A CLOCKWORK WIMP

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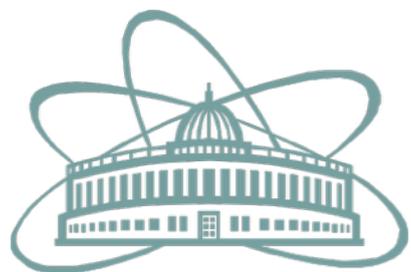
PLAN

The Clockwork Mechanism illustrated

Construction of a Clockwork WIMP

Majorana neutrino masses

Relation to dimension deconstruction

Conclusions

new physics (local) $\longrightarrow$ effective operators $\longrightarrow$ new scale

Weinberg operator $\longrightarrow$ L breaking scale

Proton decay $\longrightarrow$ GUT scale

Axion $\longrightarrow$ PQ symmetry breaking scale

Quantum gravity $\longrightarrow$ Planck scale

…

but   **MASS ≠ SCALE**

$$G_F^{-1/2} \sim \frac{M_W}{g} \qquad \Lambda_\nu \sim \frac{M_R}{y^2} \qquad M_P \sim \frac{M_s}{g_s}$$

Fermi   Weinberg   Planck

Kaplan & Ratazzi (2015); Choi & Im (2015) ; Giuidice & McCullough (2016)

$$\Lambda \sim \frac{M}{y^x}$$ large scale physics : large mass or tiny coupling?

We are often reluctant to introduce small parameters
(even if natural in the sense of 't Hooft)

The **Clockwork** is a possible mechanism
to generate **small numbers**
out of a theory with O(1) parameters
or **large effective scales** (e.g M_P)
out of
dynamics at **much lower energies (**e.g. TeV)

perhaps experimentally accessible energies…

Grand Goal : Addressing the Hierarchy Problem

Guidice & McCullough (2016)

More prosaically, a framework for model builders….

hierarchy problem   Guidice & McCullough (2016)

low scale invisible axion   Guidice & McCullough (2016); Farina *et al* (2016)

inflation   Kehagias & Riotto (2016)

neutrino physics   Hambye, Teresi & MT (2016); Carena *et al* (2017); Ibarra *et al* (2017)

dark matter   Hambye, Teresi & MT (2016) (this talk)

…

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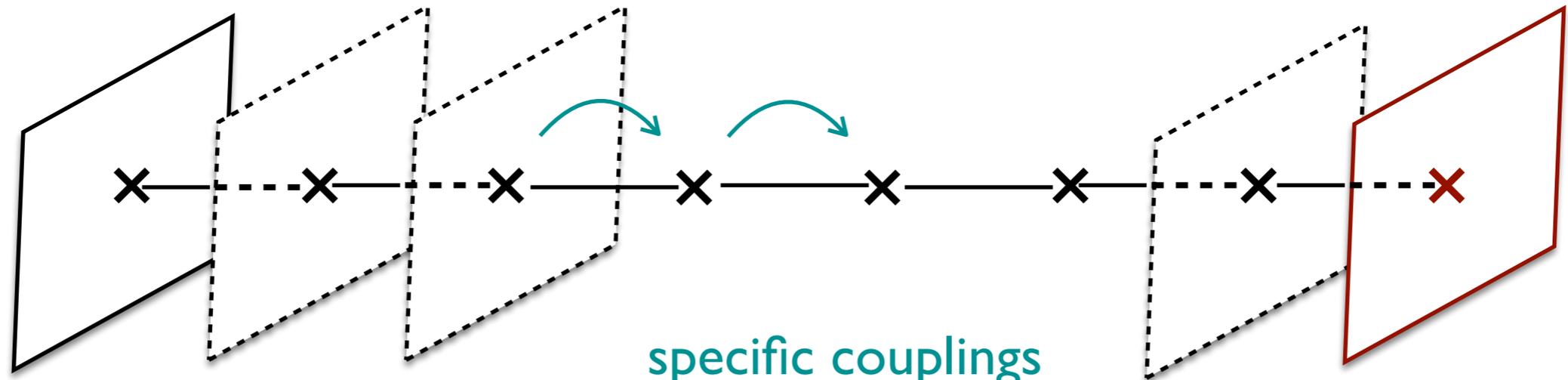
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…

Hidden Sector
(e.g. DM)

specific couplings
lead to exponential
localization of the
lightest hidden field

Visible Sector
(SM)

We first illustrate this with a scalar model
or
Scalar Clockwork

$$\Phi_0 \quad \Phi_1 \quad \cdots \quad \Phi_k \sim e^{i\phi_k/f} \quad \cdots \quad \Phi_N$$

$$\mathcal{L} = \frac{f^2}{2} \sum_{k=0}^{N} |\partial\Phi_k|^2 - \frac{m^2 f^2}{2} \sum_{k=0}^{N-1} (\Phi_k^\dagger \Phi_{k+1}^q + h.c.) \qquad \mathbf{q > 1}$$
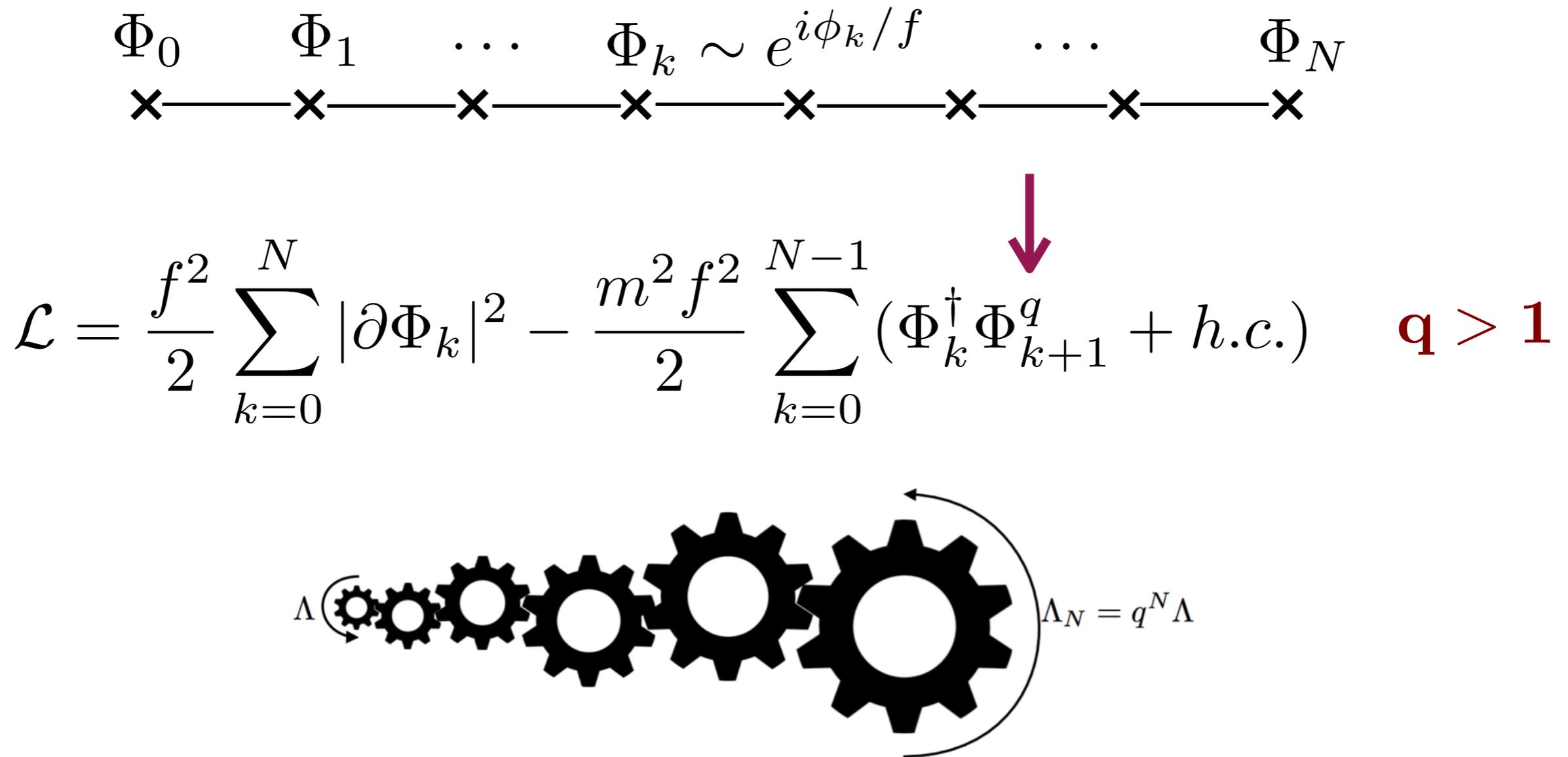
$$\Lambda \qquad \Lambda_N = q^N \Lambda$$

Figure 1: A schematic representation of the clockwork mechanism increasing the interaction scale of a non-renormalisable operator.

