

KK compactifications of Horndeski theories

based on 2405.02281

S. Mironov, A. Shtennikova, M. Valencia-Villegas

INR RAS, ITMP, ITEP (NRCKI)

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$$R^5 \longrightarrow R^4 \times S^1$$

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- * Metric + scalar \longrightarrow Metric + vector + scalar + scalar

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Natural way to obtain vector Galileon Lagrangian.

Some uniformity in vector and tensor sectors.

Results

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Vector Galileons

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Non-gauge vector galileons:

Examples: 1307.0077 (G. Tasinato, K. Koyama, and N. Khosravi)..., 1812.11134 (P. Petrov)

Gauge vector galileons:

No-go: 1312.6690 (C. Deffayet, A. E. Gümrükçüoğlu, S. Mukohyama, and Y. Wang)

Latest developments: 2312.14814, 2404.18715 (A. Colléaux, D. Langlois, and K. Noui)

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$$\text{GW170817: } \left| \frac{c_T}{c} - 1 \right| < 10^{-15}$$

$$\text{KK compactification: } \left| \frac{c_T}{c} - 1 \right| \equiv 0$$

Horndeski theory

$$S = \int d^4x \sqrt{-g} (\mathcal{L}_2 + \mathcal{L}_3 + \mathcal{L}_4 + \mathcal{L}_5),$$

$$\mathcal{L}_2 = F(\pi, X),$$

$$\mathcal{L}_3 = K(\pi, X) \square \pi,$$

$$\mathcal{L}_4 = -G_4(\pi, X)R + 2G_{4X}(\pi, X) \left[(\square\pi)^2 - \pi_{;mn}\pi^{;mn} \right],$$

$$\mathcal{L}_5 = G_5(\pi, X)G^{mn}\pi_{;mn} + \frac{1}{3}G_{5X} \left[(\square\pi)^3 - 3\square\pi\pi_{;mn}\pi^{;mn} + 2\pi_{;mn}\pi^{;ml}\pi_{;l}{}^n \right],$$

where π is the Galileon field, $X = g^{mn}\pi_{,m}\pi_{,n}$, $\pi_{,m} = \partial_m\pi$, $\pi_{;mn} = \nabla_n\nabla_m\pi$,
 $\square\pi = g^{mn}\nabla_n\nabla_m\pi$, $G_{4X} = \partial G_4/\partial X$.

$$g_{mn} = \begin{pmatrix} g_{\mu\nu} + \phi^2 A_\mu A_\nu & \phi^2 A_\mu \\ \phi^2 A_\nu & \phi^2 \end{pmatrix}$$

Horndeski theory \rightarrow U(1) vector gauge field

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$\mathcal{L}_2, \mathcal{L}_3 \rightarrow$ no dynamical A_μ

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$\mathcal{L}_4 \rightarrow$ non minimal coupling to A_μ

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$\mathcal{L}_5, \mathcal{L}_6 \rightarrow A_\mu$ is a vector Galileon

Horndeski theory

Speed of graviton

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$$\mathcal{S}_T = \int d\eta d^3x a^4 \left[\frac{1}{2a^2} \left(\mathcal{G}_\tau \left(\dot{h}_{ij} \right)^2 - \mathcal{F}_\tau (\partial_k h_{ij})^2 \right) \right]$$

$$\mathcal{G}_T = 2[G_4 - 2XG_{4,X} - X(H\dot{\pi}G_{5X} - G_{5\phi})]$$

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$$c_T^2 = c_v^2 = 1$$

$$\mathcal{G}_T = 2 \left[G_4 - \cancel{2XG_{4,X}} - \cancel{X(H\dot{\pi}G_{5X} - G_{5\phi})} \right]$$

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$$c_T^2 = c_v^2 \neq 1$$

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allowable Horndeski theory

$$S = \int d^4x \sqrt{-g} (\mathcal{L}_2 + \mathcal{L}_3 + \mathcal{L}_4),$$

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$$\mathcal{L}_4 = -G_4(\pi, X)R$$

allowable Horndeski theory

$$S = \int d^4x \sqrt{-g} (\mathcal{L}_2 + \mathcal{L}_3 + \mathcal{L}_4),$$

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$$\mathcal{L}_5 = G_5(\pi) G^{\mu\nu} \pi_{;\mu\nu}$$

$$\mathcal{S}_{5d} = \int d^5x \sqrt{-g_{(5)}} \left(G_4 R_{(5)} + G_{4,X} \left((\square\pi)^2 - (\nabla_A \nabla_B \pi)^2 \right) \right)$$

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$$\mathcal{S}_{4d} = \int d^4x \sqrt{-g} \phi G_4 R + \phi G_{4,X} \left((\square\pi)^2 - (\nabla_\mu \nabla_\nu \pi)^2 \right)$$

$$+ G_4 \left(-\frac{1}{4} \phi^3 F^2 - 2 \square \phi \right) + G_{4,X} \left(2 \nabla_\mu \phi \nabla^\mu \pi \square \pi - \frac{1}{2} \phi^3 F_\mu^\sigma F_{\nu\sigma} \nabla^\mu \pi \nabla^\nu \pi \right)$$

$$\int d^5x \sqrt{-g_{(5)}} \left(G_5(\pi) \left(R_{(5)}^{AB} - \frac{1}{2} g_{(5)}^{AB} R_{(5)} \right) \nabla_A \nabla_B \pi \right)$$

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$$\begin{aligned} & \int d^4x \sqrt{-g} \phi G_5(\pi) \left(\left(R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R \right) \nabla_\mu \nabla_\nu \pi \right. \\ & - \frac{1}{2\phi} R \nabla_\mu \phi \nabla^\mu \pi + \frac{1}{\phi} (\square \phi \square \pi - \nabla_\mu \nabla_\nu \phi \nabla^\mu \nabla^\nu \pi) \\ & + \frac{1}{8} \phi F^{\mu\nu} F^{\sigma\rho} \left(3 g_{\nu\rho} (-4 g_{\lambda\mu} g_{\beta\sigma} + g_{\lambda\beta} g_{\mu\sigma}) \nabla^\lambda \pi \nabla^\beta \phi \right. \\ & \left. \left. + \phi g_{\sigma\mu} (-4 \nabla_\nu \nabla_\rho \pi + g_{\rho\nu} \square \pi) \right) + \frac{1}{2} \phi^2 F_{\mu\nu} \nabla_\sigma F^{\nu\sigma} \nabla^\mu \pi \right) \end{aligned}$$

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Thank you for your attention!