

# Primordial Black Holes

- a couple of recent topics -

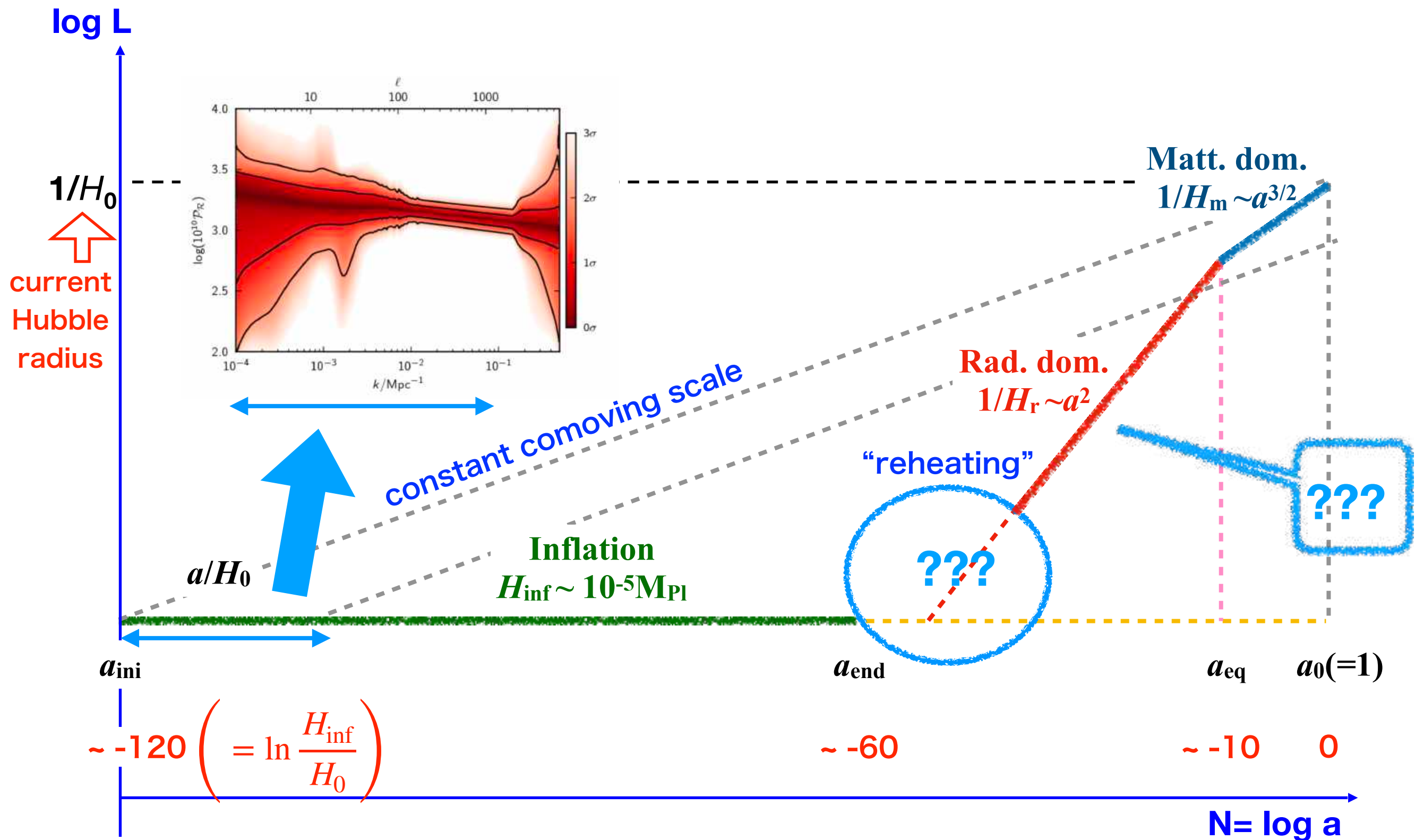
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YITP, Kyoto University  
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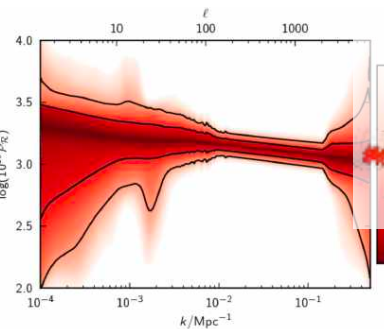
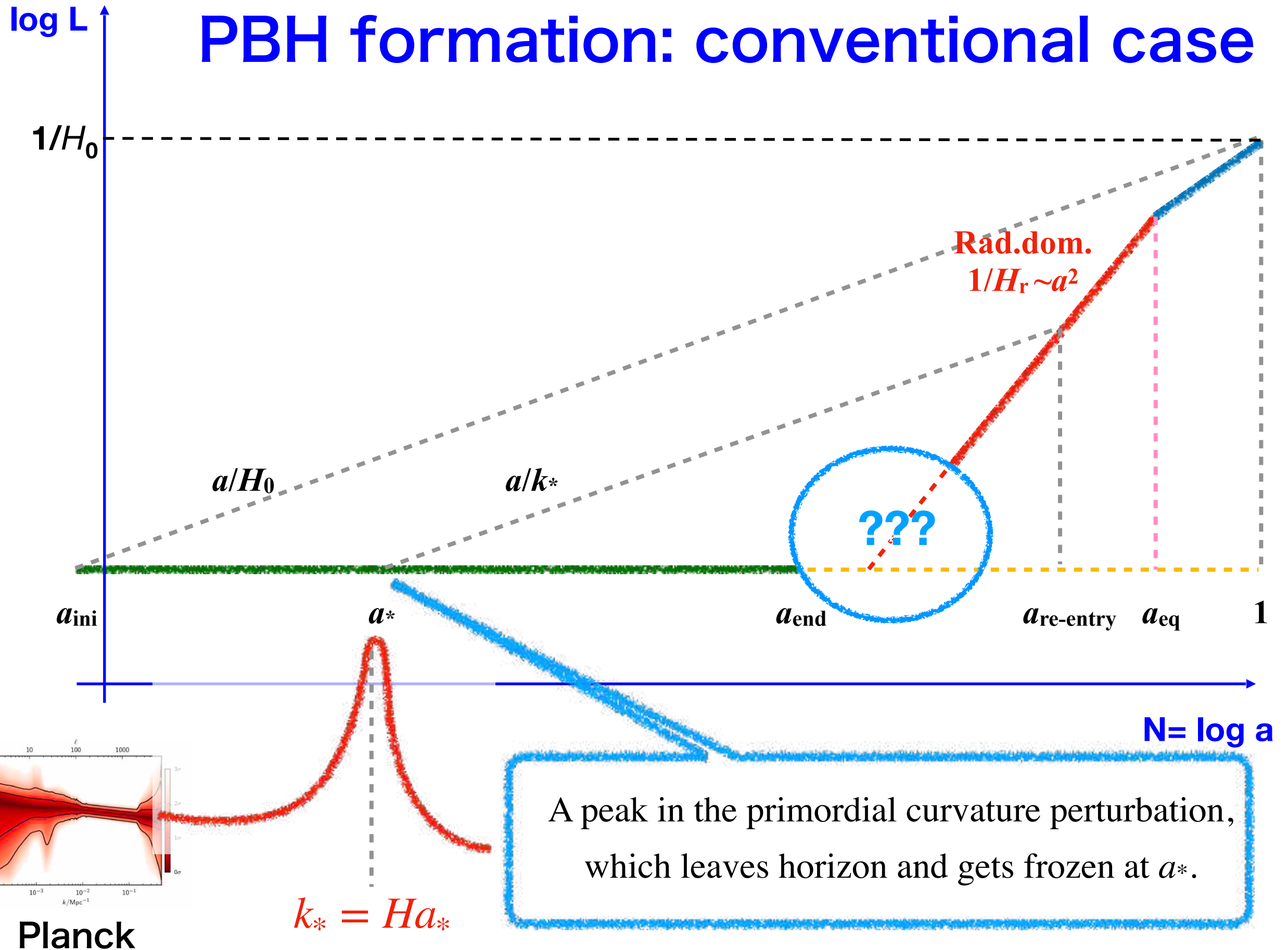
# Introduction