Guidice & McCullough (2016)

$$\Phi_0 \quad \Phi_1 \quad \cdots \quad \Phi_k \sim e^{i\phi_k/f} \quad \cdots \quad \Phi_N$$

$$\mathcal{L} = \frac{f^2}{2} \sum_{k=0}^{N} |\partial \Phi_k|^2 - \frac{m^2 f^2}{2} \sum_{k=0}^{N-1} (\Phi_k^\dagger \Phi_{k+1}^q + h.c.) \qquad \mathbf{q > 1}$$
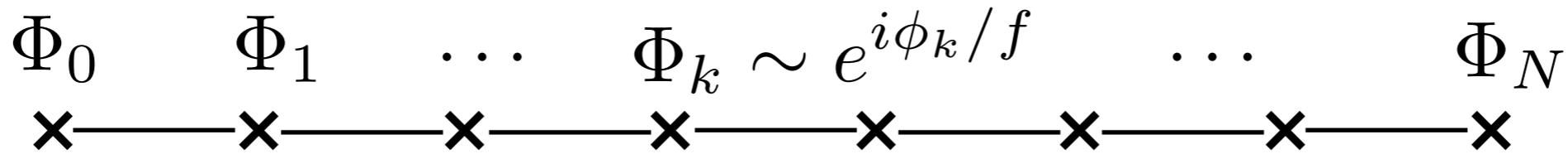
$$-\frac{m^2}{2} \sum_0^{N-1} (\phi_k - q\phi_{k+1})^2 + \ldots$$

$$m^2 \neq 0$$

$$G = U_0(1) \otimes U_1(1) \otimes \ldots \otimes U_N(1) \quad \longrightarrow \quad U(1)$$

one Goldstone mode + N massive modes

$$\Phi_0 \qquad \Phi_1 \qquad \cdots \qquad \Phi_k \sim e^{i\phi_k/f} \qquad \cdots \qquad \Phi_N$$

$$-\frac{m^2}{2}\sum_0^{N-1}(\phi_k - q\phi_{k+1})^2 + \ldots \quad \longrightarrow \quad M^2 = m^2 \begin{pmatrix} 1 & -q & 0 & \ldots & & 0 \\ -q & 1+q^2 & -q & \ldots & & 0 \\ 0 & -q & 1+q^2 & -q & & 0 \\ \vdots & \vdots & \vdots & \ddots & & \vdots \\ & & & -q & 1+q^2 & -q \\ 0 & 0 & 0 & \ldots & -q & q^2 \end{pmatrix}$$

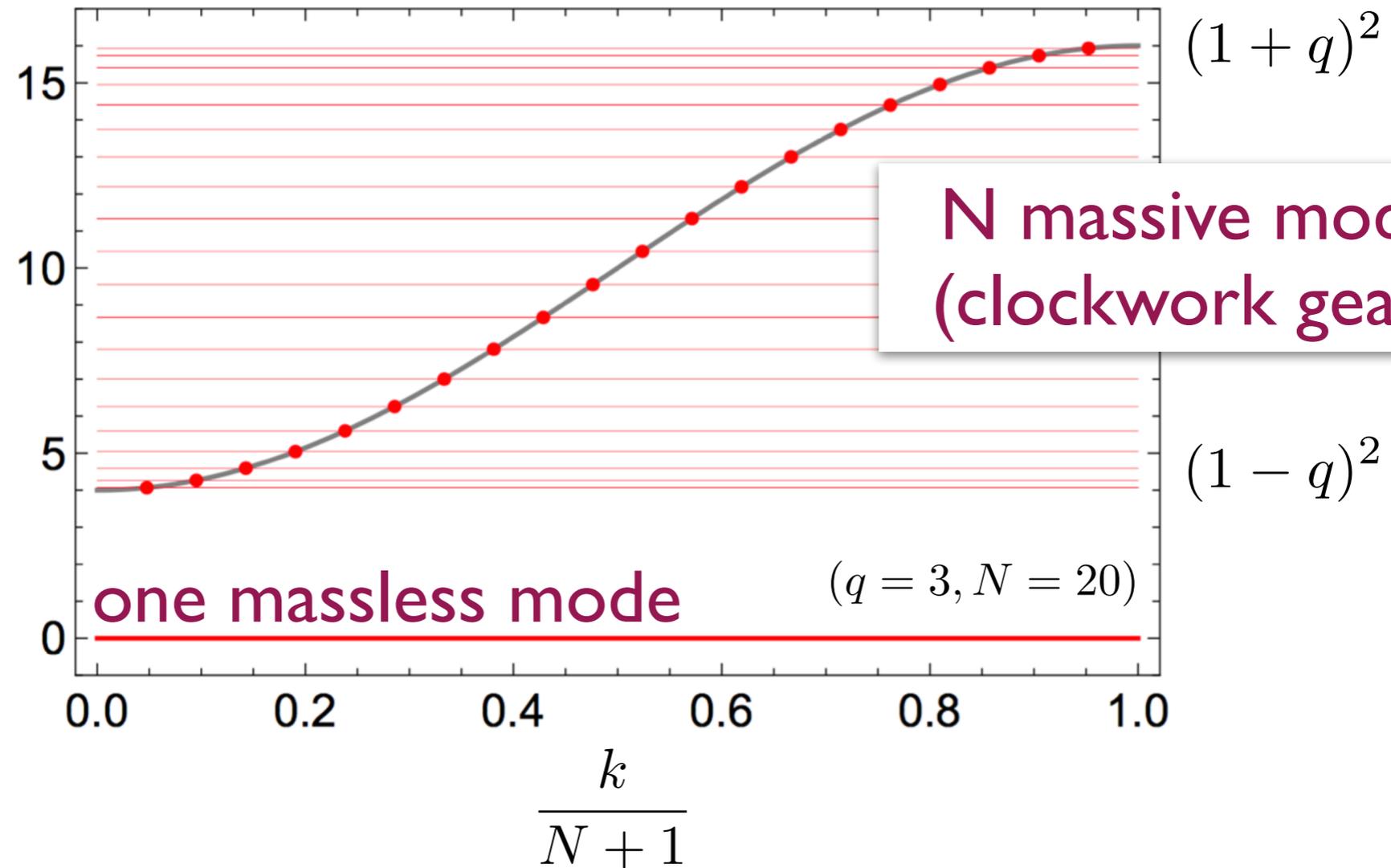$$M^2\varphi^{(0)} = 0 \quad \longrightarrow \quad \varphi^{(0)} \sim \phi_0 + \phi_1/q + \ldots + \boxed{\phi_N/q^N}$$
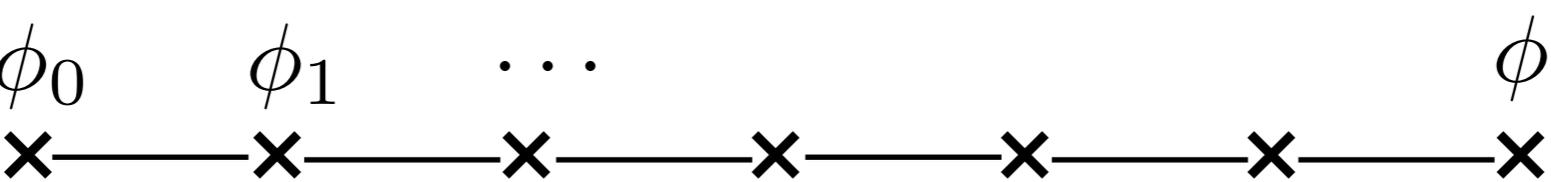
$$q > 1 \rightarrow q^N \gg 1$$

massless mode
localized towards site k=0

$$\frac{m^2_{(k)}}{m^2}$$

$(1+q)^2$

N massive modes
(clockwork gears)

$(1-q)^2$

$(q = 3, N = 20)$

one massless mode

$$\frac{k}{N+1}$$

plot from Farina *et al* (2016)

$$\varphi^{(0)} \sim \phi_0 \quad \phi_1 \quad \cdots \qquad\qquad \phi_N \text{ — SM}$$

massless mode localized towards site k=0

$$\mathcal{L} \supset \frac{\phi_N}{F} G\tilde{G} \quad \longrightarrow \quad \frac{\varphi^{(0)}}{q^N F} G\tilde{G}$$

very tiny coupling!

e.g.

$$N = 15, \quad q = 3 \quad \longrightarrow q^N F \sim 10^{10} \left(\frac{F}{\text{TeV}} \right) \quad \text{GeV}$$

thus large effective scale!

Massive particle

$$\chi$$

Real (~ Majorana neutrino)
or Complex (~ Dirac neutrino)

Neutral

(i.e. not seen, in a broad sense)

Very stable

$$\tau \gtrsim 10^{18} \text{sec}$$ to be around today

$$\tau \gtrsim 10^{26} \text{sec}$$ from constraints
on decay products

easy as Z_2

Annihilate into Standard Model particles

$$\chi + \chi \leftrightarrow f + \bar{f}$$

directly or indirectly
(i.e. secluded dark matter)

$$\langle \sigma v \rangle \sim 3 \cdot 10^{-26} \text{cm}^3 \text{s}^{-1}$$

WIMP **stability** in general protected by a **symmetry**

discrete global ~ R-parity in SUSY SM & hundreds of toy models

discrete ~ SO(10) or **continuous gauge sym.** ~ hidden photons

 accidental symmetry (like the proton)

1. Proton stability (natural because it's not a fundamental particle)

$$\mathcal{L} \supset y\,\bar{p}e^+\pi_0 \quad \longrightarrow \quad \Gamma_p \propto y^2 m_p \quad \longrightarrow \quad y \lesssim 10^{-32}$$

$$\mathcal{O} \sim \frac{1}{\Lambda^2}\bar{\psi}_q\Gamma\psi_q\bar{\psi}_q\Gamma'\psi_e \quad \longrightarrow \quad \Gamma_p \propto \frac{1}{\Lambda^4}m_p^5$$

$$\tau_p \gtrsim 10^{40}\mathrm{sec} \quad \longrightarrow \quad \Lambda \gtrsim 10^{16}\mathrm{GeV} \qquad \Lambda_{\mathrm{gut}}?$$

2. Large SU(2) representation?

e.g. a Majorana SU(2) 5-plet (aka **Minimal Dark Matter**) $M_5 \sim \mathrm{TeV}$

dim-6 decay op. $\tau \sim \dfrac{\Lambda^4}{M^5}$ $\qquad \tau_5 \gtrsim 10^{26}\mathrm{sec} \longrightarrow \Lambda \gtrsim 10^{16}\mathrm{GeV}$

Q: can we make a WIMP long lived
with $\Lambda \sim$ TeV ?

