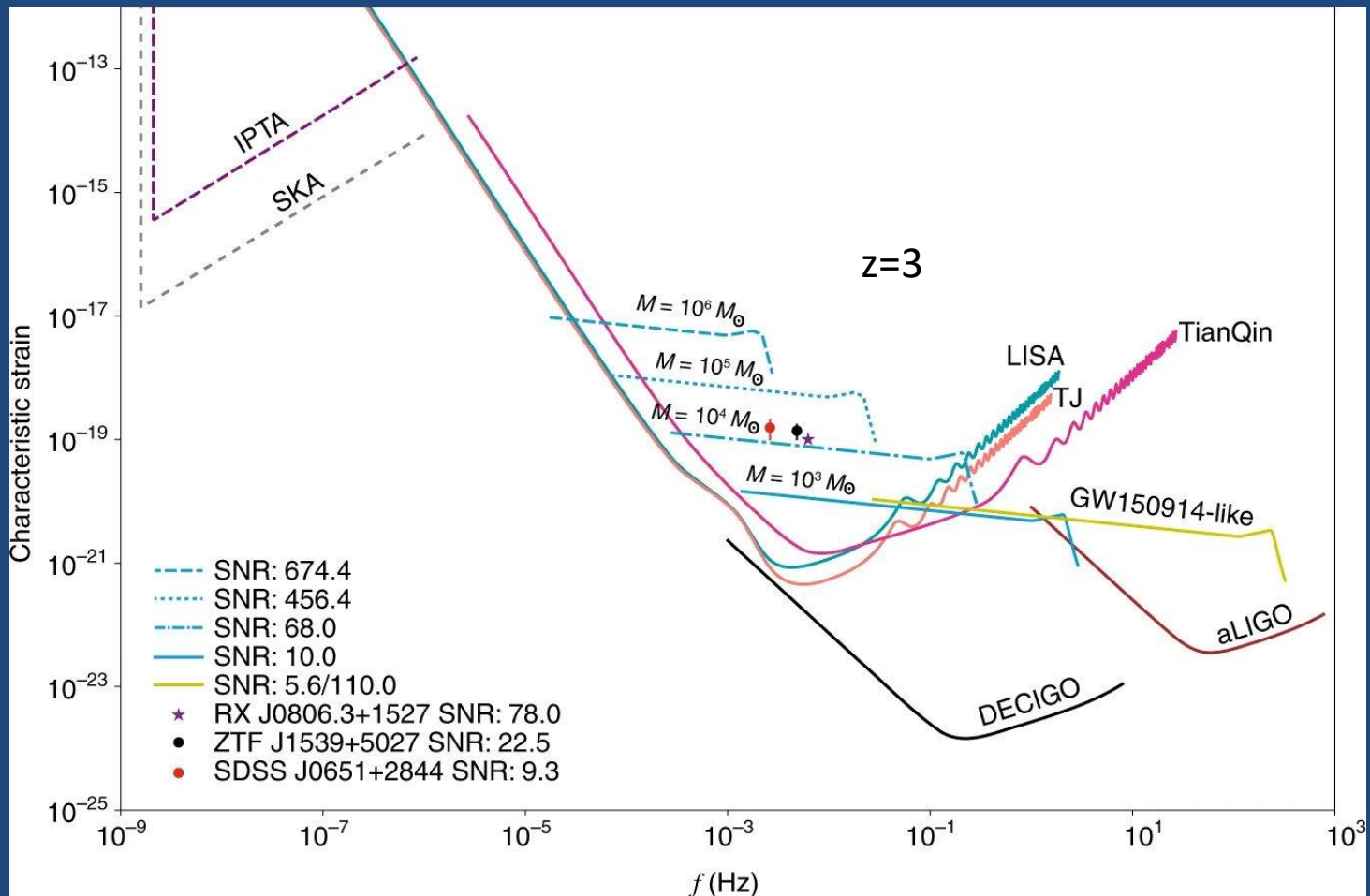
$$e^{3H(t_e - t_1)} = 10^{37+5-10+15} \left(\frac{T_h}{10^{14} \text{GeV}} \right)^3 \frac{M_{\text{BH}}^{(\text{max})}}{10^4 M_{\odot}} \frac{100 \text{ MeV}}{T_{\text{MD}}}$$

- $H(t_e - t_1) = 36$ is required to create PBH up to $10^4 M_{\odot}$ (Blinnikov, Dolgov, Porayko, PK 2016 JCAP)