curvature perturbation,  
formation of PBHs,  
and gravitational waves

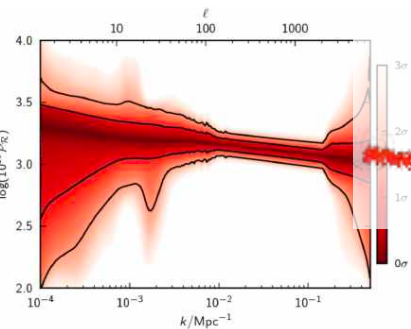
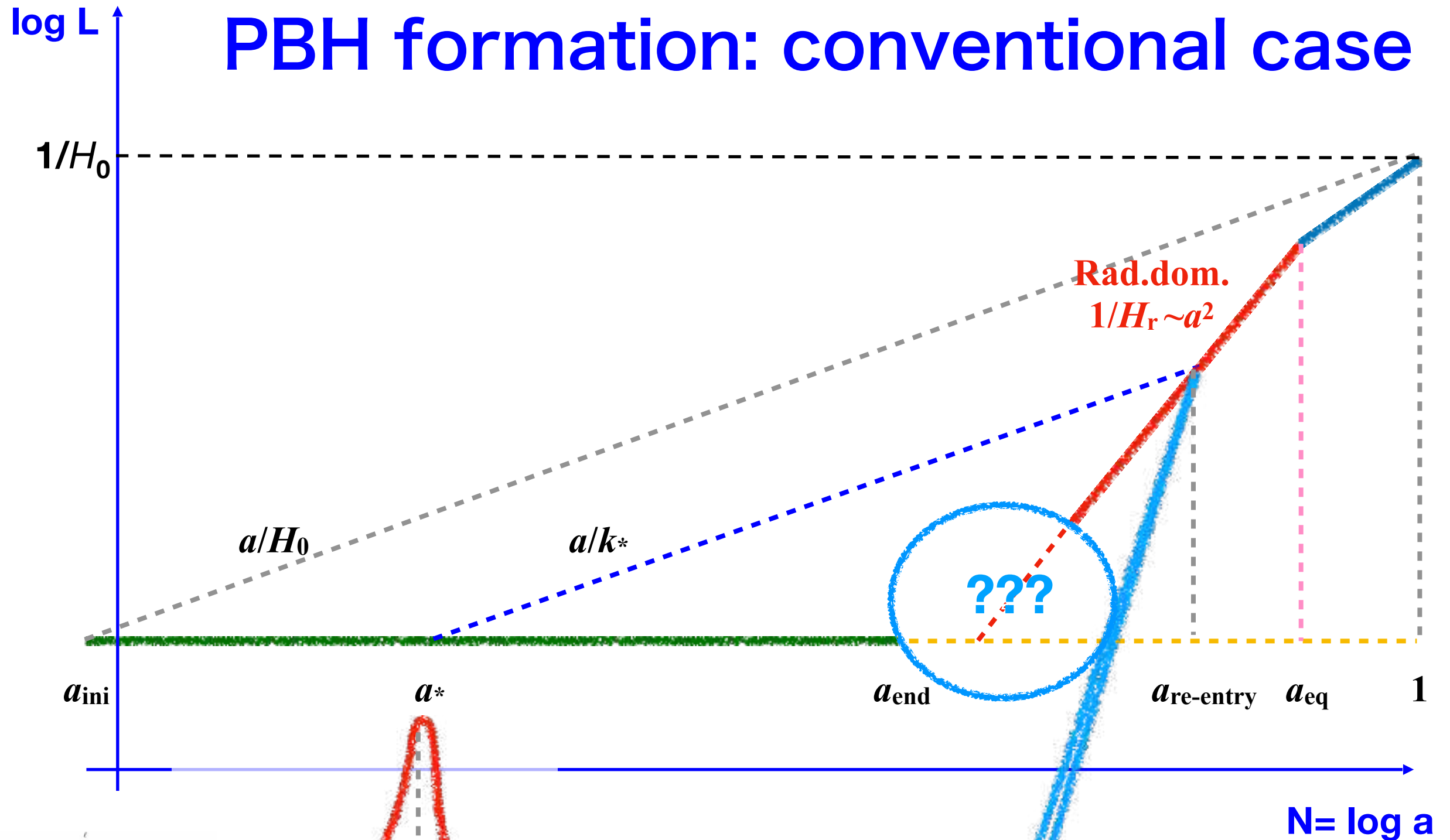
# cosmic spacetime diagram



# PBH formation: conventional case



# PBH formation: conventional case



$k_* = Ha_*$

The peak re-enters horizon during radiation era.  
If the amplitude  $> O(0.1)$ , PBH will form.

# Curvature perturbation to PBH

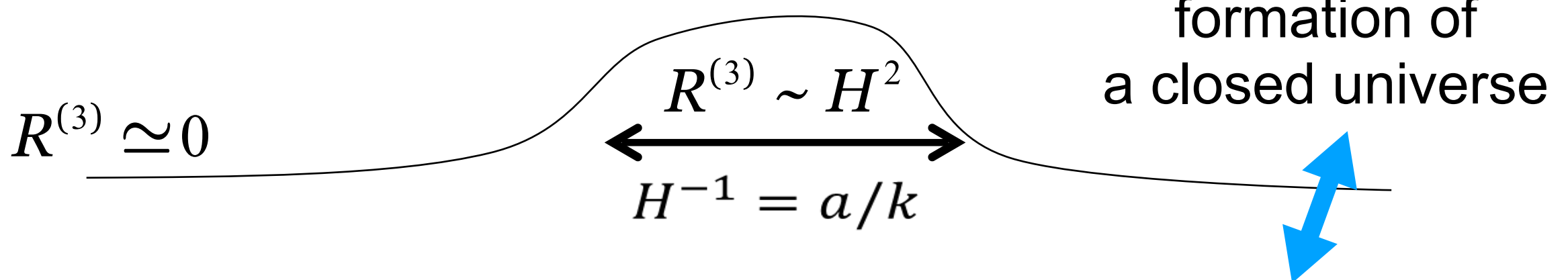
**conventional (PBH formation at rad-dominance) case**

➤ gradient expansion/separate universe approach

$$6H^2(t, \mathbf{x}) + R^{(3)}(t, \mathbf{x}) = 16\pi G\rho(t, \mathbf{x}) + \dots$$

Hamiltonian constraint  
(Friedmann eq.)

$$\Rightarrow R^{(3)} \approx -\frac{4}{a^2} \nabla^2 \mathcal{R}_c \approx \frac{8\pi G}{3} \delta\rho_c \Rightarrow \frac{\delta\rho_c}{\rho} \approx \mathcal{R}_c \text{ at } \frac{k^2}{a^2} = H^2$$



➤ If  $R^{(3)} \sim H^2$  ( $\Leftrightarrow \delta\rho_c / \rho \sim 1$ ), it collapses to form BH