chiral
chain

$$R_0 \quad L_1 R_1 \quad \ldots \quad L_N R_N$$

$$\times \quad \times \quad \times \quad \times \quad \times \quad \times \quad \times$$

complex
scalars
$$\begin{cases} S_i \sim (1, -1) & \text{under} \quad U(1)_{L_{i+1}} \times U(1)_{R_i} \\ C_i \sim (1, -1) & \text{under} \quad U(1)_{L_i} \times U(1)_{R_i} \end{cases}$$
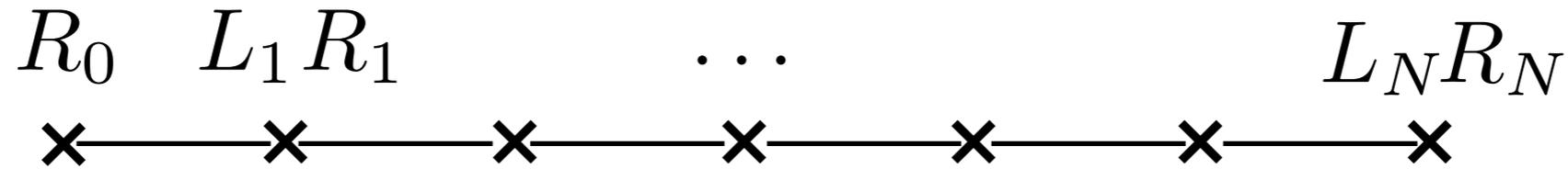
basic ingredients:

One chiral chain (L and R Weyl spinors)
+ 2 N complex scalars (spurions and/or dynamical)

The construction goes through 4 steps

chiral chain

$$R_0 \quad L_1 \, R_1 \qquad \ldots \qquad\qquad L_N R_N$$

$$\times \!\!-\!\!\times \!\!-\!\!\times \!\!-\!\!\times \!\!-\!\!\times \!\!-\!\!\times \!\!-\!\!\times$$

complex scalars

$$\begin{cases} S_i \sim (1,-1) & \text{under} \quad U(1)_{L_{i+1}} \times U(1)_{R_i} \\ C_i \sim (1,-1) & \text{under} \quad U(1)_{L_i} \times U(1)_{R_i} \end{cases}$$

step 1: break the chiral symmetries

$$\begin{cases} m \to y_S \langle S_i \rangle \\ mq \to y_C \langle C_i \rangle \end{cases}$$

These fields are spurions in the original framework **They are dynamical in our case**

$$U(1)_{R_0} \times U(1)_{L_1} \times U(1)_{R_1} \times \ldots \times U(1)_{L_N} \times U(1)_{R_N} \longrightarrow U(1)_R$$

one **massless** mode

$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c.$$

$$R_0 \quad L_1 R_1 \quad \cdots \quad L_N R_N$$
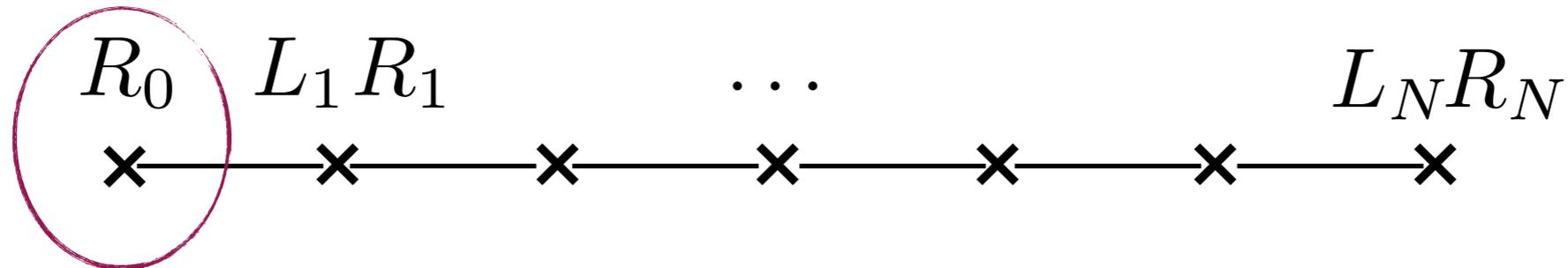
$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c.$$

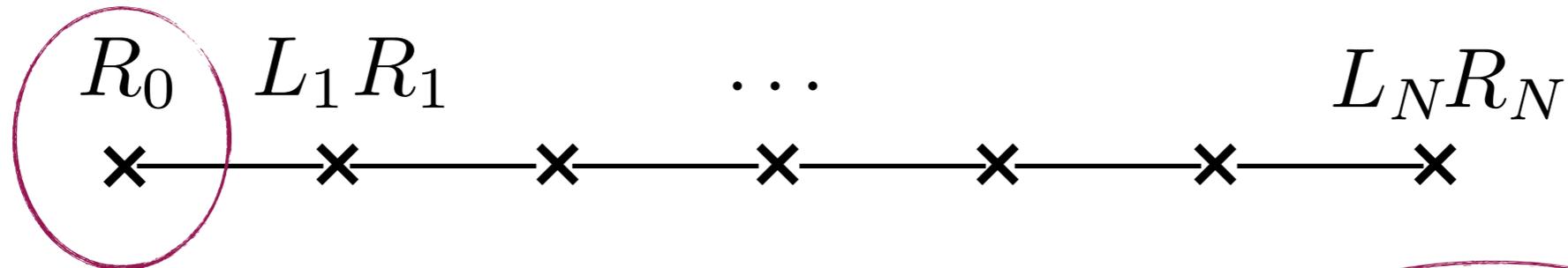$$N \approx R_0 + \frac{1}{q} R_1 + \frac{1}{q^2} R_2 + \ldots + \frac{1}{q^N} R_N$$

massless
chiral mode

$$\psi^{(n)} \approx \frac{1}{\sqrt{N}} \sum_k \left[\mathcal{O}(1) L_k + \mathcal{O}(1) R_k \right]$$

N Dirac gears
mass ~ q m

i.e. pretty much like the scalar clockwork

$$R_0 \quad L_1 R_1 \quad \cdots \quad L_N R_N$$

$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q\bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$

step 2: break the residual chiral symmetry
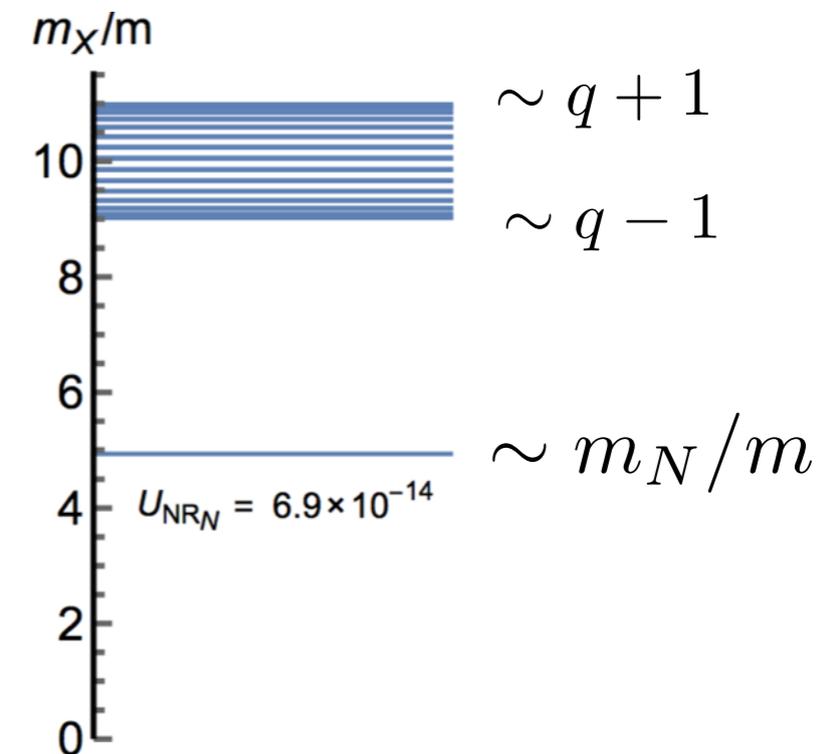
i.e.

give a mass to the N ~ DM state
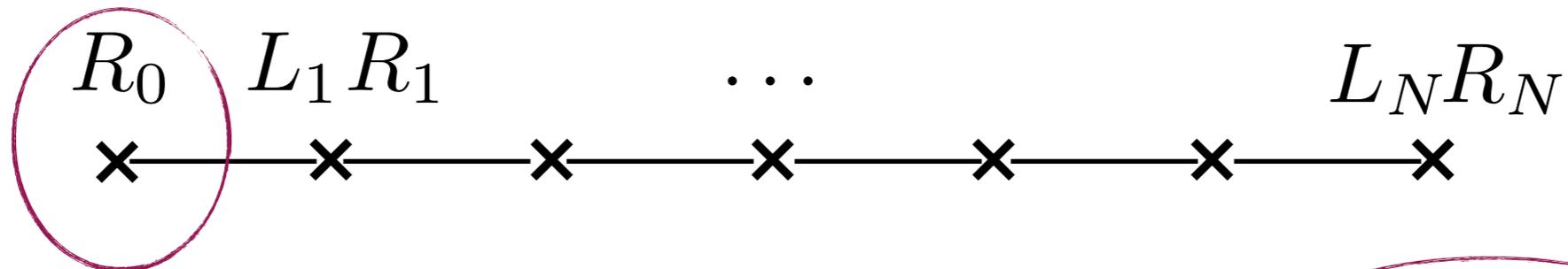
$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$