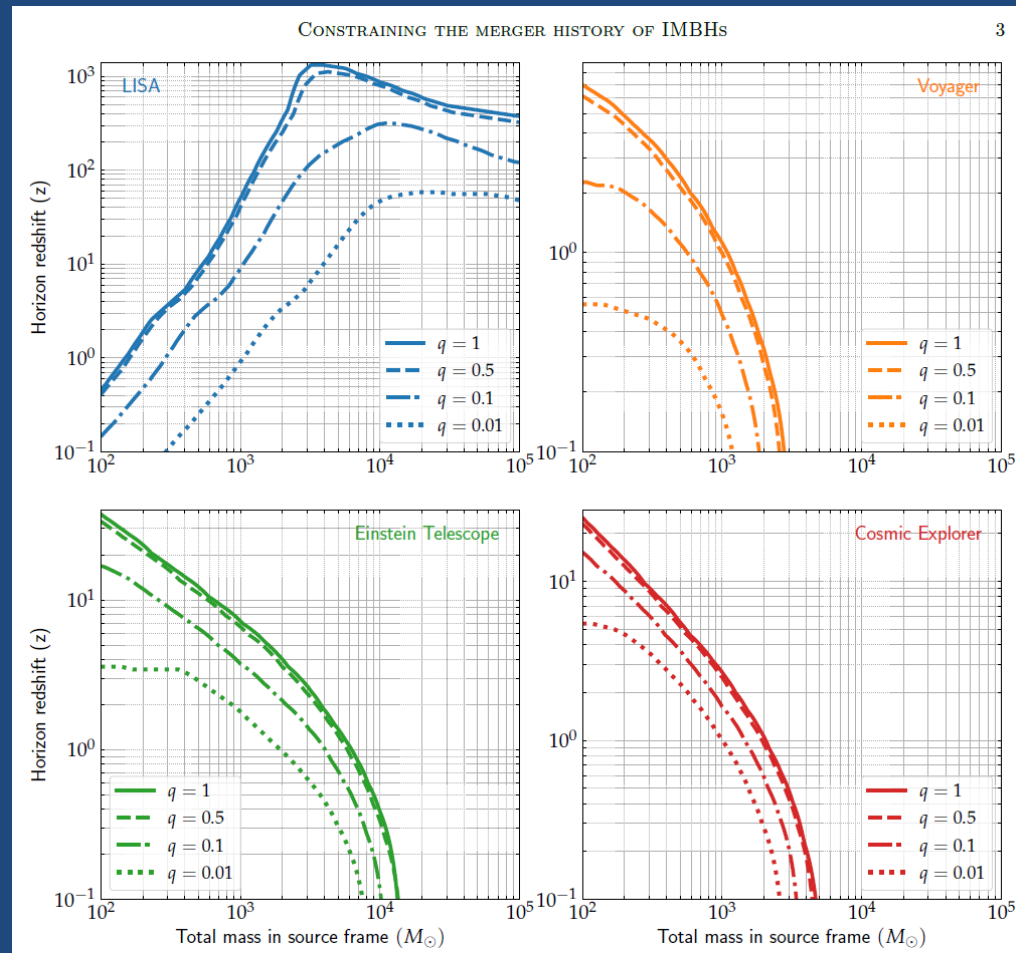
Space laser interferometers



Coalescing binary IMBH



Detection horizon for merging binary IMBH



TianQin project (SYSU, China)

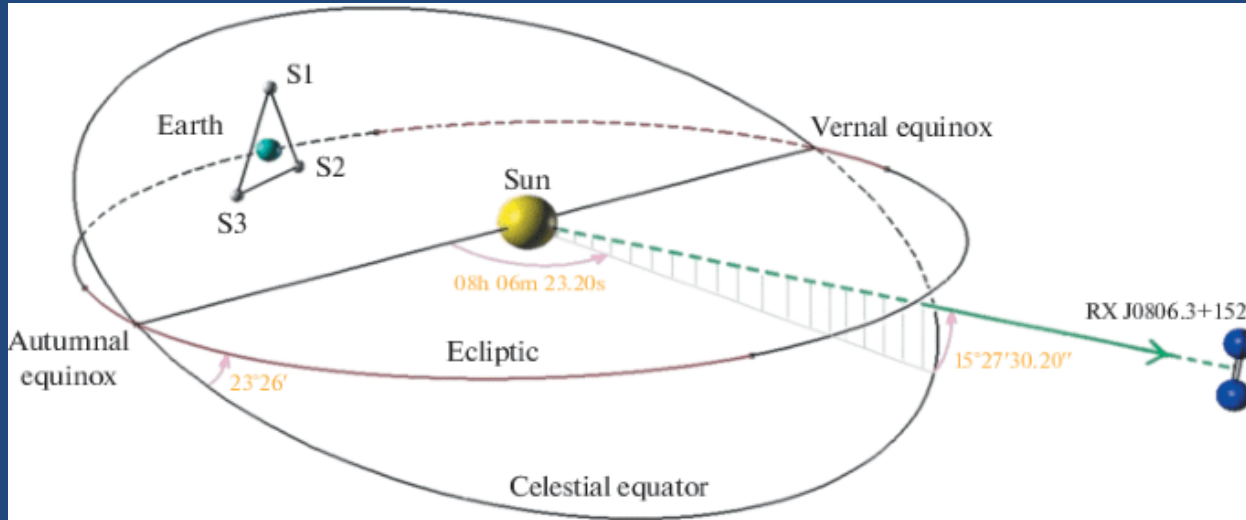
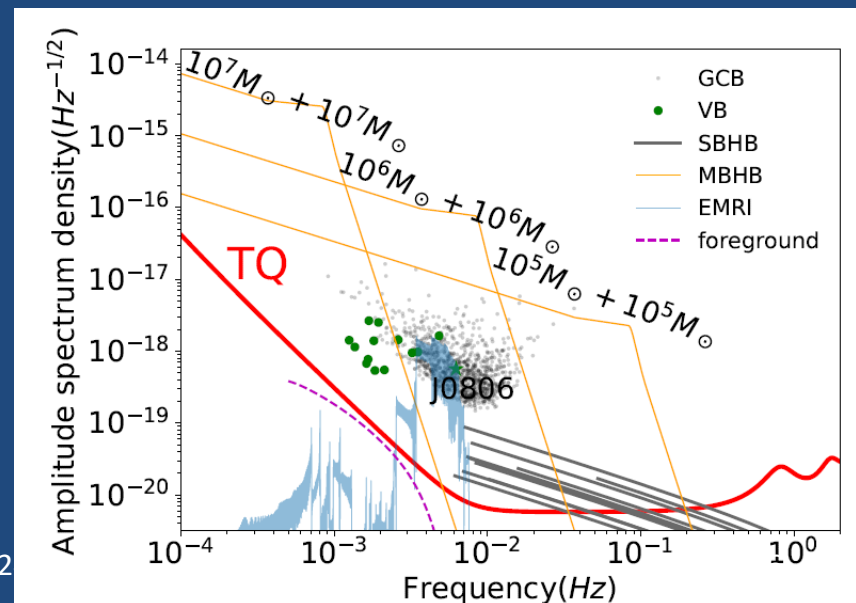