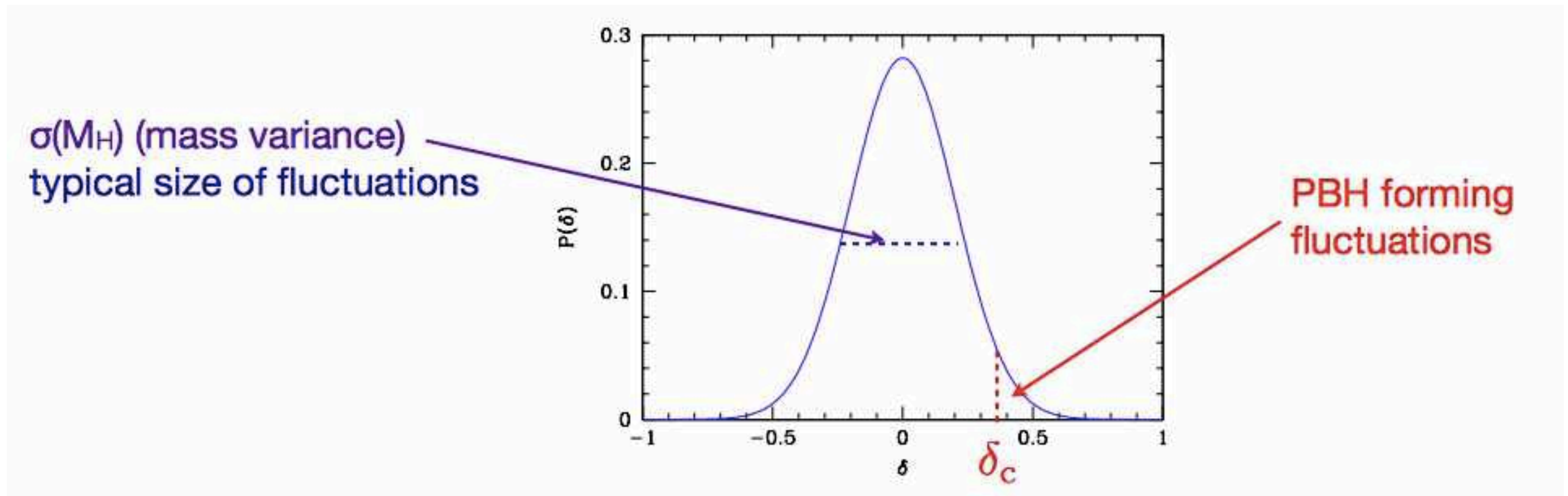
Young, Byrnes & MS '14

➤ Spins of PBHs are expected to be very small

De Luca et al. 1903.01179, ...

# fraction $\beta$ that turns into PBHs

for **Gaussian** probability distribution



- When  $\sigma_M \ll \delta_c$ ,  $\beta$  can be approximated by exponential:

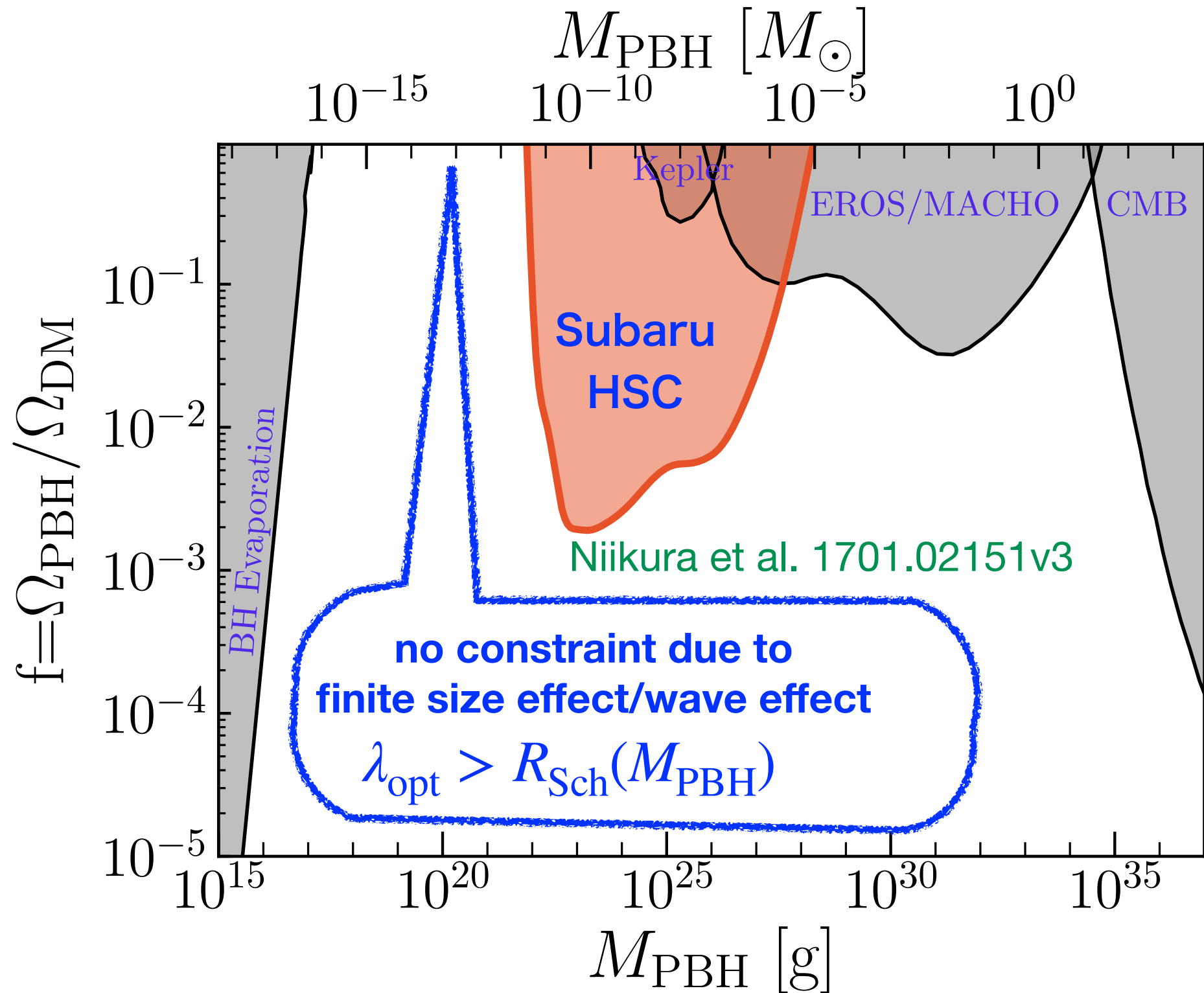
$$\beta \approx \sqrt{\frac{2}{\pi}} \frac{\sigma_M}{\delta_c} \exp\left(-\frac{\delta_c^2}{2\sigma_M^2}\right)$$

$$\delta_c \equiv \left(\frac{\delta\rho_c}{\rho}\right)_{\text{crit}} \sim 0.4$$

Carr '75, ...



# PBH constraints



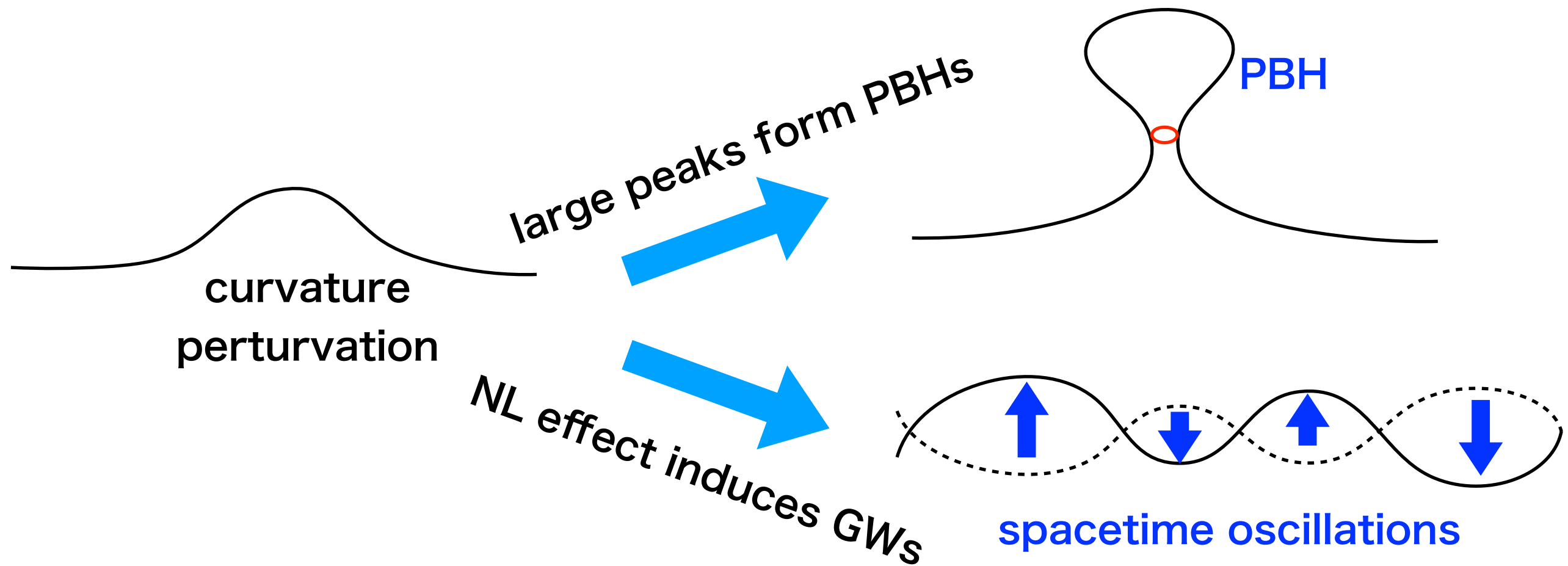
big window at  $M_{\text{PBH}} \approx 10^{17} - 10^{22} \text{ g}$