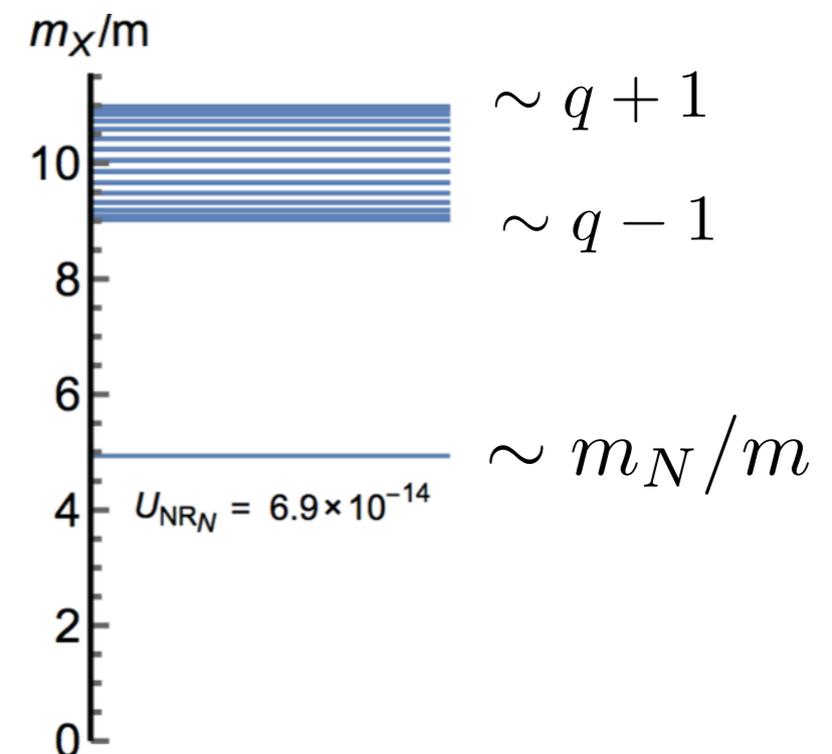$(q = 10, N = 15, m_N = 0)$

$(q = 10, N = 15, m_N = 5m)$

$\sim q + 1$

$\sim q - 1$

$m_N = 0$

$\sim q + 1$

$\sim q - 1$

$\sim m_N / m$

$U_{NR_N} = 6.9 \times 10^{-14}$

$$R_0 \quad L_1 \, R_1 \qquad \ldots \qquad\qquad L_N R_N$$
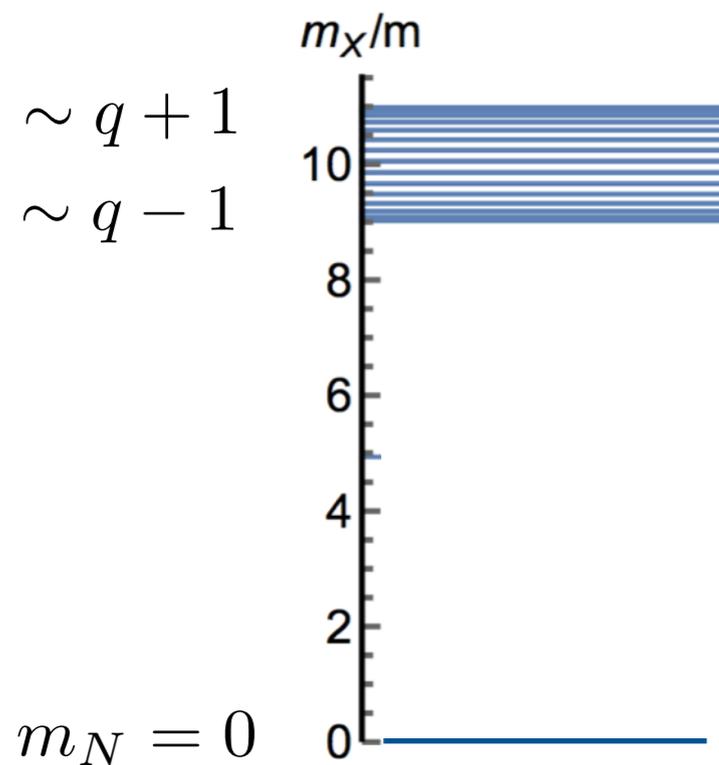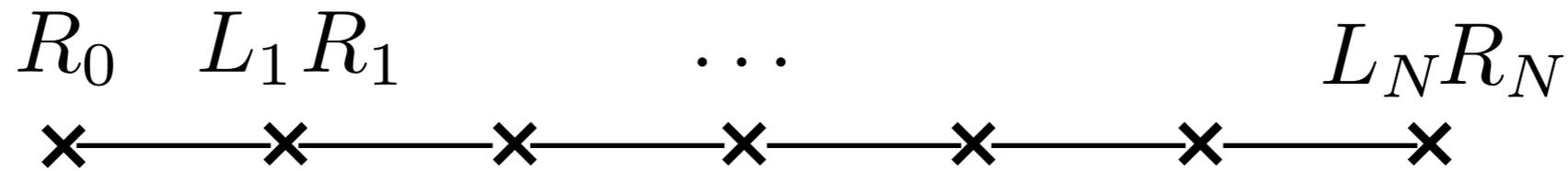
$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$

clockwork mechanism
unaffected provided

$$m_N \lesssim qm$$

$\begin{cases} \text{one light \& localized mode} \\ \text{N gears with O(1) couplings} \end{cases}$

(pseudo-Dirac if $q \gg m_N/m$)

$(q = 10, N = 15, m_N = 5m)$
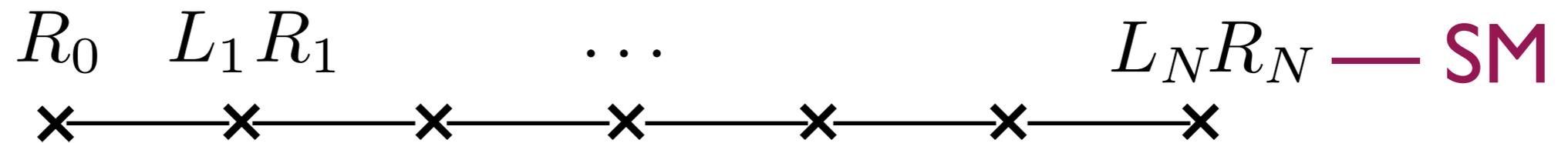
m_X/m

$\sim q+1$

$\sim q-1$

$\sim m_N/m$

$U_{NR_N} = 6.9 \times 10^{-14}$

Hambye, Teresi & MT (2016)

$$R_0 \quad L_1 R_1 \qquad \cdots \qquad\qquad L_N R_N \text{ — SM}$$

$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$
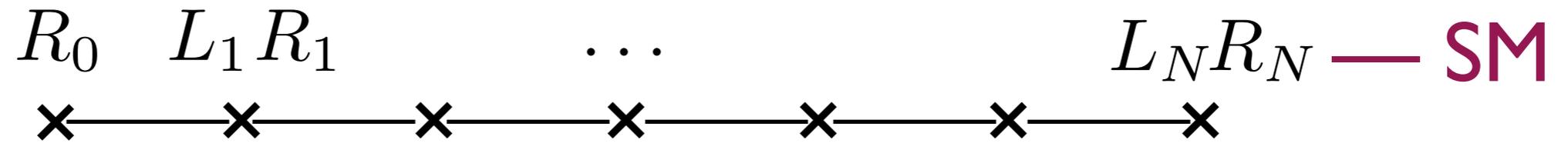
$$-y(\bar{L}_{SM} \tilde{H} R_N + h.c.)$$

step 3: couple the chain to the SM

$$R_0 \quad L_1 R_1 \qquad \ldots \qquad\qquad L_N R_N \text{ —— SM}$$

$$\times \text{———} \times \text{———} \times \text{———} \times \text{———} \times \text{———} \times \text{———} \times$$
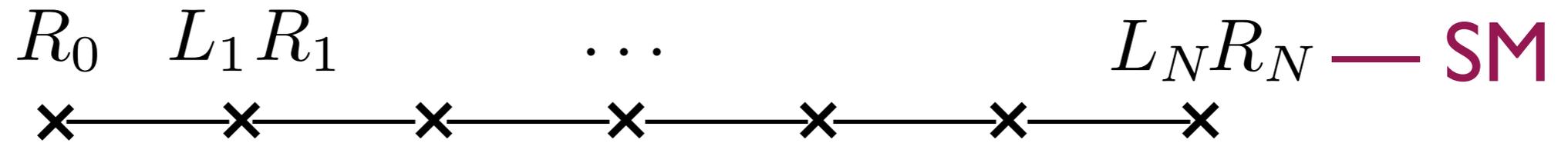
$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$

$$-y(\bar{L}_{SM} \tilde{H} R_N + h.c.)$$

step 3: couple the chain to the SM

$$N \approx R_0 + \frac{1}{q} R_1 + \frac{1}{q^2} R_2 + \ldots + \frac{1}{q^N} R_N \rightarrow \mathcal{L} \supset -\frac{y}{q^N} \bar{L}_{SM} \tilde{H} N + h.c.$$

tiny
coupling

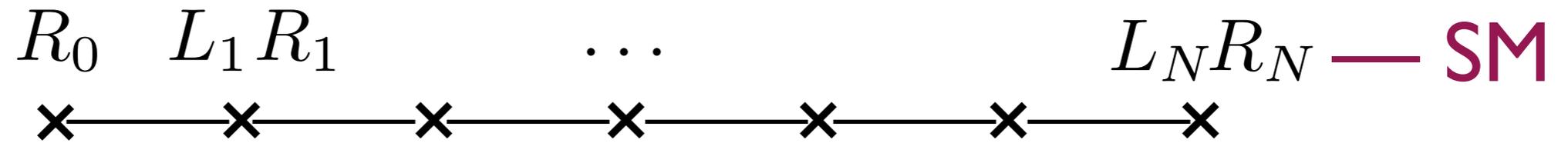$$R_0 \quad L_1 R_1 \qquad \ldots \qquad\qquad L_N R_N \text{ — SM}$$

$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$

$$-y(\bar{L}_{SM} \tilde{H} R_N + h.c.)$$

$$\mathcal{L} \supset \frac{-y}{q^N} \bar{L}_{SM} \tilde{H} N + h.c. \quad \longrightarrow \quad \Gamma(N \to \nu h, \nu Z, lW) \sim \frac{m_N}{8\pi} \frac{y^2}{q^{2N}}$$

The N mode is **unstable** but lifetime $\gtrsim 10^{26} \sec$ (gamma's, etc.)