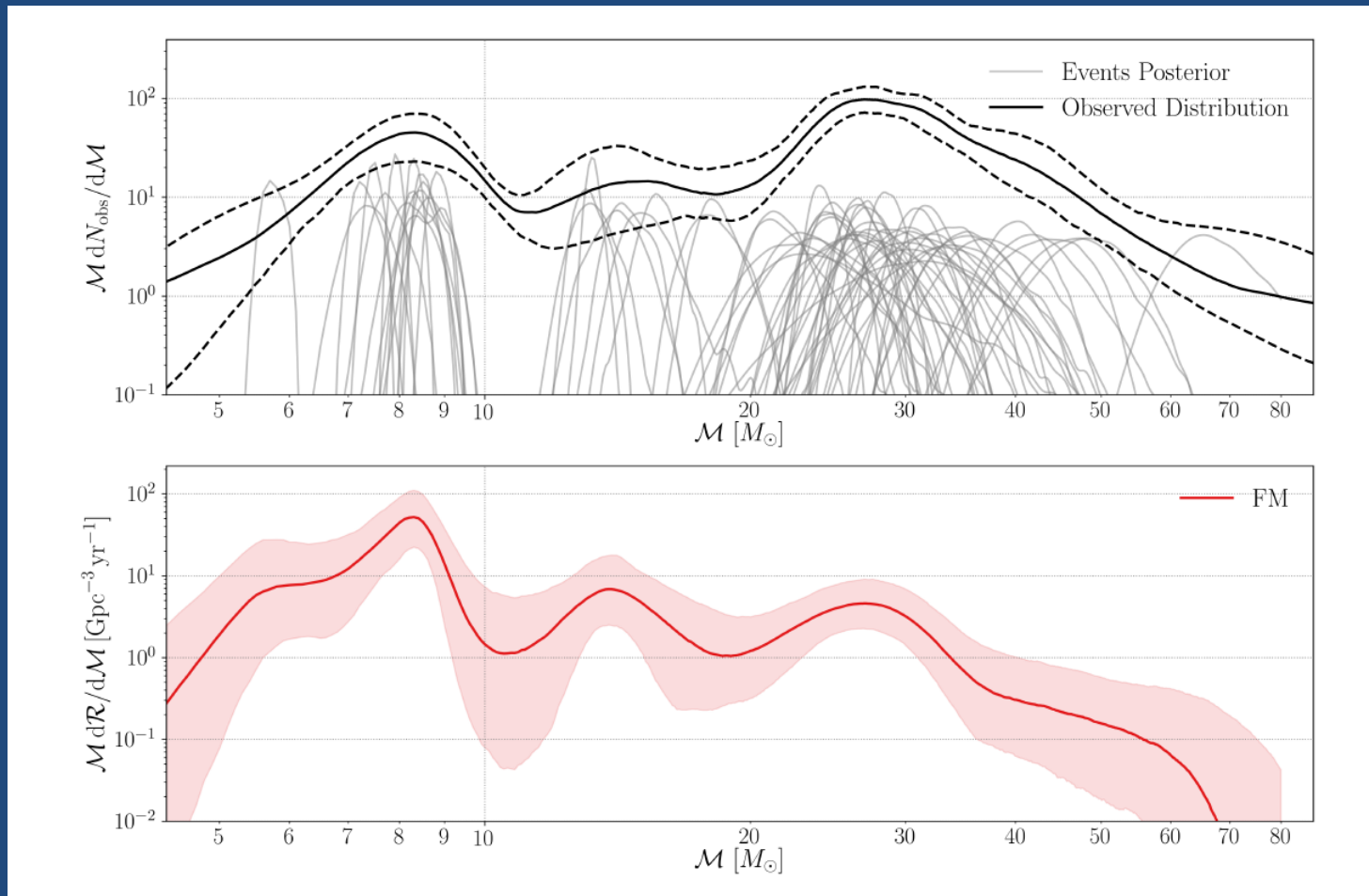
Table 1. Basic parameters of TianQin.

Parameters	Value
Type of orbit	Geocentric
Number of satellites	$N=3$
Arm length	$L = \sqrt{3} \times 10^5$ km
Displacement measurement noise	$S_x^{1/2} = 1 \times 10^{-12}$ m/Hz ^{1/2}
Residual acceleration noise	$S_a^{1/2} = 1 \times 10^{-15}$ m s ⁻² /Hz ^{1/2}



Connecting LVK coalescing BH+BH with IMBH

GWTC-3 features (wait for O4 results)



Binary PBH (distinctive feature: zero spins, only sources with $z > 20$)

- Merging rate (e.g. Raidal+ 2404.08416)

$$\frac{d\mathcal{R}(z)}{dm_1 dm_2} = \frac{3.2 \times 10^6}{\text{Gpc}^3 \text{yr}^{-1}} S(m_1, m_2, z) f_{pbh}^{\frac{53}{57}} F(m_1) F(m_2) \frac{m_1 m_2}{\langle m \rangle^2} \left(\frac{t(z)}{t_0} \right)^{-\frac{34}{37}} M^{-\frac{32}{37}} \eta^{-\frac{34}{37}}$$