$T_{\text{re-entry}} \sim 10^4 - 10^8 \text{ GeV}$



# GWs can capture PBHs!



PBHs = CDM with  $M_{\text{PBH}} \sim 10^{21} \text{g}$   
generates GWs with  $f \sim 10^{-3} \text{Hz}$

Background GWs  
at LISA band

LIGO-Virgo : 10 - 1000 Hz

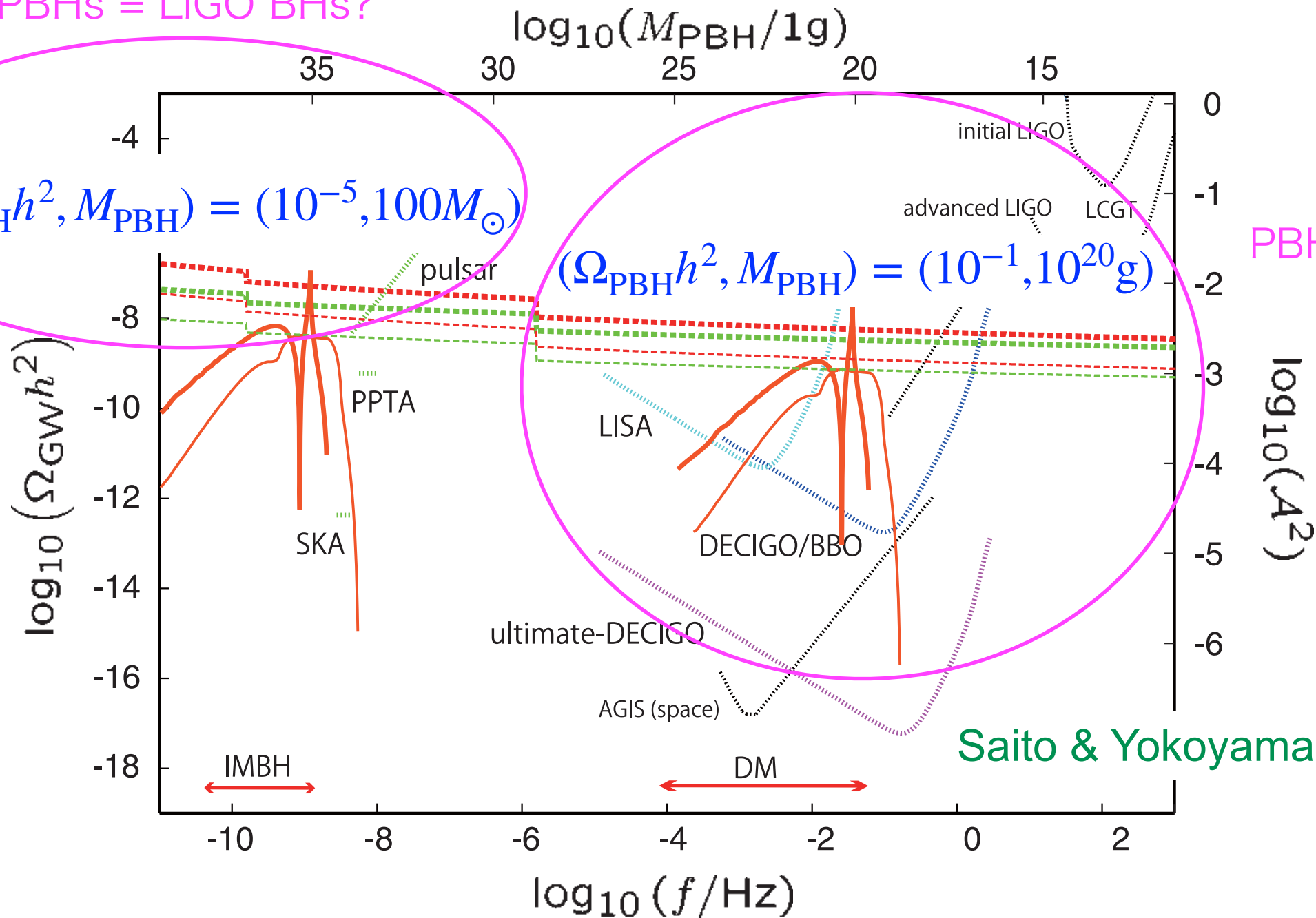
# GWs can test PBH scenario!

PBHs = LIGO BHs?

$(\Omega_{\text{PBH}} h^2, M_{\text{PBH}}) = (10^{-5}, 100 M_{\odot})$

$(\Omega_{\text{PBH}} h^2, M_{\text{PBH}}) = (10^{-1}, 10^{20} \text{g})$

PBHs = CDM?



Saito & Yokoyama 0812.4339

$$M_{\text{PBH}} \sim 0.1 M_{\odot} \left( \frac{1 \text{GeV}}{T} \right)^2 \sim 10 M_{\odot} \left( \frac{1 \text{pc}^{-1}}{k} \right)^2$$

So far, we have focused on PBH formation  
from primordially **adiabatic** perturbation

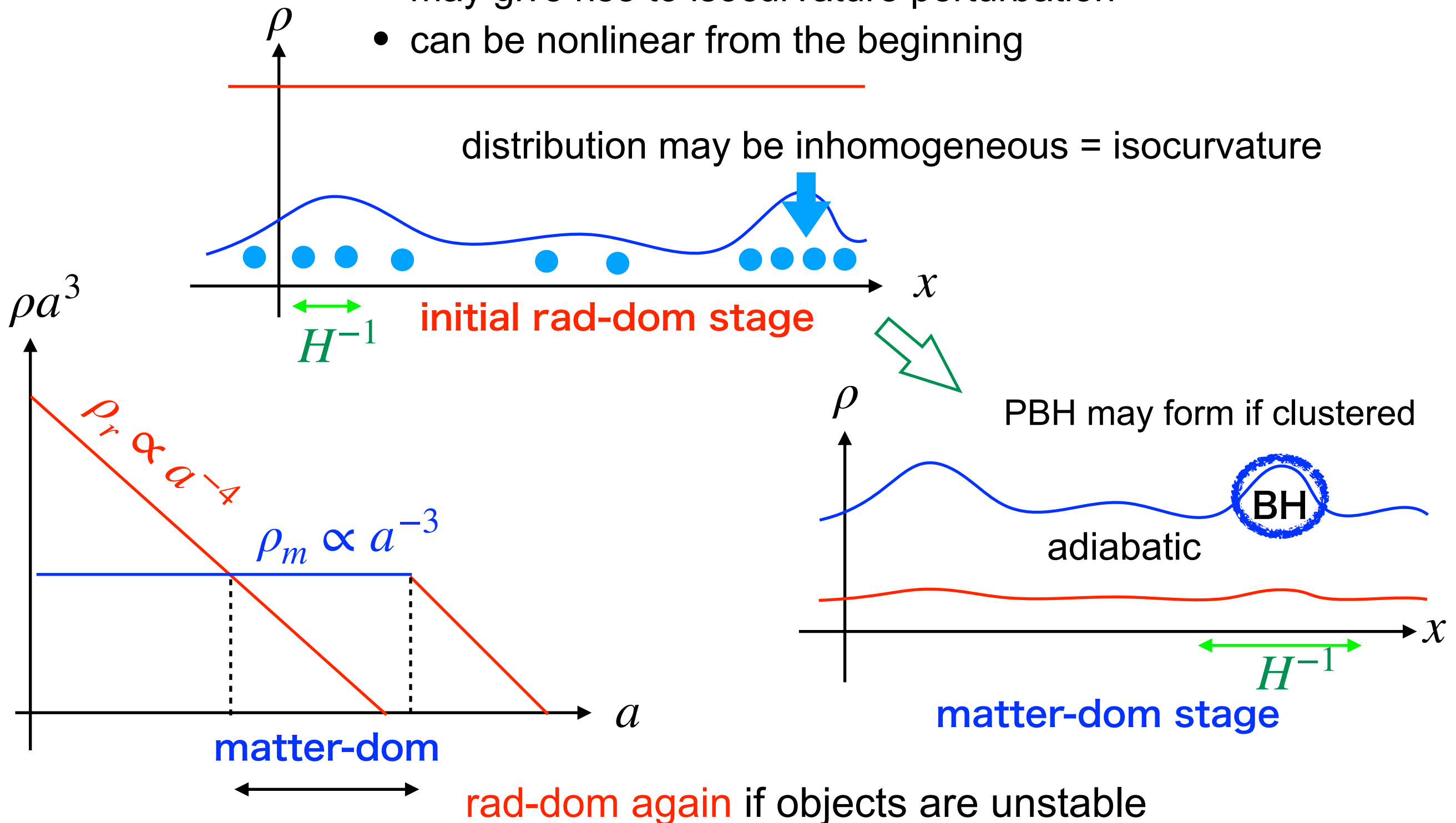
How about primordially  
**isocurvature** perturbation?