if $\quad q^{2N} \gtrsim 10^{52} \left(\frac{m_N}{100 \text{GeV}} \right) y^2 \qquad$ **e.g.** $(q = 10, N = 26, y = 1)$

$$R_0 \quad L_1 R_1 \qquad \ldots \qquad\qquad L_N R_N \text{ --- SM}$$

$$\mathcal{L} \supset -m \sum_{i=1}^{N} \left(\bar{L}_i R_{i-1} - q \bar{L}_i R_i \right) + h.c. - \frac{1}{2} m_N \bar{R}_0^c R_0$$

$$-y(\bar{L}_{SM} \tilde{H} R_N + h.c.)$$
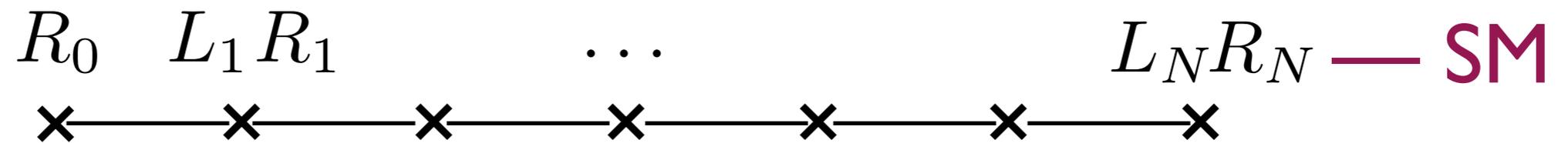
last (but not the least) **step**:
    abundance from thermal freeze-out?

$$NN \rightarrow hh, hZ, l\bar{l}, \ldots \quad \text{but rate} \quad \propto \frac{y^4}{q^{4N}}$$

i.e. very much suppressed…

$$R_0 \quad L_1 R_1 \qquad \cdots \qquad\qquad L_N R_N \; \text{—} \; \text{SM}$$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{kinetic}}$$

$$- \sum_{i=1}^{N} \left(y_S S_i \bar{\tilde{L}}_i R_{i-1} - y_C C_i \bar{\tilde{L}}_i R_i + h.c. \right)$$

$$- \left(y \, \bar{L}_{SM} \widetilde{H} R_N + h.c. \right) - \frac{1}{2} \left(m_N \overline{R_0^c} \, R_0 + h.c. \right)$$

$$- \sum_{i=1}^{N} \left(\lambda_S \, S_i^\dagger S_i \, H^\dagger H + \lambda_C \, C_i^\dagger C_i \, H^\dagger H \right)$$

χSB

N mass
& DM-SM

Higgs
portal

not sure I would put that Lagrangian on a t-shirt

$$\mathcal{L} \supset -\frac{\xi_j^S}{\sqrt{2}} S_1 \bar{\psi}_j P_R N + h.c.$$
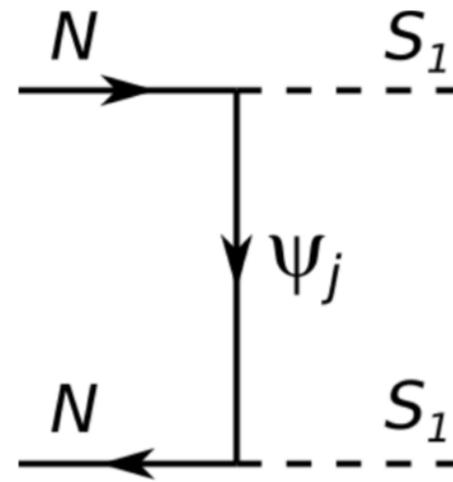
$$-\frac{\xi_j}{\sqrt{2}} h \bar{\psi}_j P_R N + h.c.$$
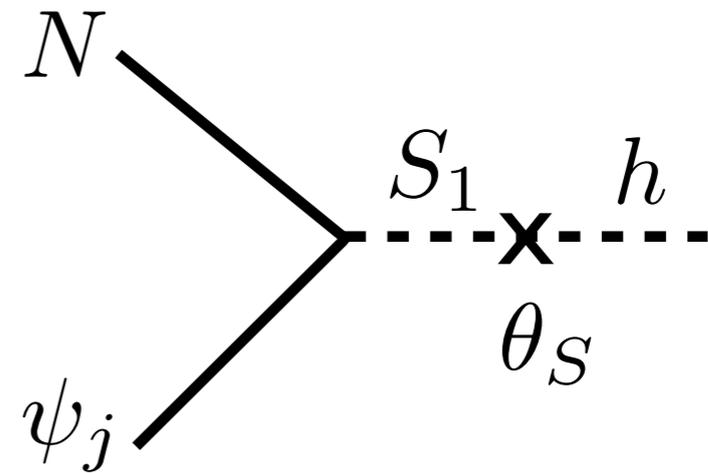
with $\quad \xi_j^S \approx \sqrt{\dfrac{2}{N+1}} \sin\left(\dfrac{j\pi}{N+1}\right) y_S \quad$ **and** $\quad \xi_j \approx \theta_S \xi_j^S$
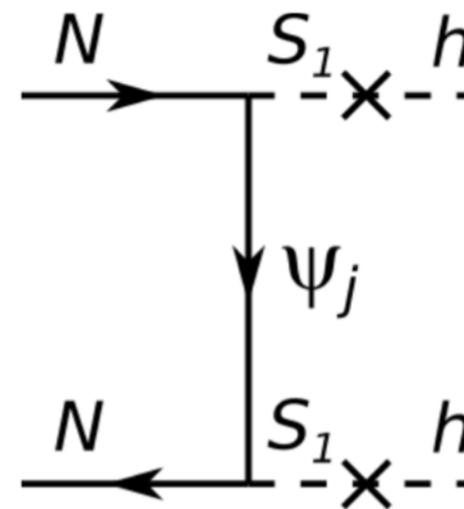

O(1) couplings with gears can be
used to make N a WIMP!

Abundance

(Secluded scenario)   (Higgs portal scenario)
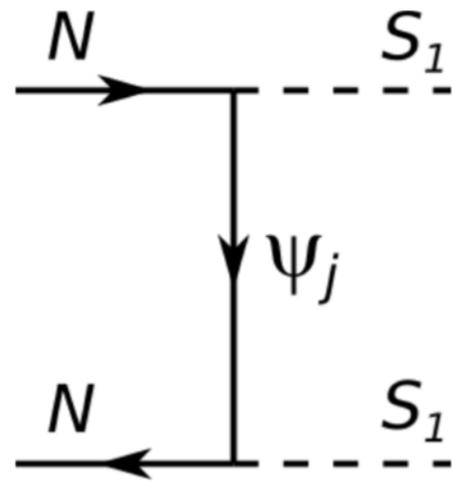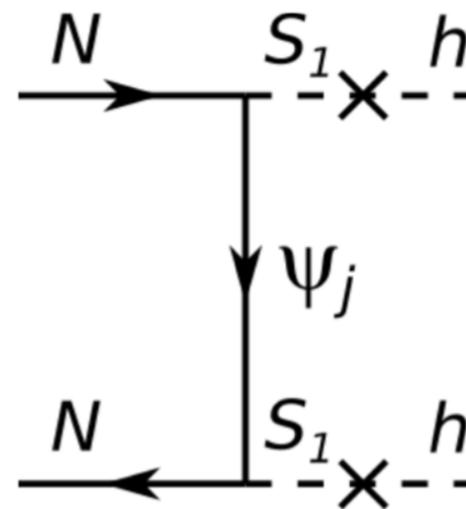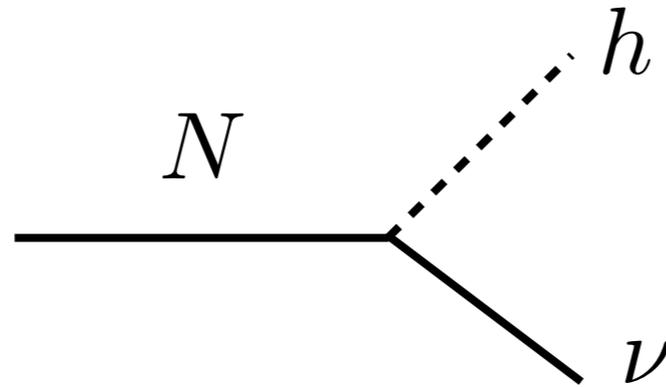
annihilation through the chiral gears

$$\langle \sigma v \rangle \sim 3 \cdot 10^{-26} \text{cm}^3 \text{s}^{-1}$$

Abundance

(secluded scenario)   (Higgs portal scenario)

annihilation through the chiral gears

$$\langle \sigma v \rangle \sim 3 \cdot 10^{-26} \mathrm{cm}^3 \mathrm{s}^{-1}$$

Stability/decay

$$\frac{1}{q^N}$$

suppressed by the chiral chain

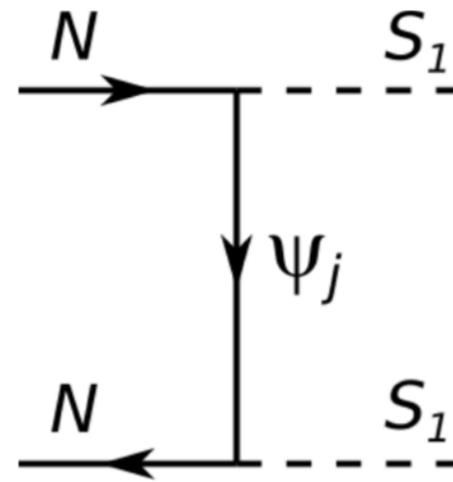👉 **Diagonal coupling to Higgs**

(taking into account the pseudo-Dirac nature of the clockwork gears)