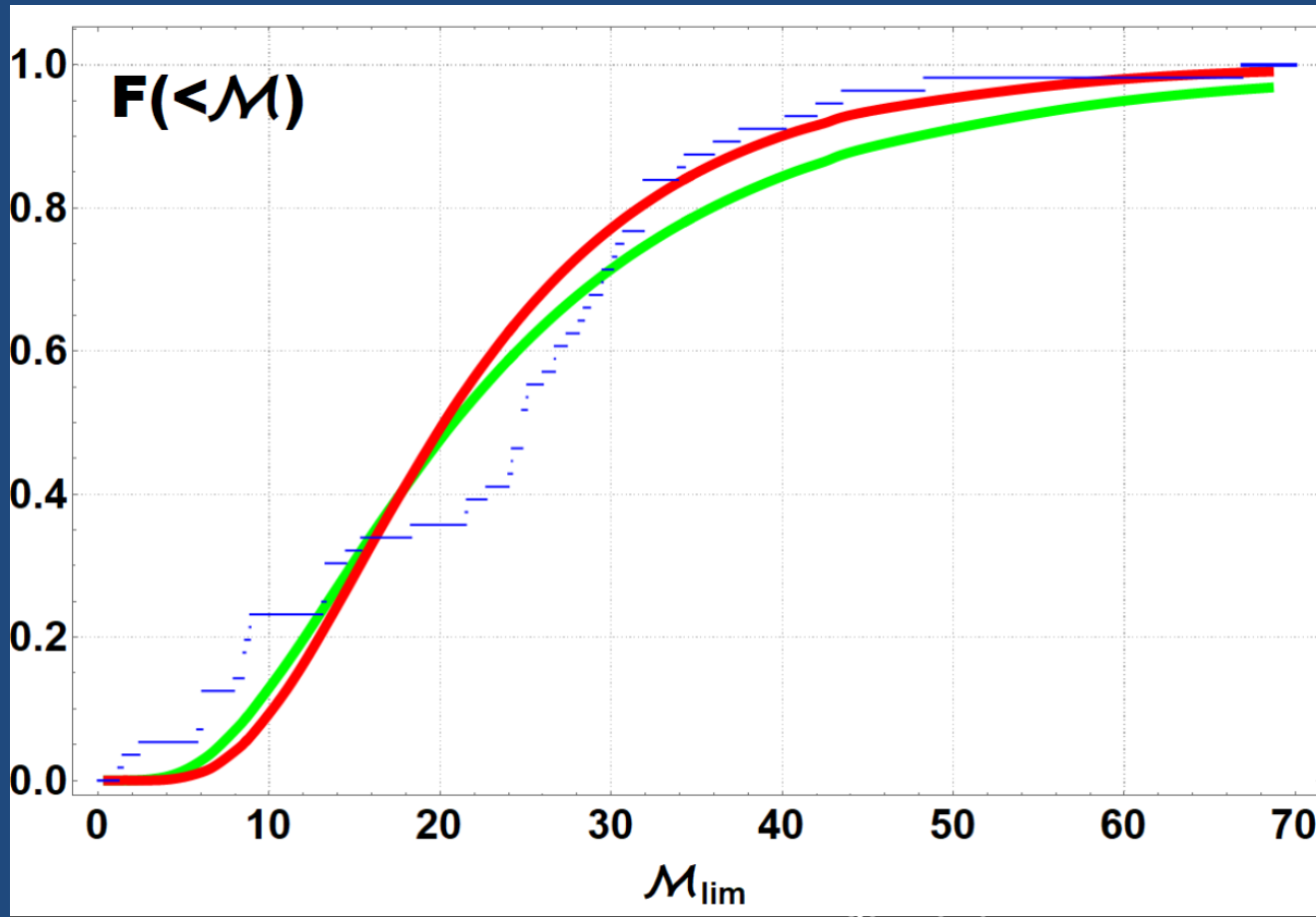
$$\frac{d\mathcal{R}(z)}{d\mathcal{M} dq} \propto f_{pbh}^{\frac{53}{57}} F(\mathcal{M}) F(q) \frac{(1+q)^{\frac{8}{5}}}{q^{\frac{9}{5}}} \frac{\mathcal{M}^{\frac{79}{37}}}{M_c^2 e^{\frac{1}{2\gamma}}} \left(\frac{t(z)}{t_0} \right)^{-\frac{34}{37}}$$

- Chirp mass cumulative distribution

$$\mathcal{R}(< \mathcal{M}) = \int_0^{\mathcal{M}} \int_0^{z_{\text{lim}}(\mathcal{M})} \int_0^1 \frac{d\mathcal{R}}{d\mathcal{M} dq} \frac{dV}{dz} \frac{dz}{1+z} dq, \quad F(< \mathcal{M}) = \frac{\mathcal{R}(< \mathcal{M})}{\mathcal{R}(< \infty)}$$

Example: log-normal PBH mass function fits GWTC data

$$F(M) = A \exp[-\gamma \ln^2(M/M_0)]$$



$$\gamma = 0.7, M_0 = 19$$

$$DR(M_{ch})dM_{ch} \sim Ae^{-B \ln^2\left(\frac{M_{ch}}{M_0}\right)}dM_{ch}$$

Dolgov+'19

Independent
analysis:

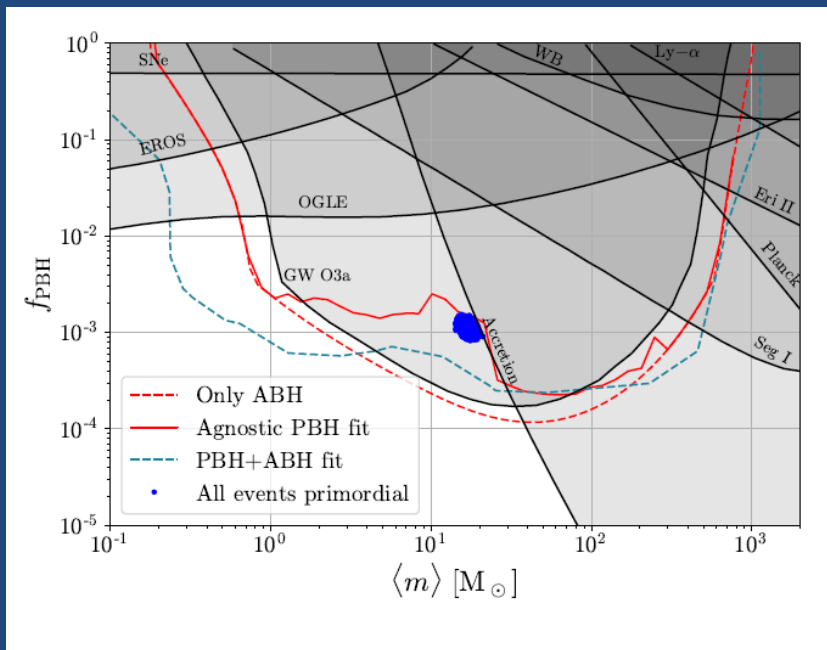
Lang Liu et al
2023

(2210.16094)

$M_0 \sim 17, \gamma \sim 1,$

$f_{\text{pbh}} \sim 10^{-3}$

PBH abundance constraints from GWTC-3



$$\psi(m) = \frac{1}{\sqrt{2\pi}\sigma} \exp\left\{-\frac{\ln^2(m/m_c)}{2\sigma^2}\right\}$$