# PBHs from Isocurvature Perturbation

eg, E. Cotner, A. Kusenko, MS & V. Takhistov, 1907.10613

non-gravitational formation of compact objects/Q-balls/etc inside horizon.

- may give rise to isocurvature perturbation
- can be nonlinear from the beginning



# linear theory

H. Kodama & MS, IJMPA 1 (1986) 265, ibid 2 (1987) 491

matter isocurvature perturbation

$$S \equiv \delta_m - \frac{3}{4}\delta_r \rightarrow \delta_m \text{ at } a \rightarrow 0 \text{ (on, say, uniform total density slices)}$$

evolution for  $\omega \ll 1$      $\omega \equiv \left(\frac{k}{Ha}\right)_{eq}$ ,     $R \equiv \frac{a}{a_{eq}}$     modes that are superhorizon at equality

$R \ll 1$  (rad dom)

$$\begin{cases} \mathcal{R}_c = \frac{R}{4} S & \left( \Phi = \frac{R}{8} S \right) \\ \delta = \frac{1}{6} \omega^2 R^3 S \end{cases}$$

$1 \ll R$  (matter dom)

$$\begin{cases} \mathcal{R}_c = \frac{1}{3} S & \left( \Phi = \frac{1}{5} S \right) \\ \delta = \frac{4}{15} \omega^2 R S \end{cases}$$

$\mathcal{R}_c$  : curv pert on comoving slice

$\Phi$  : curv pert on Newton slice

horizon crossing:  $\omega^2 R = \frac{1}{2}$

BH formation criterion:  $\delta(k = aH) = \frac{2}{15} S > \delta_{cr} (\sim 0.5) ?$

# linear theory continued

evolution for  $\omega \gg 1$  (modes that enters horizon before equality)

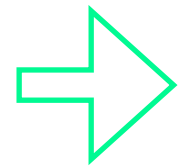
$$\omega \equiv \left( \frac{k}{Ha} \right)_{eq}, \quad R \equiv a/a_{eq}$$

$$R \ll \omega^{-1} \text{ (rad dom)}$$

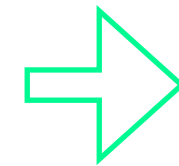
$$\omega^{-1} \ll R \ll 1 \text{ (rad dom)}$$

$$1 \ll R \text{ (matter dom)}$$

$$\begin{cases} \mathcal{R}_c = \frac{R}{4} S \\ \delta = \frac{1}{6} \omega^2 R^3 S \end{cases}$$



$$\begin{cases} \mathcal{R}_c = \frac{3}{4\omega^2 R} S \\ \delta = RS \end{cases}$$



$$\begin{cases} \mathcal{R}_c = \frac{5}{4\omega^2} S \\ \delta = \frac{3R}{2} S \\ \left( \Phi = \frac{3}{4\omega^2} S \right) \end{cases}$$

horizon crossing:  $\omega R = 1/2$

$$\delta(k = aH) = \frac{1}{2\omega} S, \quad \mathcal{R}_c = \frac{3}{2\omega} S$$

$\Phi = O(1)$  implies  $S = O(\omega^2) \gg 1 !!$



highly nonlinear initial condition

conventional growth rate at matter-dom stage

PBH formation criterion?

need more studies!

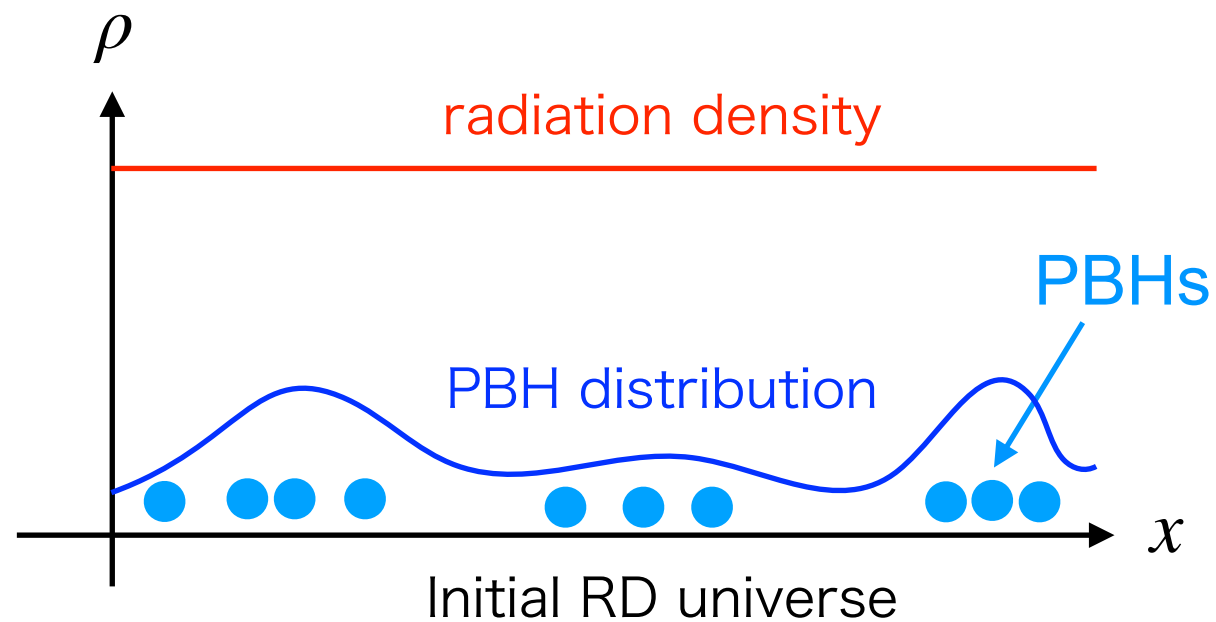
# Isocurvature Perturbation due to inhomogeneous PBH distribution