$$\mathcal{L} \supset \frac{m_N}{\sqrt{2}m} \frac{\xi}{q^2} \bar{N}^c N \, h + h.c.$$
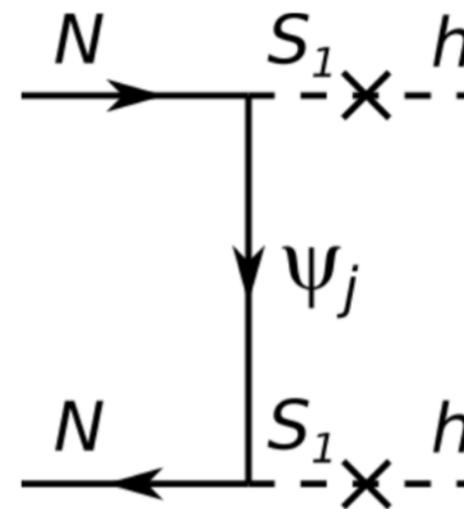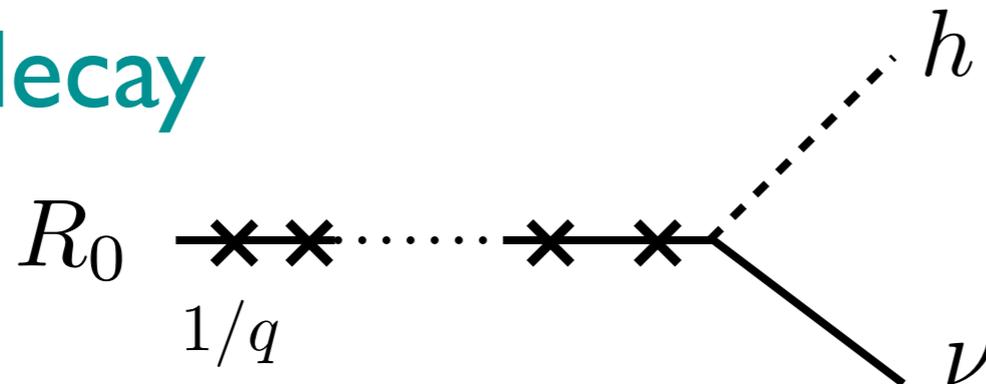
Abundance

(secluded scenario)   (Higgs portal scenario)

annihilation through the chiral gears

$$\langle \sigma v \rangle \sim 3 \cdot 10^{-26} \mathrm{cm}^3 \mathrm{s}^{-1}$$

Stability/decay

R_0

$1/q$

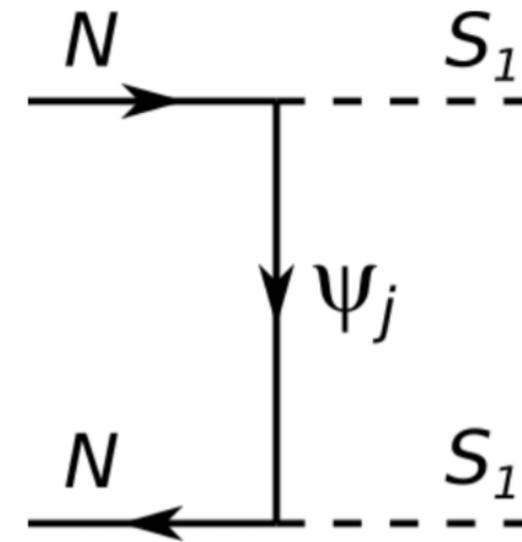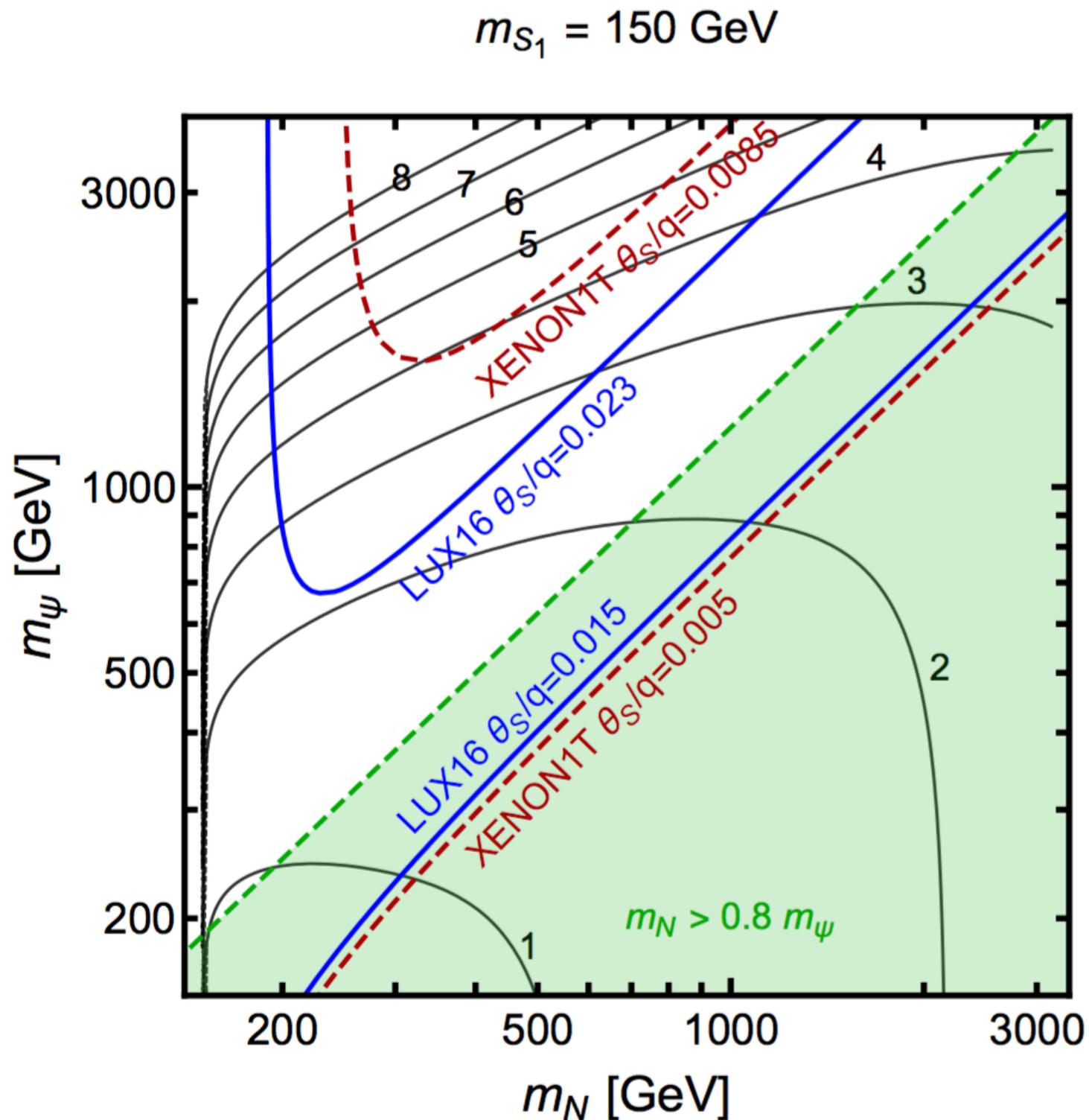$\dfrac{1}{q^N}$ suppressed by the chiral chain

Diagonal coupling to Higgs
(taking into account the pseudo-Dirac nature of the clockwork gears)

$$\mathcal{L} \supset \frac{m_N}{\sqrt{2}m} \frac{\xi}{q^2} \bar{N}^c N \, h + h.c.$$

$m_{S_1} = 150$ GeV

black solid: y_S required for $\Omega_{dm} \approx 0.25$

blue solid: LUX16 exclusions (below lines)

red dashed: Xenon1T reach

Higgs mediated $\propto (\theta_S/q)^2$

The clockwork gears ψ_i are essentially

a bunch of **TeV pseudo-Dirac sterile neutrinos**

with O(1) coupling to the SM leptons and the Higgs

<u>L conserving searches @ LHC</u>   **~ 3 charged leptons**

$$pp \to l_1^+ \Psi_i \to l_1^+ l_2^- W^+ \to l_1^+ l_2^- l_3^+ \nu \qquad m_\psi \lesssim 200\,\text{GeV} \quad (300\,\text{fb}^{-1})$$

Das, Dev & Okada (2014)

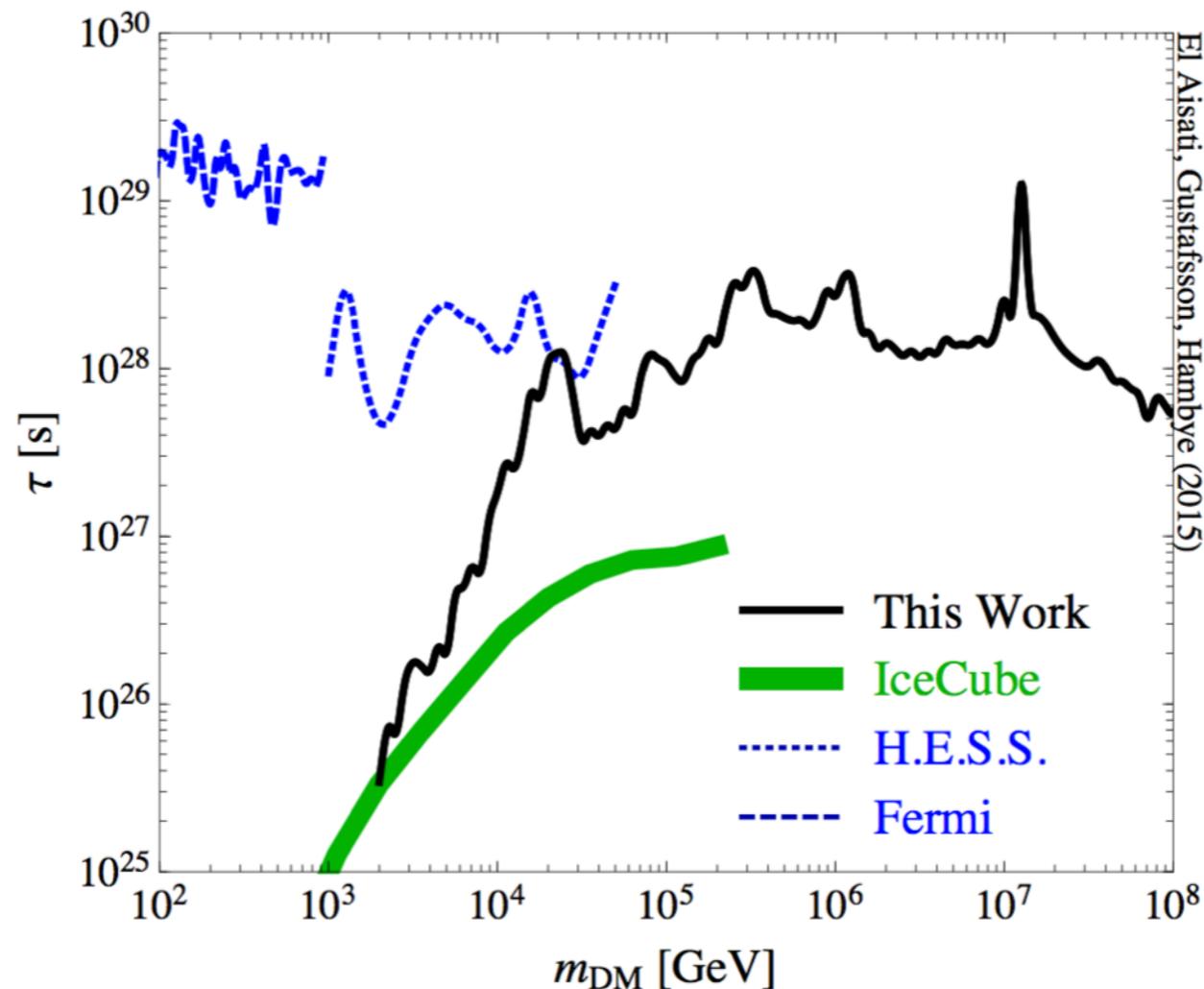<u>L violating searches @ LHC</u>   **~ 2 same sign leptons**