$$\langle m \rangle = m_c e^{-\sigma^2/2}$$

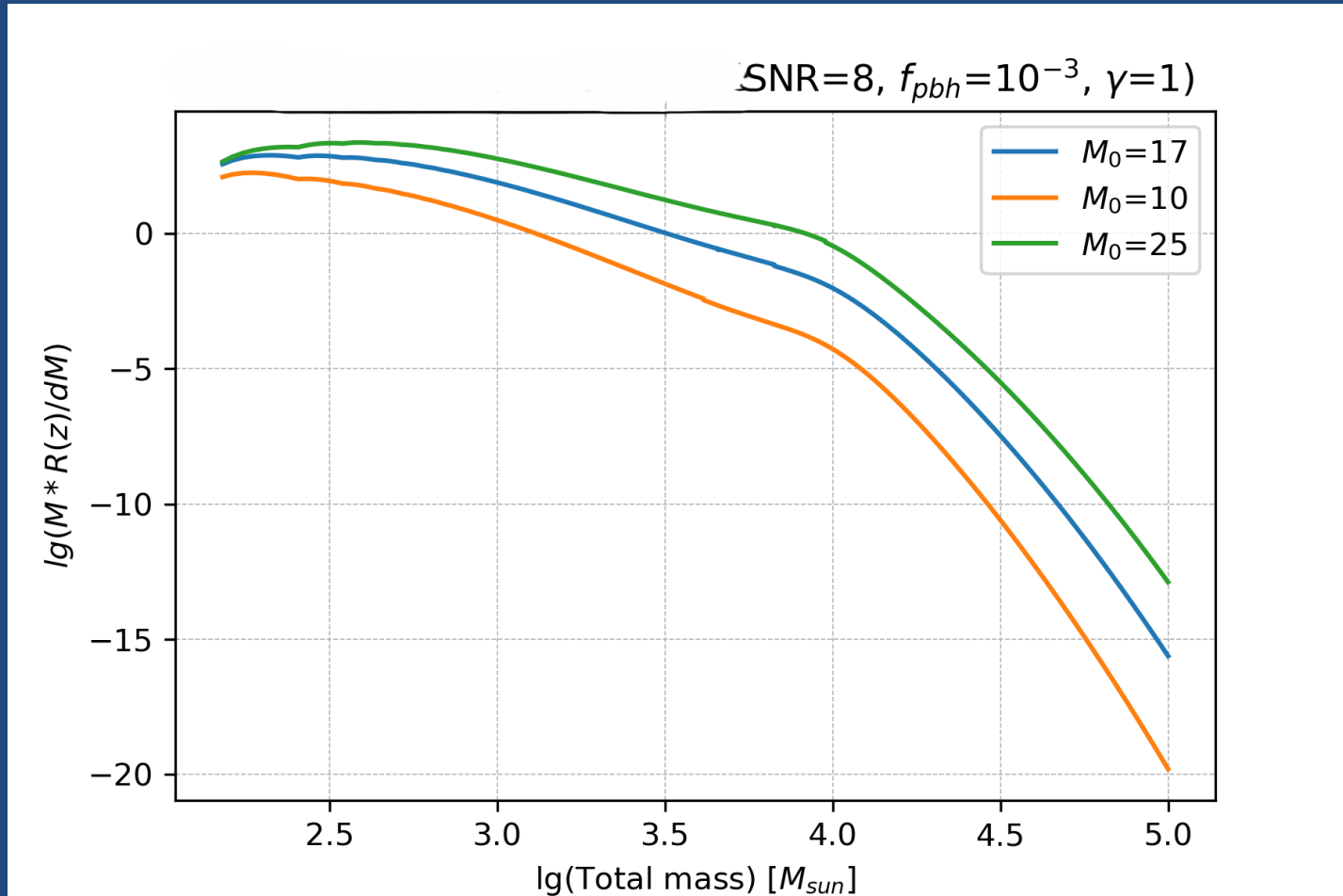
$$M_0 = 17 M_\odot \quad \sigma = 0.6$$

Andres-Carcasona et al 2405.05732

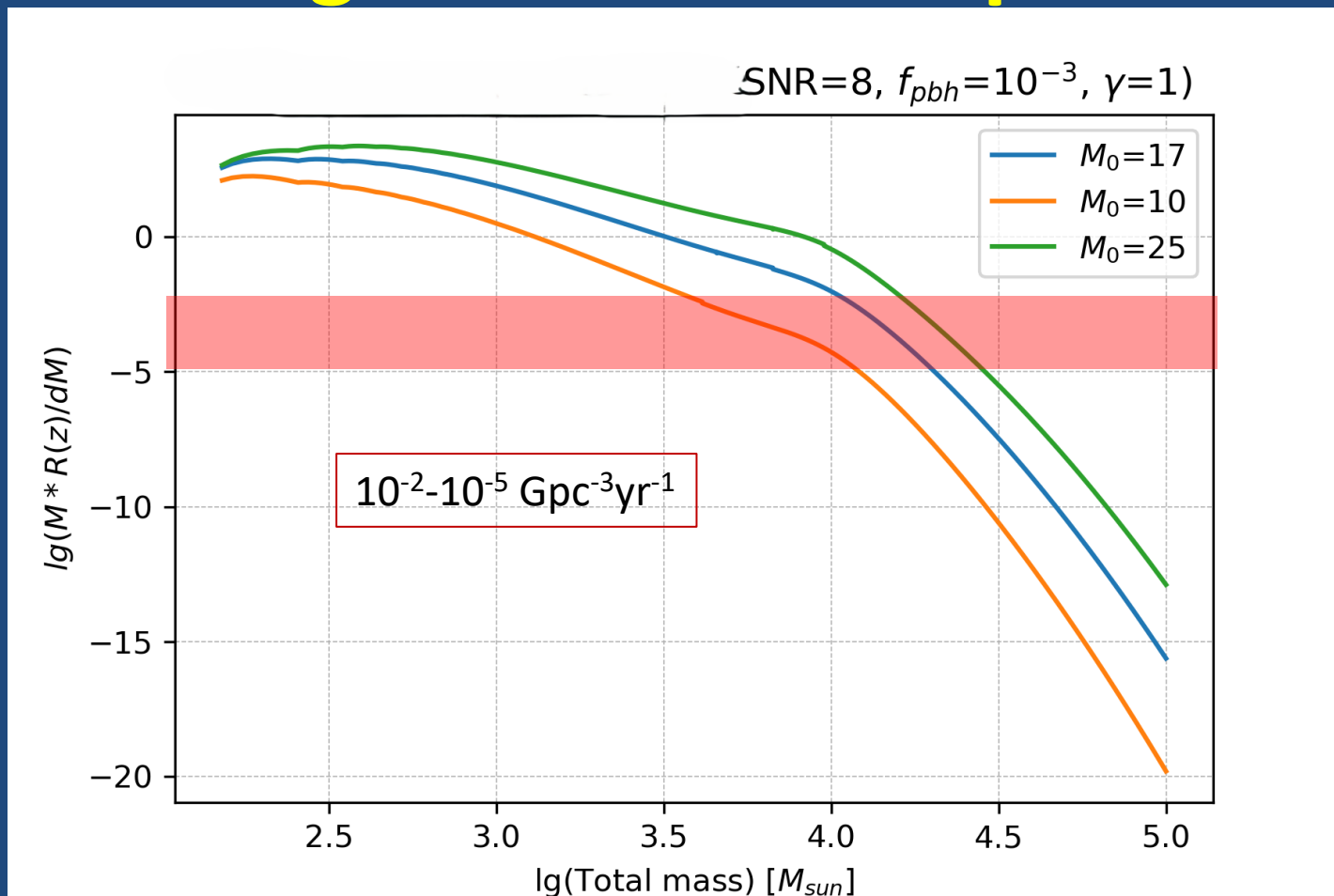
- Initial log-normal PBH mass function \rightarrow log-normal PBH mass distribution
- Rate estimate

$$m = \frac{M}{M_0} \quad \frac{d\mathcal{R}(z)}{dm dq} \sim F(q)F(m)e^{-2\gamma \ln^2(m)} \left(\frac{t(z)}{t_0} \right)^{-34/37}$$