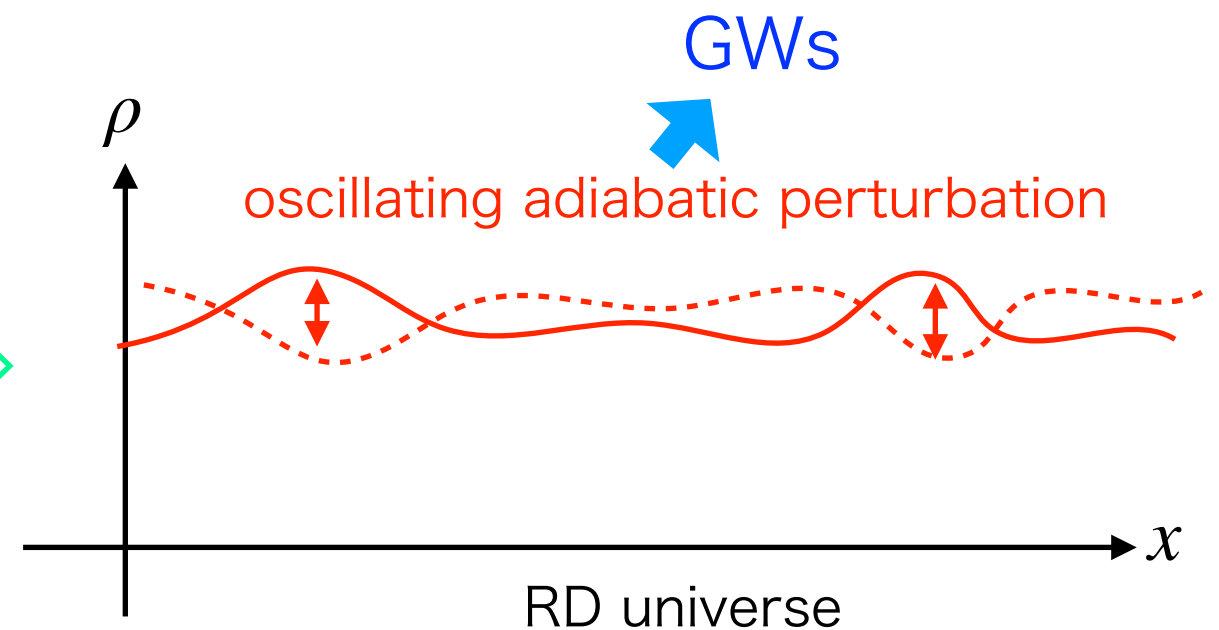
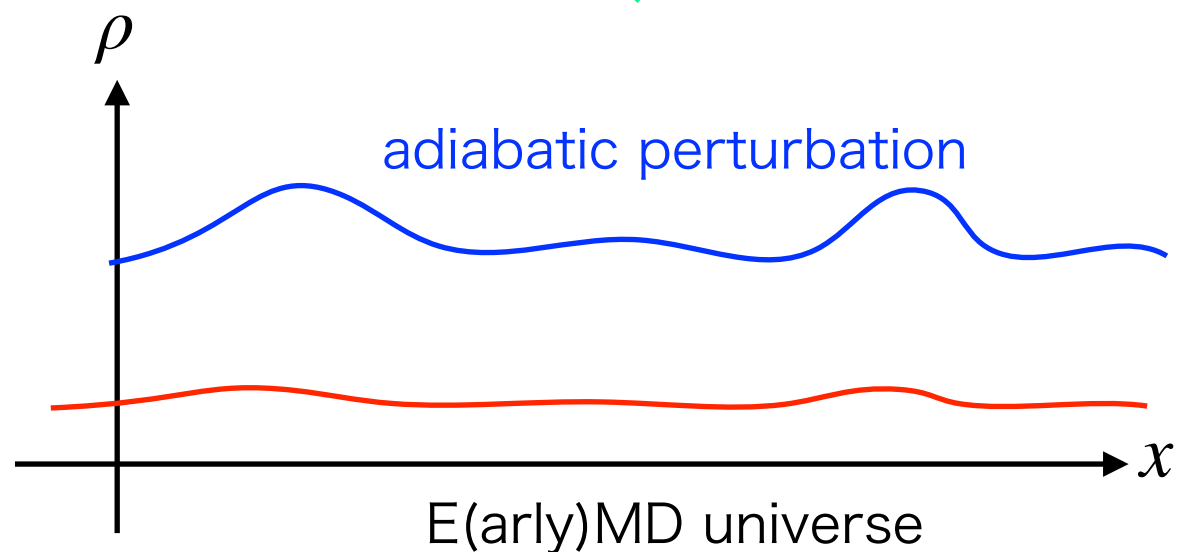
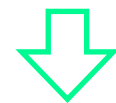
# What if formed objects are PBHs?

Papanikolaou et al., arXiv:2010.11573

Domenech, Lin & MS, arXiv:2012.08151



Even if PBHs are unclustered, randomly distributed, the inhomogeneities may induce GWs when the universe is reheated by PBH evaporation



# Induced GWs from PBH evaporation

Domenech, Lin & MS, arXiv:2012.0851

- If the transition from EMD to RD is slow ( $\Delta t \sim H^{-1}$ ) as in the case of decaying particles, there will be **no significant production** of induced GWs.

Inomata et al., arXiv:1904.12878

$$Q = Q_0 e^{-\Gamma t} \quad \rightarrow \quad \frac{1}{\Delta t} = \frac{1}{Q} \frac{dQ}{dt} = -\Gamma = \text{const.}$$

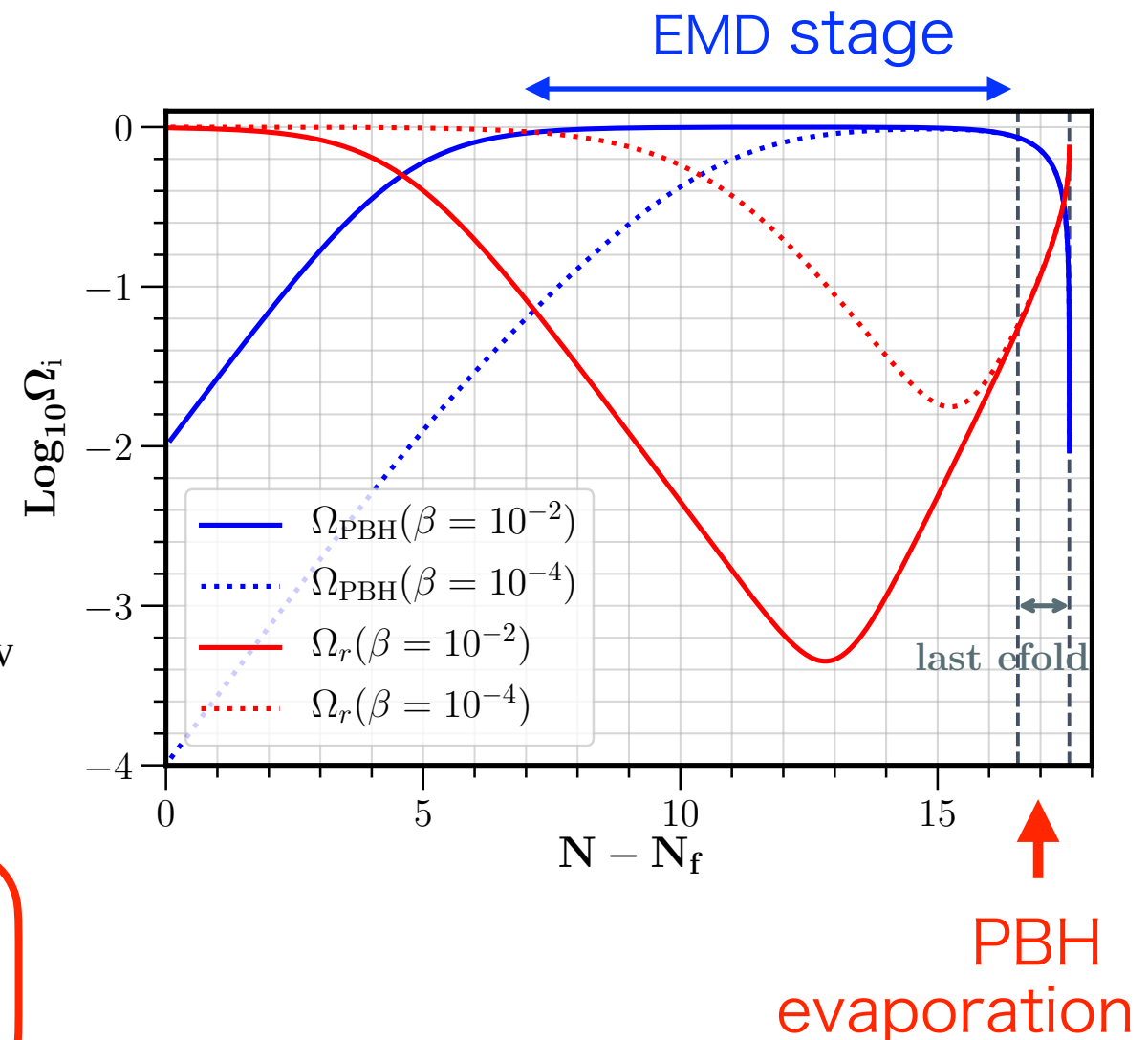
- A **fast transition** leads to **strong enhancement of induced GWs** on sub-horizon scales, which is the case for **PBH evaporation**.