$$pp \to l_1^+ \Psi_i \to l_1^+ l_2^+ W^- \to l_1^+ l_2^+ \text{jets} \qquad m_\psi \lesssim 300\text{GeV} \quad (300\,\text{fb}^{-1})$$

Deppisch, Dev & Pilaftsis (2015)

DM indirect detection

p-wave annihilation, but decay into mono-energetic neutrinos

FIG. 8: Lifetime limits on DM particle decay into monochromatic lines $\nu + \gamma$. From Fermi-LAT [78] (blue, dashed) and H.E.S.S. [12, 28] (blue, dotted) using gamma-ray data compared to the neutrino line bounds derived in this study (solid black) as well as previous IceCube limits [31] (thick green).

$$N \rightarrow \nu + H$$

To recap, the <u>same</u> degrees of freedom that fix N abundance make N long-lived!

El Aisati, Gustafsson & Hambye (2015)

LOT OF FIELDS, WITH **VERY** SPECIFIC COUPLINGS…

☞ our Lagrangian is related to a discretized **flat** 5th dimension (Z)

1) $\mathcal{L}_5 \supset \bar{\psi}(i\overleftrightarrow{\partial\!\!\!/} - M)\psi = i\bar{\psi}\gamma^\mu \partial_\mu \psi + \left[\frac{1}{2}\left(\bar{L}\partial_Z R - \partial_Z \bar{L}R \right) - M\bar{L}R + h.c. \right]$

2) naive discretization $\Rightarrow$ fermion doubling

$\Rightarrow$ add a Wilson term $\quad -\dfrac{a}{2}\partial_Z\bar{\psi}\partial_Z\psi \quad$ with lattice spacing $\quad a = \dfrac{\pi R}{N} \to 0$

3) Dirichlet condition $\quad L(0) = 0 \quad \Rightarrow \quad$ surviving chiral mode $\quad R_0$

4) Standard Model degrees of freedom at $\quad Z = \pi R$ (~ braneworld)

$\mathcal{L} \supset \displaystyle\sum_{i=0}^{N-1} \frac{1}{a}\bar{L}_{i+1}R_i - \sum_{i=1}^{N}\left(\frac{1}{a} + M \right)\bar{L}_i R_i$

$\downarrow$

$m \qquad\qquad qm$

well-defined continuum limit

$q^N = \left(1 + \dfrac{\pi R M}{N} \right)^N \quad \to \quad e^{\pi R M}$

$1/M \ll \pi R$

Hambye, Teresi & MT (2016)

Clockwork (CW) scalar ☞ Linear Dilaton (LD)

warped, conformally flat 5D
$$ds^2 = e^{\frac{4}{3}k|y|}\left(\eta_{\mu\nu}dx^\mu dx^\nu + dy^2\right)$$

LargeExtraDim   <   CW/LD   <   Randall-Sundrum

**hierarchy
from volume**

$$M_P^2 \propto R M_5^3$$

**hierarchy
from warping**

Clockwork (CW) scalar ☞ Linear Dilaton (LD)

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hierarchy
from volume

$$M_P^2 \propto R M_5^3$$

**hierarchy
from
volume
& warping**

hierarchy
from warping

CONCLUSIONS

We have built a clockwork WIMP
It is accidentally stable (not protected by a symmetry)

**Its decay is mediated by the very same
degrees of freedom that determine its abundance**

These many degrees of freedom lie in the 100 GeV-TeV range

They could be seen @ LHC as sterile RHN

There are also (possibly) plenty of scalars to search for @ LHC

It is also a framework for low scale SM neutrino Majorana mass and
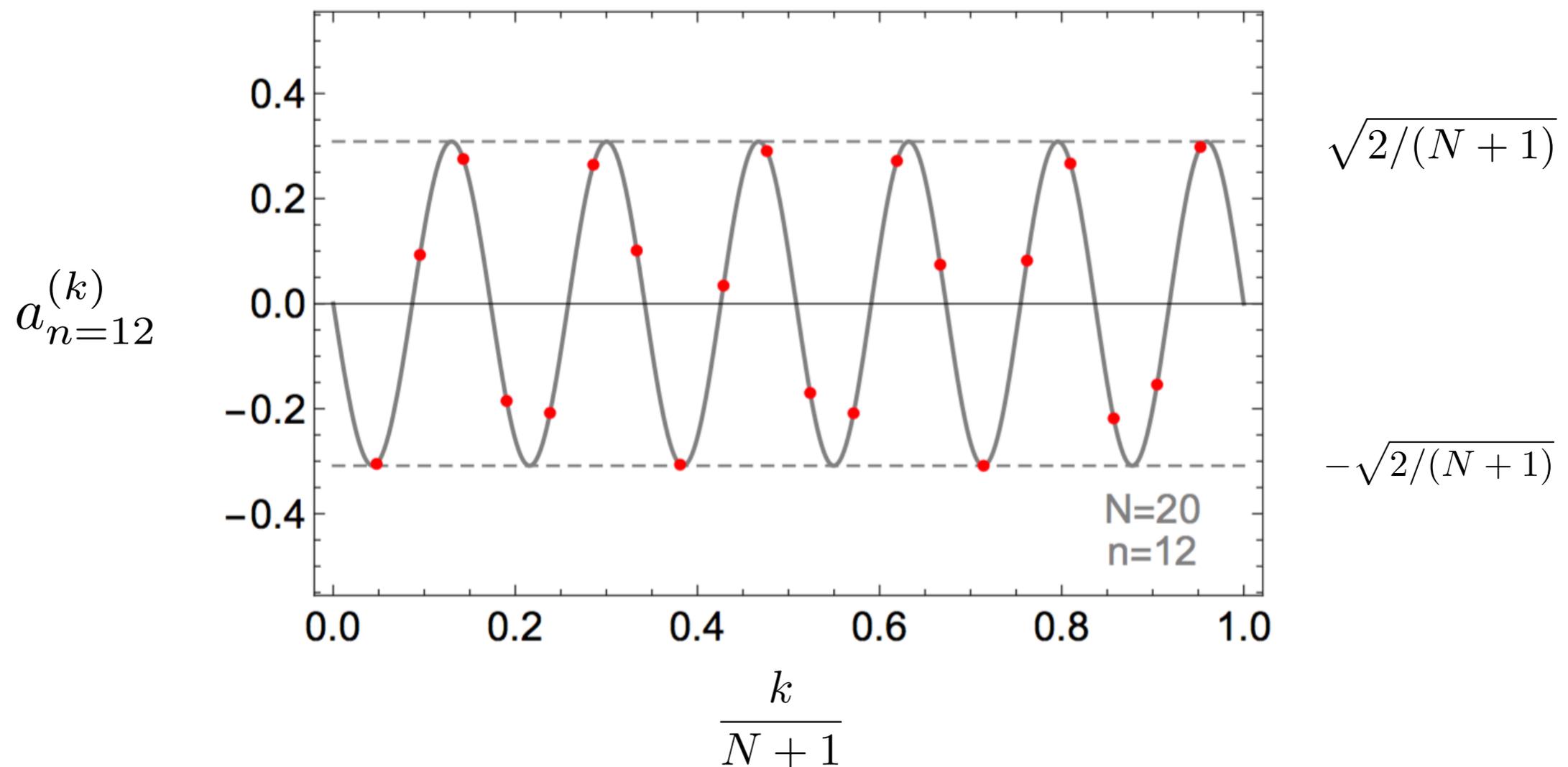potentially for much of BSM physics

The clockwork could be seen as a (discrete) spatial dimension
May lead to unusual signatures, which are worth being studied further