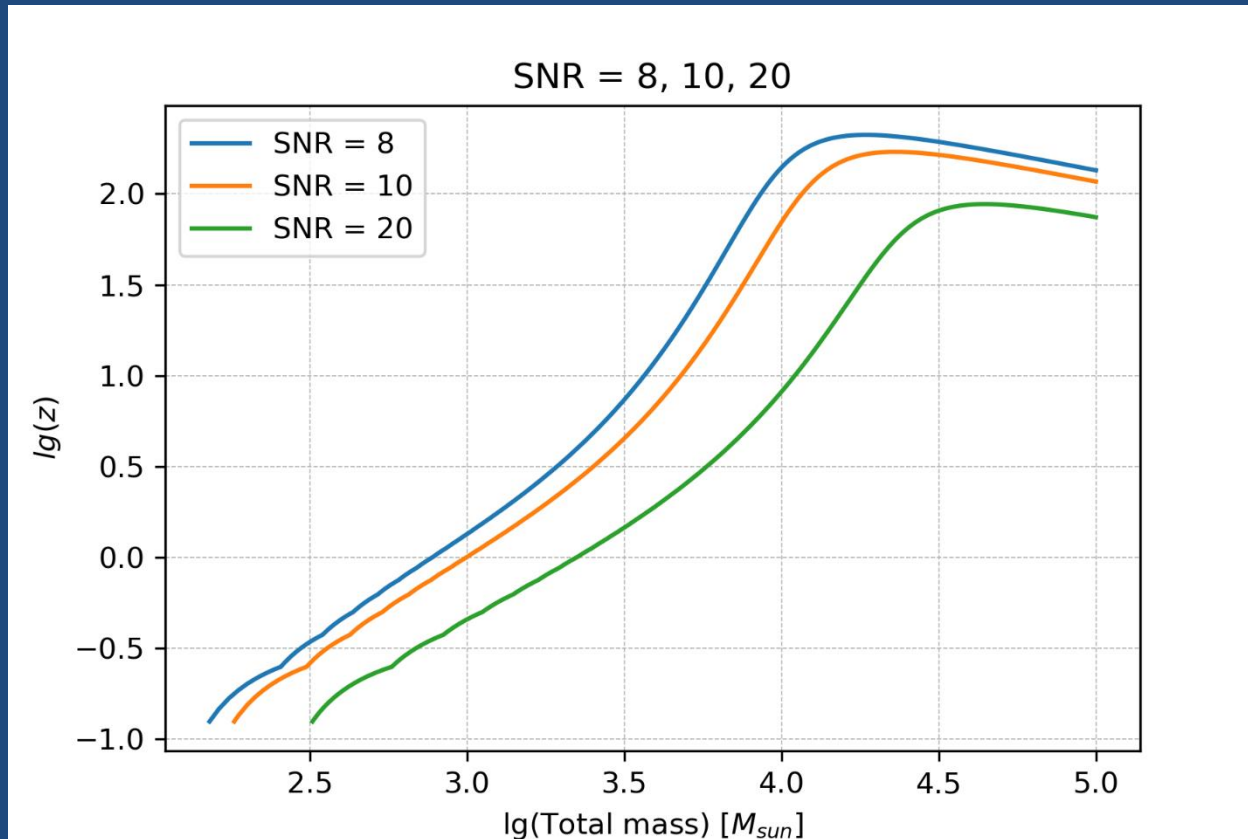
MdR/dM [$\text{Gpc}^{-3} \text{yr}^{-1}$], $z=0$