Inomata et al., arXiv: 2003.10455

$$\frac{1}{\Delta t} = \left| \frac{1}{M} \frac{dM}{dt} \right| = \frac{1}{3(t_{\text{ev}} - t)} \gg H \text{ as } t \rightarrow t_{\text{ev}}$$



may lead to strong constraints on early PBH dominance model



# Constraints on early PBH dominated universe

Domenech, Lin & MS, arXiv:2012.08151

Domenech, Takhistov & MS, arXiv:2105.06816

- Assumptions

- Monochromatic mass function for PBHs.

- Poisson distribution for  $\delta n_{\text{PBH}}/n_{\text{PBH}}$  :  $\mathcal{P}_s(k) = \frac{2}{3\pi} (k/k_{\text{UV}})^3$ ;  $k < k_{\text{UV}} = n_{\text{PBH}}^{-1/3}$

- Resulting spectrum

- sharp rise  $\sim k^5$  near the peak.

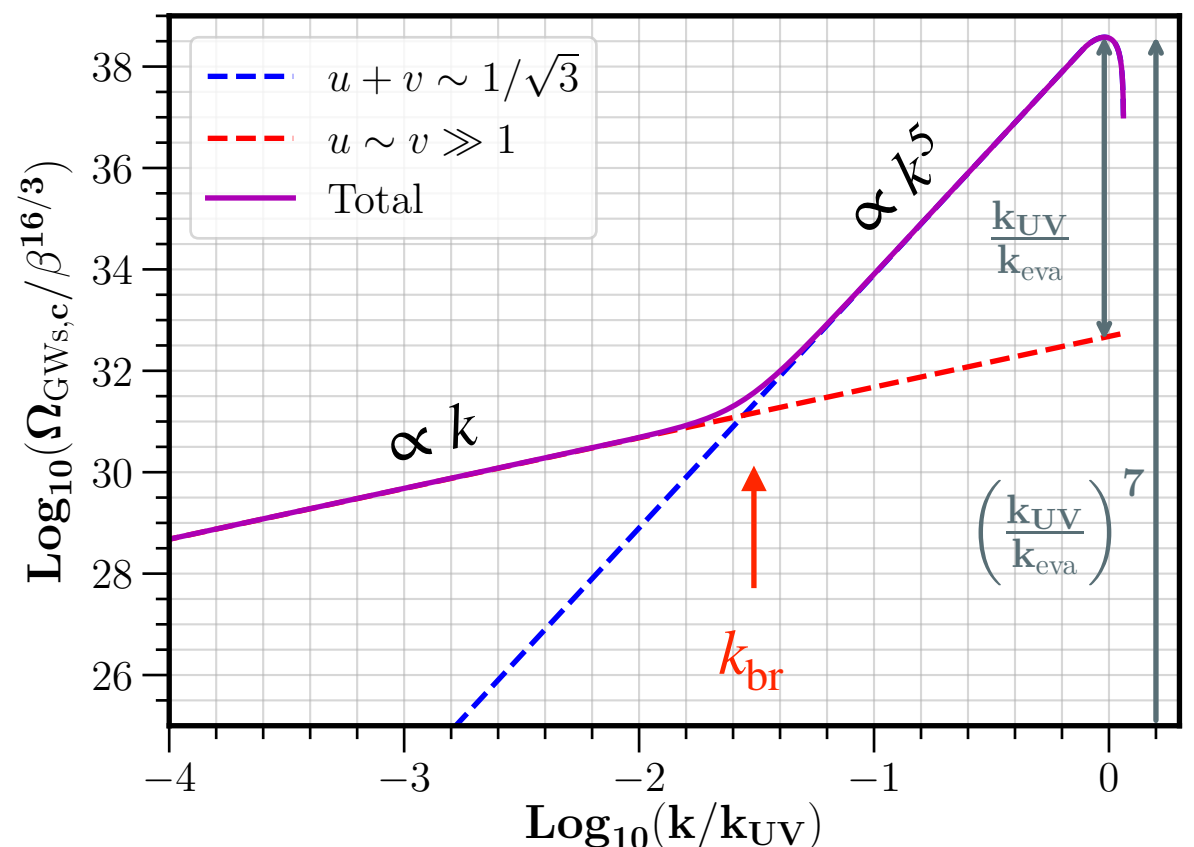
- Peak value:

$$\left( \frac{\Omega_{\text{GW,max}}}{\Omega_{r,0}} \right) \approx 5 \times 10^{34} \beta^{16/3} \left( \frac{M}{10^4 \text{ g}} \right)^{14/3}$$