Backup slides

$$\varphi^{(k)} = \sum_{n=0}^{N} a_n^{(k)} \phi_n$$

**Clockwork Gears
have unsuppressed
overlap (eventually coupling) at all sites!**

$a_{n=12}^{(k)}$
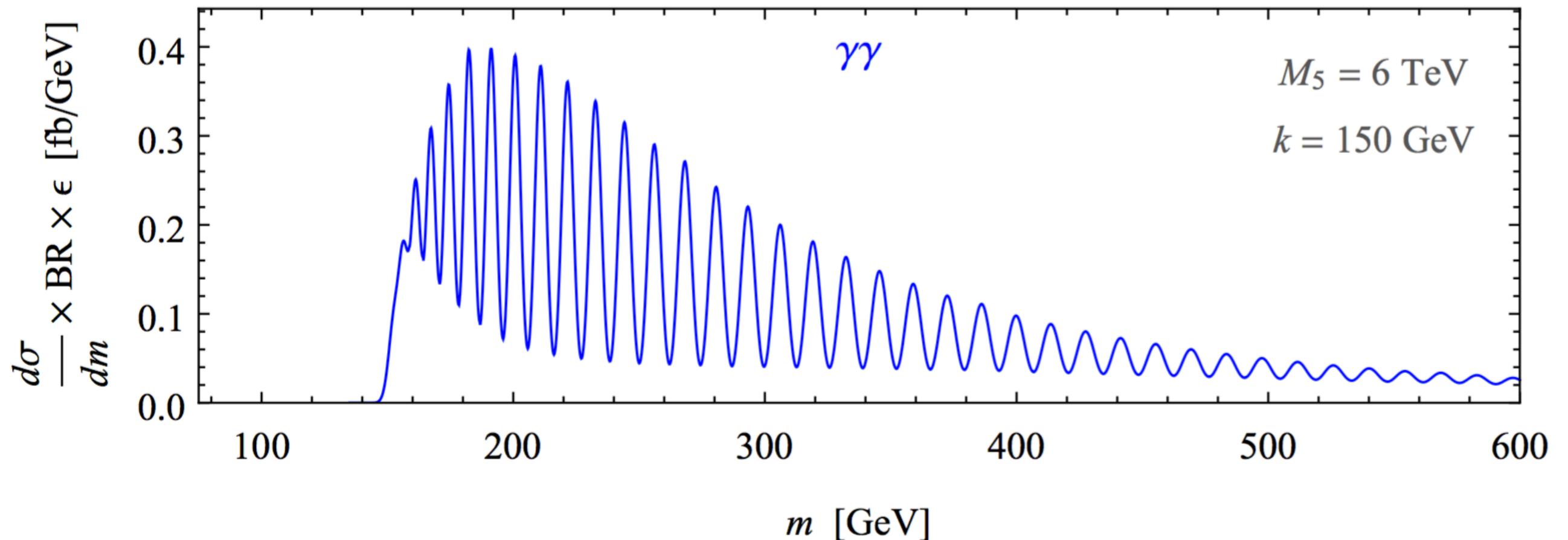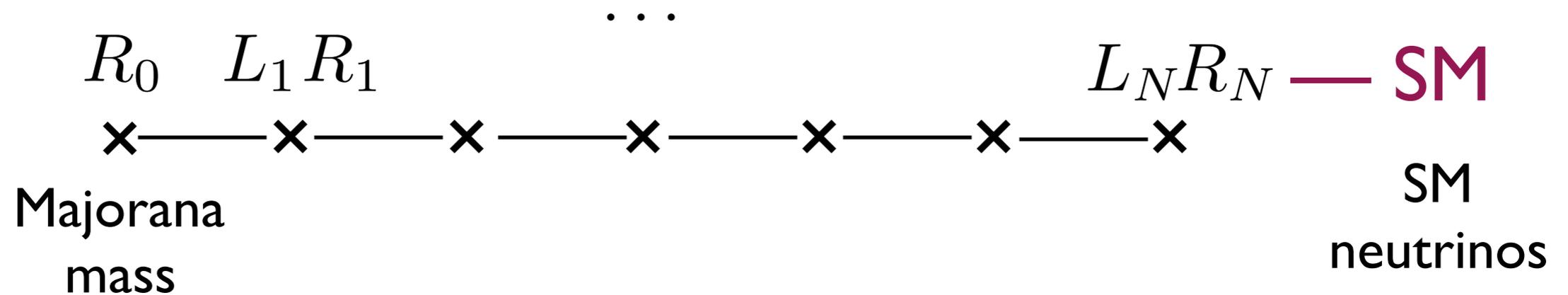
$\sqrt{2/(N+1)}$

$-\sqrt{2/(N+1)}$

N=20
n=12

$\dfrac{k}{N+1}$

plot from Farina, Pappadopulo, Rompineve & Tesi (2016)

Clockwork (CW) scalar ☞ Linear Dilaton (LD)

KK graviton → γγ

plot from Giuidice, Kats, McCullough, Torre & Urbano (2017)

another chain…

$$R_0 \quad L_1\, R_1 \qquad\qquad \cdots \qquad\qquad L_N R_N \text{ — SM}$$

×——×———×———×———×———×———×

Majorana
mass

SM
neutrinos

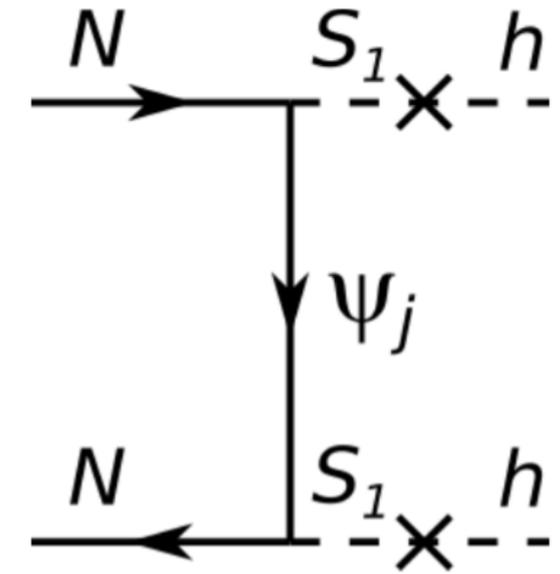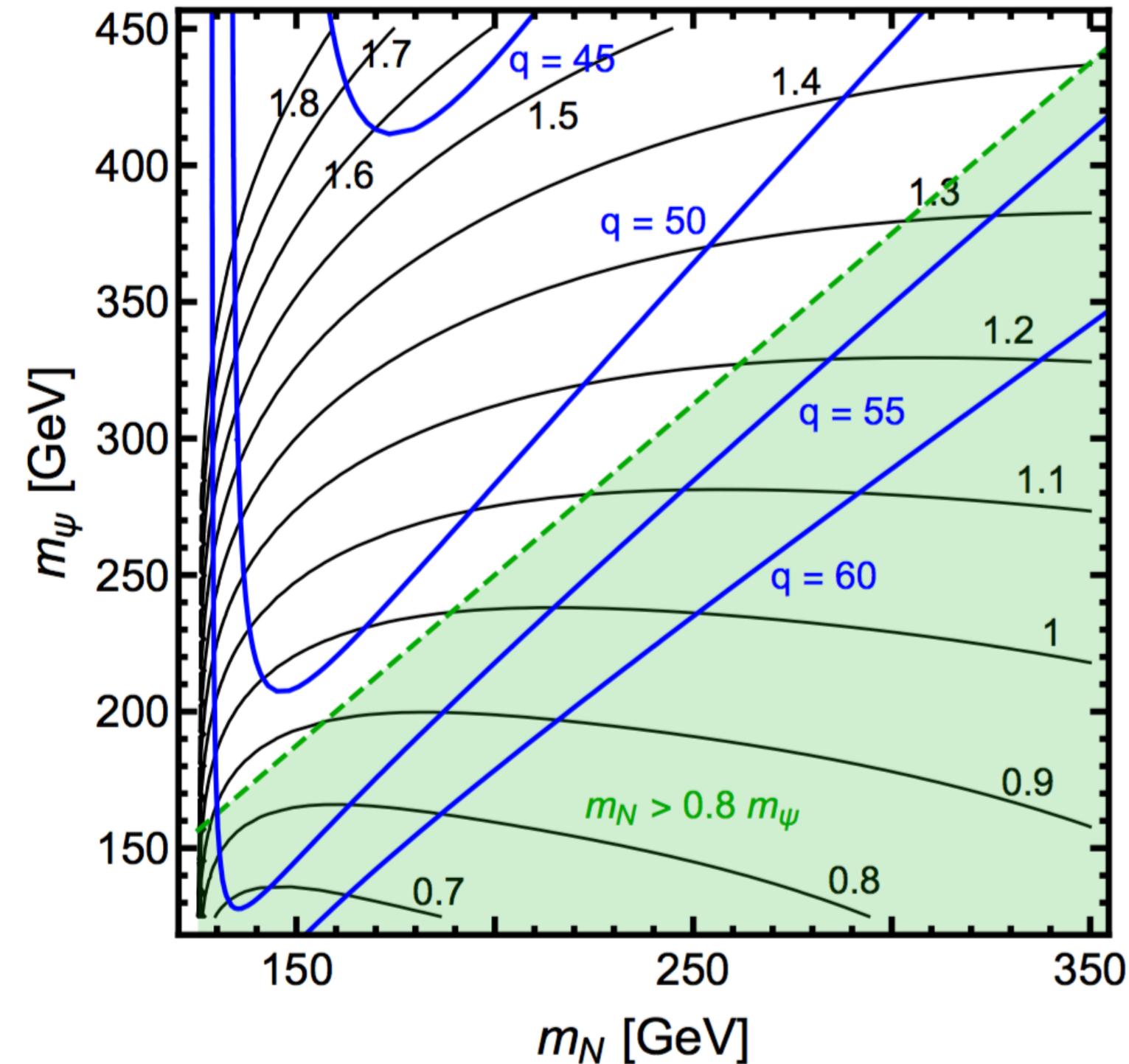👉 Majorana character has to go through
the whole chain

$$m_\nu \sim y^2 \frac{v^2}{q^N M}$$

Smaller chain than for DM stability $(q = 10, N = 7, M = 1\,\text{TeV})$

but need one per SM neutrino mass!

(i.e. at least 2 chains for neutrino, one for the DM…)

See also Carena, Li, Machado, Machado & Wagner (2017); Ibarra, Kushwaha & Vempati (2017)

$$\theta_S \lesssim \frac{0.4}{\sqrt{N}}$$

black solid: $y_S \theta_S$ required for $\Omega_{dm} \approx 0.25$

blue solid: LUX16 exclusions (below lines) (no time to put the new Xenon1T limits, sorry)