Comoving merging rate of binary PBH with log-normal mass spectrum



TQ detection horizon for BH+BH

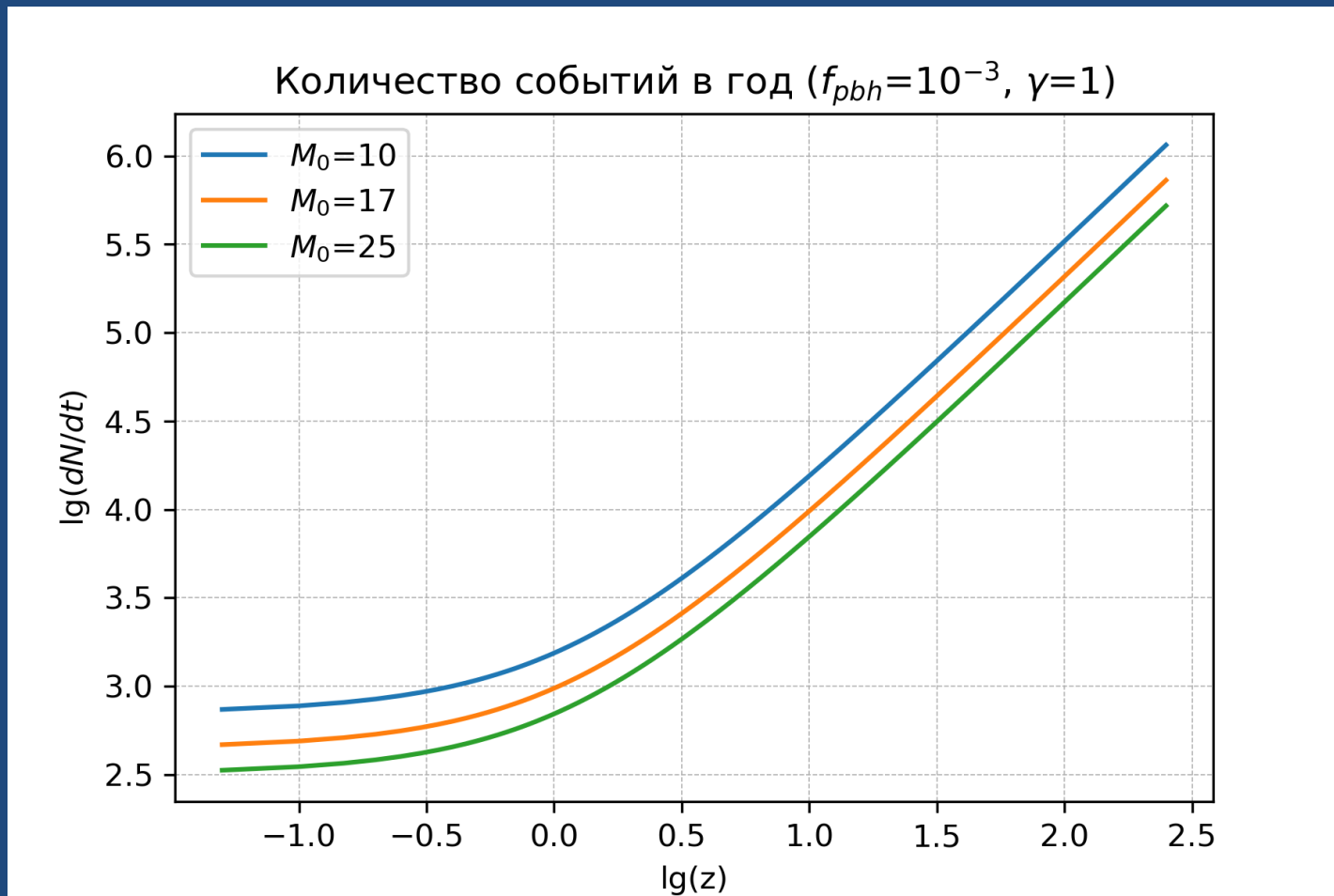


$$SNR \propto \frac{[\mathcal{M}/(1+z)]^{5/6}}{d_l(z)}$$

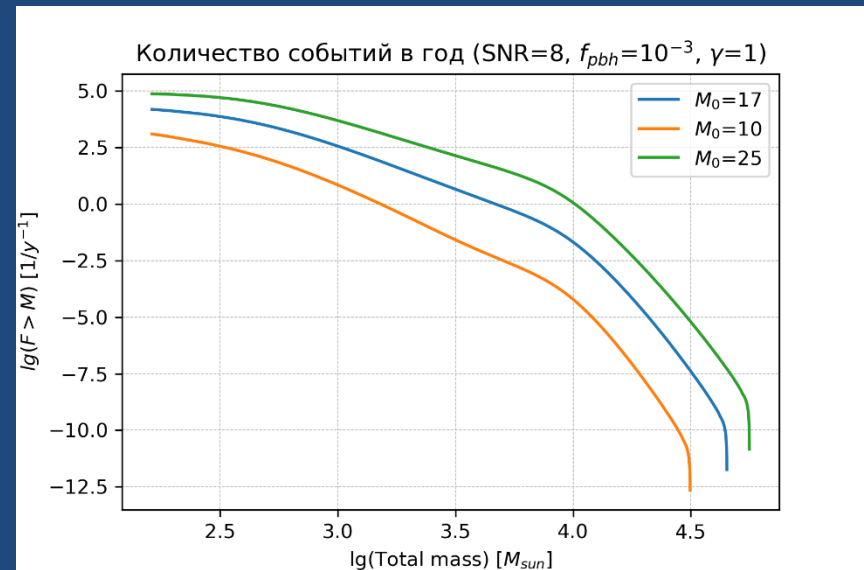
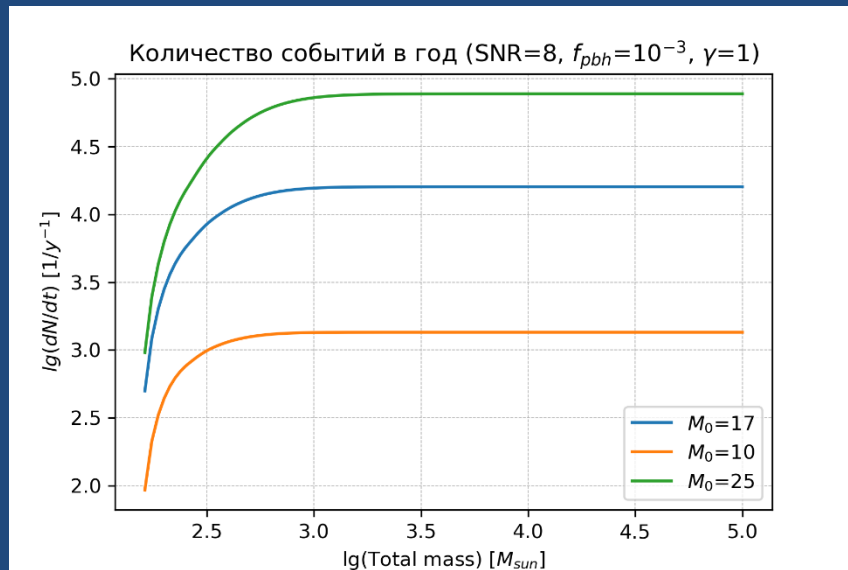
$$\mathcal{M} = \frac{(M_1 M_2)^{3/5}}{M^{1/5}}$$

(chirp-mass)

$$dN/dt/dz \sim t(z)^{-34/37}$$



Expected TQ detection rate (yr^{-1})



Conclusions

- Astrophysical IMBH with masses $\sim 1000-10000 M_{\odot}$ are (almost) found in centers of dwarf galaxies
- Emerging evidences for **binary IMBHs** in galactic centers (JWST observations)
- Binary IMBHs are among primary sources for future space laser GW interferometers
- **IMBHs can have primordial origin.** In Dolgov&Silk (1993) baryogenesis, IMBHs can have masses up to $10^4 M_{\odot}$ and can be seeds for early SMBH formation in galactic centers.
- **Distinctive feature of binary primordial IMBH – zero spins & high redshift ($z>20$)**
- Assuming $\sim 0.1\%$ of DM in PBH with log-normal mass spectrum consistent with LVK detections, IMBH mergings with masses $> 1000 M_{\odot}$ can be detected in a few year of observations by planned laser GW interferometers like LISA or TianQin.