$\beta$  : PBH fraction at formation

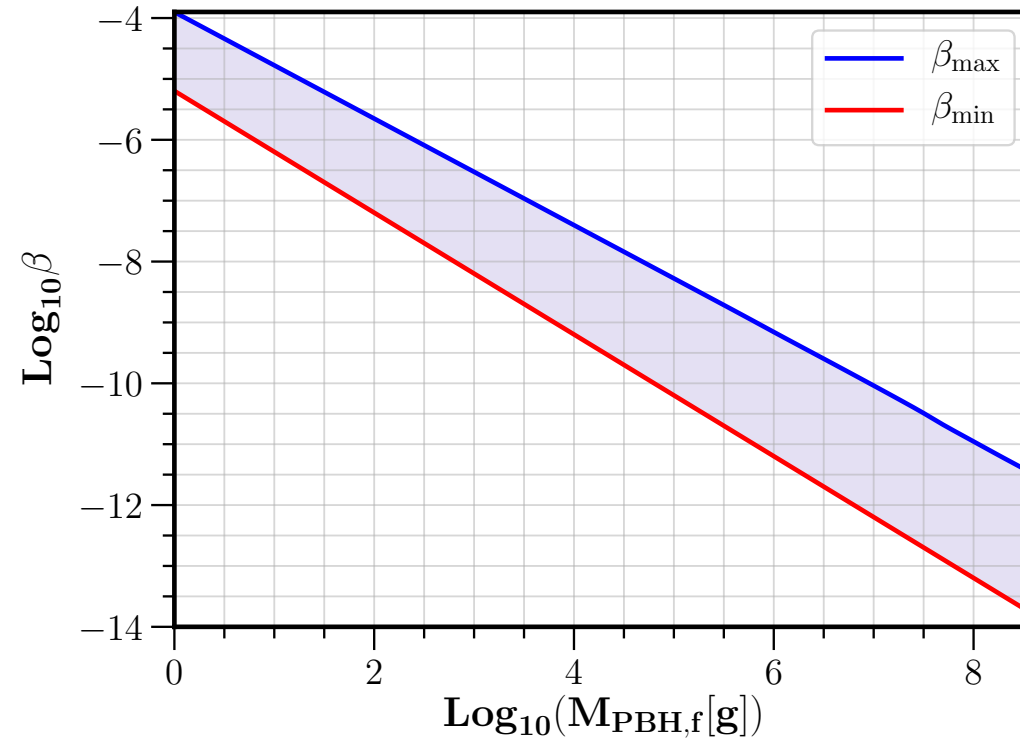


constraints on  $\beta$



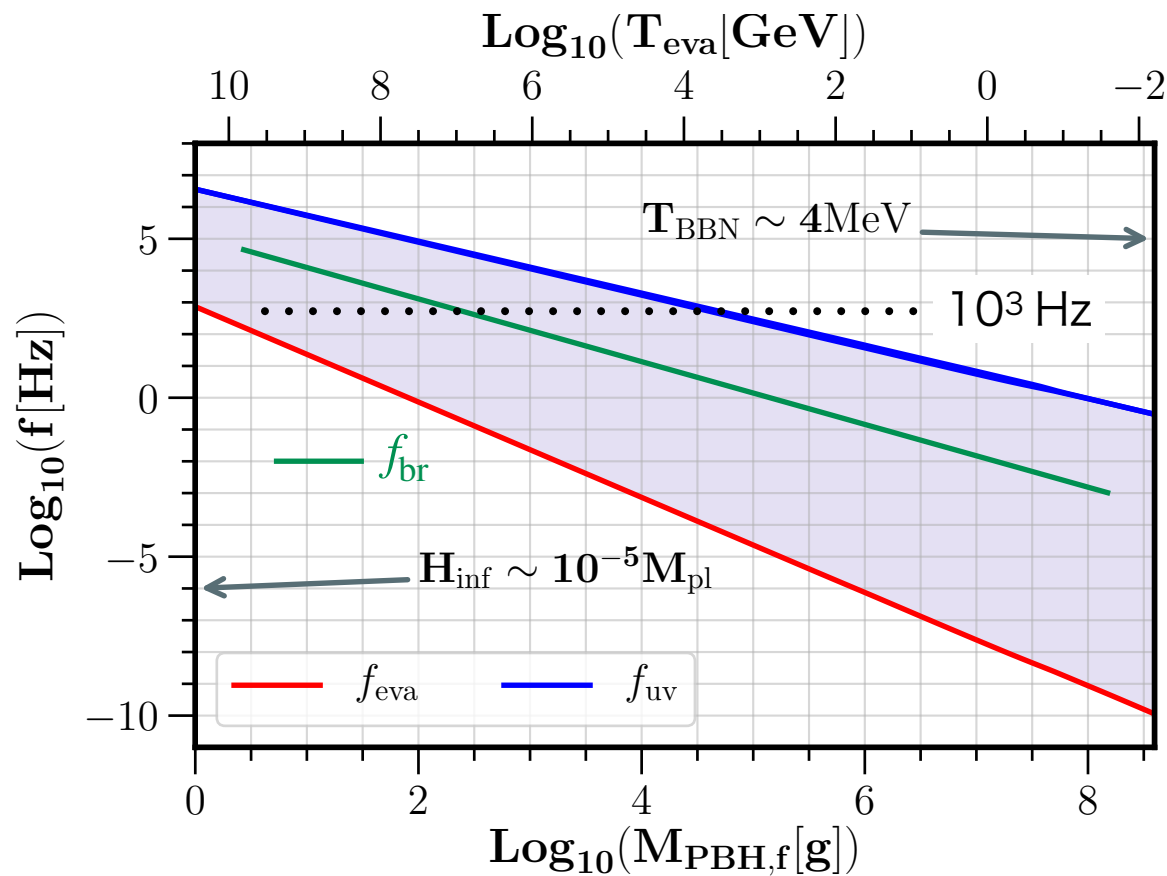
$$k_{\text{br}} \approx 0.04 k_{\text{UV}} \left( M_{\text{PBH}} / 10^4 \text{ g} \right)^{-1/6}$$

# Constraints on $\beta$ and frequencies

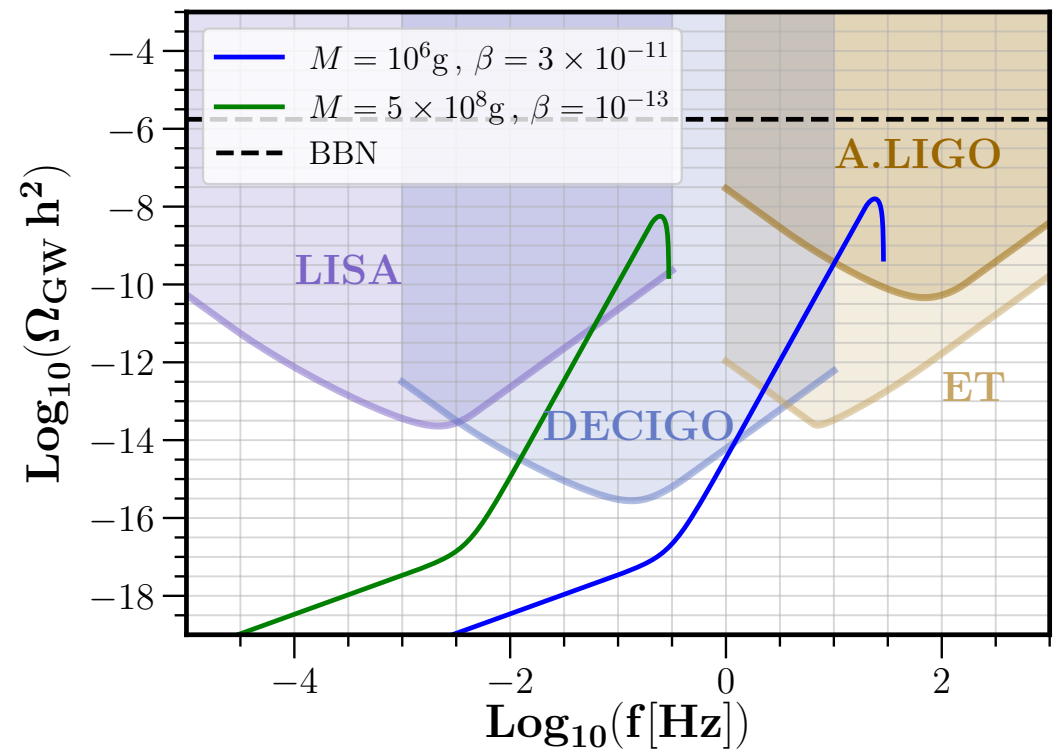


$$\beta_{\text{max}} \approx 3 \times 10^{-8} \left( \frac{M_{\text{PBH}}}{10^5 \text{ g}} \right)^{-7/8}$$

$$\beta_{\text{min}} \approx 6 \times 10^{-10} \left( \frac{M_{\text{PBH}}}{10^5 \text{ g}} \right)^{-1}$$



frequency range vs  $M_{\text{PBH}}$



GW detectors sensitivity curves

## Caviat . . .

For the primordial isocurvature perturbation,

$$\mathcal{P}_S(k) = \frac{2}{3\pi} (k/k_{\text{UV}})^3; \quad k < k_{\text{UV}} = n_{\text{PBH}}^{-1/3}$$

the resulting curvature perturbation at PBH dominated Universe is

$$\Phi = \frac{3}{4} \left( \frac{k_{\text{eq}}}{k} \right)^2 S \sim 0.3 \left( \frac{k_{\text{eq}}}{k_{\text{UV}}} \right)^2 \left( \frac{k}{k_{\text{UV}}} \right)^{-1/2} \quad \text{for } k_{\text{eq}} < k < k_{\text{UV}}$$

➔ The density perturbation becomes **nonlinear for  $k > k_{\text{NL}}$** :

$$\frac{\delta\rho}{\rho} = \frac{2}{3} \left( \frac{k}{aH} \right)^2 \Phi \sim 0.1 \left( \frac{a_{\text{evap}}}{a_{\text{eq}}} \right) \left( \frac{k}{k_{\text{UV}}} \right)^{3/2} \gtrsim 1$$

for  $k > k_{\text{NL}} \sim 5 \left( \frac{a_{\text{evap}}}{a_{\text{eq}}} \right)^{-2/3} k_{\text{UV}}$

$$\log \left( \frac{a_{\text{evap}}}{a_{\text{eq}}} \right)^{2/3} \approx 2 + \frac{8}{9} \left( \log \frac{\beta}{10^{-7}} + \log \frac{M}{10^4 \text{ g}} \right) \quad \uparrow$$

# take-home messages:

- **PBHs** may play central roles in **GW** cosmology



**PBH-GW Cosmology!**

- **(nonlinear) isocurvature** perturbations may play important roles in **PBH-GW** cosmology